



# 1 **Contrasting the Penultimate and Last Glacial Maxima (140 and 21 ka** 2 **BP) using coupled climate-ice sheet modelling**

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10 **Abstract.** The configuration of the Northern Hemisphere ice sheets during the Penultimate Glacial Maximum differed to the  
11 Last Glacial Maximum. However, the reasons for this are not yet fully understood. These differences likely contributed to the  
12 varied deglaciation pathways experienced following the glacial maxima and may have had consequences for the interglacial  
13 sea level rise. Therefore, a better understanding of how and why these two glacial maxima differed is crucial for developing  
14 the full picture on why the Last Interglacial sea level was up to 9 meters higher than today, and thus may help constrain future  
15 sea level rise. To understand the differences between the North American Ice Sheet at the Last and Penultimate Glacial Maxima  
16 (21 and 140 ka BP), we perform two perturbed-physics ensembles of 62 simulations using a coupled climate-ice sheet model  
17 FAMOUS-ice, in which the North American and Greenland ice sheets are dynamically simulated with the Glimmer ice sheet  
18 model. We select six ensemble members that match reconstructed ice extent and volumes at the Last and Penultimate Glacial  
19 Maxima. To understand the role of orbit, greenhouse gases and initial conditions on the final ice sheet configurations, we use  
20 a factor decomposition technique. This reveals that the initial ice sheet conditions used in the model are extremely important  
21 in determining the difference in final ice volumes between both periods due to the large effect of the ice-albedo feedback. In  
22 contrast to evidence of a smaller Penultimate North American Ice Sheet, our model shows that the climate boundary conditions  
23 at these glacial maxima, if considered in isolation, imply a larger Penultimate Glacial Maximum North American Ice Sheet  
24 than at the Last Glacial Maximum, of around 6 meters sea level equivalent. This suggests the growth of the ice sheet prior to  
25 the glacial maxima is key in explaining the differences in North American ice volume.

## 26 **1 Introduction**

27 The Penultimate Glacial Maximum (PGM) occurred around 140,000 years ago, within Marine Isotope Stage 6 (MIS 6).  
28 Greenhouse gas (GHG) concentrations and global average insolation were similar to the Last Glacial Maximum (LGM; ~21  
29 ka BP) (Berger and Loutre, 1991; Loulergue et al., 2008; Bereiter et al., 2015) but the orbital configuration differed, affecting  
30 the seasonal and latitudinal distribution of incoming shortwave radiation (Berger, 1978; Colleoni et al., 2011). The global total



31 ice sheet volume, and thus the global mean sea level, was likely similar between the two glacial maxima (~120-130 m below  
32 present), with larger uncertainty at the PGM (Rabineau et al., 2006; Masson-Delmotte et al., 2010; Rohling et al., 2017). Both  
33 geological evidence and numerical modelling suggest that despite the similarities in total ice volume between the PGM and  
34 the LGM, the configurations of the Northern Hemisphere ice sheets differed significantly (e.g. Svendsen et al., 2004; Colleoni  
35 et al., 2016; Batchelor et al., 2019).

36 Some reconstructions suggest the Eurasian Ice Sheet (EIS) may have been up to ~50 % larger during the Penultimate Glacial  
37 Cycle (MIS 6: ~190-130 ka BP) than during the Last Glacial Cycle (~115-12 ka BP) (Svendsen et al., 2004), however evidence  
38 of multiple advances and uncertainties in dating proxy records means that the maximum extent mapped at 140 ka BP could  
39 correspond to previous advances during MIS 6 (Svendsen et al., 2004; Margari et al., 2014; Colleoni et al., 2016; Ehlers et al.,  
40 2018). The extent of the North American Ice Sheet (NAIS) during the PGM is even less well constrained due to a lack of  
41 glaciological evidence (e.g. moraines and till). The scarcity of empirical data in itself suggests that it was smaller in most areas  
42 than at the LGM because the subsequent larger ice sheet could have largely erased the evidence of prior glaciations (Dyke et  
43 al., 2002; Rohling et al., 2017). Additionally, evidence of reduced ice rafted debris (IRD) discharge from the Hudson Strait in  
44 the North Atlantic IRD belt (e.g. Hemming, 2004; Naafs et al., 2013; Obrochta et al., 2014), relative sea level assessment  
45 studies (e.g. Rohling et al., 2017) and climate, ice sheet and glacial isostatic adjustment modelling (e.g. Colleoni et al., 2016;  
46 Dyer et al., 2021) all point to a smaller volume PGM NAIS. For example, assuming a similar global mean sea level fall (and  
47 Antarctic ice sheet volume) at the PGM as at the LGM but with a larger volume EIS at the PGM (estimated at 33-53 m sea  
48 level equivalent (SLE) versus 14-29 m SLE at the LGM), this follows that the NAIS must have been smaller than at the LGM  
49 to compensate (39-59 m SLE versus 51-88 m SLE) (Rohling et al., 2017).

50 The reason for these differences is likely complex and is not yet fully understood. The evolution and surface mass balance  
51 (SMB) of ice sheets depends on many factors such as; background climate, climate and ice sheet histories, dust deposition,  
52 vegetation, ice albedo and sea surface temperatures, as well as the interactions and feedbacks between them all (Kageyama et  
53 al., 2004; Krinner et al., 2006; 2011; Colleoni et al., 2009a; 2011; Liakka et al., 2012; Stone and Lunt, 2013). The ice sheets  
54 themselves also strongly influence the climate through their interactions with atmospheric and oceanic circulation and the  
55 energy balance. This alters global and local temperature and precipitation patterns which in turn affects ice sheet ablation and  
56 accumulation (i.e. SMB) (e.g. Kageyama and Valdes, 2000; Abe-Ouchi et al., 2007; Beghin et al., 2014; 2015; Ullman et al.,  
57 2014; Liakka et al., 2016; Gregoire et al., 2015; 2018; Snoll et al., 2022; Izumi et al., 2023). These interactions between the  
58 vast ice sheets and other components of the climate system exerted an important control on the initial climate state for the  
59 deglaciations, and hence on the subsequent chain of events, thus impacting the climate, ocean and sea level evolution during  
60 deglaciation. Thus, the contrasting configurations of the Northern Hemisphere ice sheets at the glacial maxima may have  
61 contributed to the different deglaciation pathways that followed. The timings and magnitudes of the climate and ocean  
62 circulation changes that occurred during the Penultimate Deglaciation (~138-128 ka BP) differed to the Last Deglaciation  
63 (~19-11 ka BP) (Landais et al., 2013; Menviel et al., 2019). For example the Last Deglaciation experienced two abrupt climate



64 changes associated with a weakened Atlantic Meridional Overturning Circulation, Heinrich Stadial 1 and the Younger Dryas  
65 (Denton et al., 2010; Ivanovic et al., 2016; 2018), compared to evidence of only one, much longer abrupt change towards the  
66 end of the Penultimate Deglaciation, Heinrich Stadial 11 (Cheng et al., 2009; Govin et al., 2015; Marino et al., 2015; Jimenez-  
67 Amat and Zahn, 2015). The deglaciations also led to interglacials with very different characteristics to one another, including  
68 average global surface temperatures 1-2 °C higher and sea level up to 9 m higher than the pre-industrial during the Last  
69 Interglacial (~129-116 ka BP) (Kopp et al., 2009; Turney and Jones, 2010; Grant et al., 2012; Dutton and Lambeck, 2012;  
70 Otto-Bliesner et al., 2013; Dutton et al., 2015; Dyer et al., 2021).

71 In this context, it is important to examine the complex physical interactions between the climate and the ice sheets to better  
72 understand why the last two glacial maxima had different ice sheet configurations and evaluate the ice sheets' sensitivities to  
73 changes in climate in relation to different orbits and greenhouse gas concentrations.

74 Despite the challenges in coupling Atmosphere Ocean General Circulation Models (AOGCMs) with ice sheet models due to  
75 the mismatch between the required spatial and temporal scales, recent technical advances have meant that this is now possible.  
76 A combination of increased computer power, the development of more computationally efficient, lower resolution AOGCMs  
77 and sub-grid scale schemes translating ice sheet relevant atmospheric processes onto the higher resolution ice sheet grid, has  
78 made bi-directional, coupled climate-ice sheet simulations over longer timescales, and in large ensembles, feasible (Fyke et  
79 al., 2011; Vizcaino et al., 2013; Ziemer et al., 2014; Sellevold et al., 2019; Smith et al., 2021).

80 This study uses a coupled climate-ice sheet model, called FAMOUS-ice (Smith et al., 2021), to perform ensemble simulations  
81 of the PGM and LGM to explore input climate and ice sheet parameter uncertainties, their effects on the North American ice  
82 sheet volume during each period, and find parameter combinations that give a reasonable ice sheet configuration for both  
83 glacial maxima. The ensembles are also constrained based on volume and extent metrics and the 'Not Ruled Out Yet' (NROY)  
84 simulations are analysed to try and understand the similarities and differences between both periods. We find that the initial  
85 conditions used in the LGM and PGM experiments played an important role in some of the differences seen and we quantify  
86 this impact through the use of sensitivity tests and factor decomposition analysis.

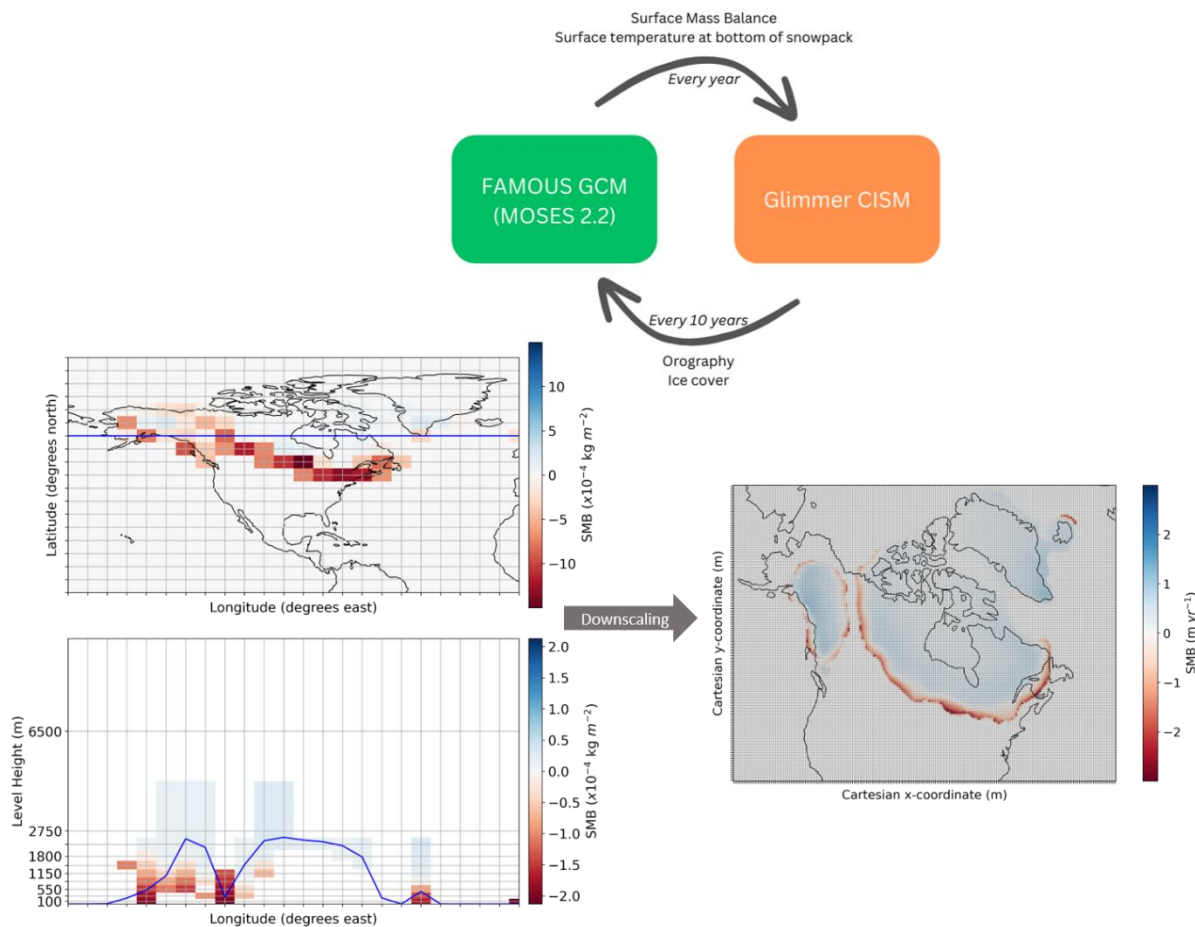
## 87 **2 Methods**

### 88 **2.1 Model description**

89 FAMOUS is a fast, low resolution AOGCM that is based on Hadley Centre coupled model HadCM3 and therefore retains all  
90 the complex processes represented in an AOGCM but uses only half the spatial resolution and a longer time step. Since it  
91 requires only 10 % of the computational costs of HadCM3, it has been successfully used for long transient palaeo simulations  
92 (Smith and Gregory, 2012; Gregory et al., 2012; Gregoire et al., 2012; Roberts et al., 2014; Dentith et al., 2019) and large  
93 ensembles for uncertainty quantification (Gregoire, 2010; Gandy et al., 2023). This study uses the atmospheric component,  
94 which is a quasi-hydrostatic, primitive equation grid point model with a horizontal resolution of 7.5° longitude by 5° latitude



95 with 11 vertical levels and a 1-hour time step (Williams et al., 2013). Land processes are modelled using the MOSES2.2 land  
96 surface scheme (Essery et al., 2003), which uses a set of sub-gridscale tiles in each grid box to represent fractions of nine  
97 different surface types, including land ice (Smith et al., 2021). Whilst this study prescribes sea surface temperatures and sea  
98 ice concentrations, FAMOUS can also be run fully coupled with a dynamical ocean (e.g. Dentith et al., 2019).  
99 FAMOUS now allows the direct two way coupling to an ice sheet model in the configuration FAMOUS-ice (Smith et al.,  
100 2021). Here, we use FAMOUS in combination with Glimmer to interactively simulate the North American and Greenland ice  
101 sheets at 40km resolution. Glimmer is a fast running, 3D thermomechanical ice sheet model which uses the shallow ice  
102 approximation. This allows it to model ice sheet evolution over long timescales as it is more computationally efficient, and  
103 therefore has been used to simulate continental ice sheets over glacial-interglacial cycles (Rutt et al., 2009; Gregoire et al.,  
104 2016).  
105 FAMOUS-ice accounts for the mismatch between atmosphere and ice sheet grid sizes by using a multilayer surface snow  
106 scheme to calculate SMB on ‘tiles’ at 10 set elevations within each grid box that contains land ice in FAMOUS. This SMB is  
107 then downscaled from the coarse FAMOUS grid to the much finer Glimmer grid at each model year (Smith et al., 2021).  
108 Glimmer uses this SMB field to calculate ice flow and surface elevation and passes this back to FAMOUS in which orography  
109 and ice cover is updated. In this study, to reduce computational costs further, FAMOUS-ice runs at 10 times ice sheet  
110 acceleration: for every year of climate integrated in FAMOUS, the simulated SMB field forces 10 years of ice sheet integration  
111 in Glimmer. Figure 1 shows a simplified diagram of this coupling process and full details can be found in Smith et al.,  
112 (2021). The current computational cost of this set up is around 50 decades (of climate years) per wallclock day using 8  
113 processors.  
114 FAMOUS-ice has been shown to perform well in simulations of past and future ice sheets including Greenland and North  
115 America (Gregory et al., 2020; Smith et al., 2021; Gandy et al., 2023). In particular, the LGM North American Ice Sheet study  
116 of Gandy et al., (2023) was able to utilise the useful constraints of the LGM to infer the importance of parameters controlling  
117 ice sheet albedo on ice sheet configuration in this model.



118

119 **Figure 1. Schematic illustrating the calculation of SMB at different elevations on the FAMOUS grid followed by downscaling onto**  
120 **the Glimmer grid.**

## 121 2.2 Experiment design

122 Our FAMOUS-ice simulations are set up following the Paleoclimate Modelling Intercomparison Project Phase 4 (PMIP4)  
123 protocols for the LGM (Kageyama et al., 2017) and PGM (Menviel et al., 2019). These protocols prescribe climatic boundary  
124 conditions, including orbital parameters and GHG concentrations, the values of which can be found in Table 1. Concentrations  
125 of  $\text{CO}_2$ ,  $\text{CH}_4$  and  $\text{N}_2\text{O}$  are very similar between the LGM and PGM but orbital parameters are significantly different. The larger  
126 eccentricity at the PGM enhances the effect of precession compared to the LGM which affects the seasonal and latitudinal  
127 distribution of insolation. These changes are important for ice sheet surface mass balance since melting is particularly sensitive  
128 to spring and summer temperatures (Huybers, 2006; Niu et al., 2019). The PGM received lower insolation in the Northern  
129 Hemisphere in late winter to early summer but higher levels in late summer to early winter, compared to the LGM (Fig. 2a).  
130 Subsequent to the completion of this work, it was discovered that the equation for the role of eccentricity on solar insolation



131 was incorrect in the model code. The magnitude of the error is larger for periods with higher eccentricity values and so a  
 132 sensitivity test was run to determine the effect this correction has on SMB and ice volume at the PGM. Details of this error  
 133 and the results of the sensitivity test can be found in Appendix A, but the impact was shown to be minimal (Fig. A1).

134  
 135 **Table 1. Climatic boundary conditions used in the LGM and PGM experiments as prescribed by the PMIP4 protocols for each**  
 136 **period (Kageyama et al., 2017; Menviel et al., 2019).**

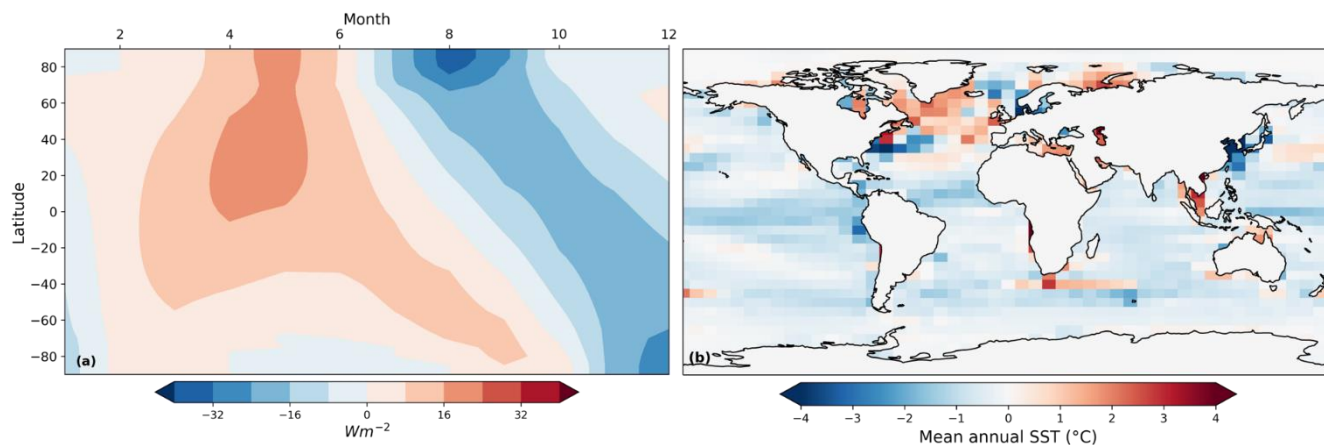
	<i>Eccentricity</i>	<i>Obliquity</i> (°)	<i>Perihelion</i> – 180 (°)	<i>Solar Constant</i> ( $Wm^{-2}$ )	<i>CO<sub>2</sub></i> (ppm)	<i>CH<sub>4</sub></i> (ppb)	<i>N<sub>2</sub>O</i> (ppb)
<b>LGM</b> <b>(21 ka)</b>	0.019	22.949	114	1360.7	190	375	200
<b>PGM</b> <b>(140 ka)</b>	0.033	23.414	73	1360.7	191	385	201

137  
 138 In the climate model, the global orography (including ice sheets) and land-sea mask for the LGM are calculated from the  
 139 GLAC1D 21 ka BP reconstruction (Tarasov et al., 2012) which is one of the two options in the PMIP4 protocol (Kageyama et  
 140 al., 2017). For the PGM simulations we used the 140 ka BP combined ice sheet reconstruction (Tarasov et al., 2012; Abe-  
 141 Ouchi et al., 2013; Briggs et al., 2014) detailed in the PGM PMIP4 protocol (Menviel et al., 2019). Vegetation is prescribed  
 142 based on a pre-industrial distribution and kept constant. As ice cover changes, the fractions of grid cells that are land ice versus  
 143 other surface types changes proportionally, altering albedo. However, since there is no dynamical vegetation component, some  
 144 important climate-ice-vegetation feedbacks are neglected which could have a significant impact on ice sheet evolution (Stone  
 145 and Lunt, 2013). Sea Surface Temperature (SST) and sea ice concentration is also prescribed and constant and is taken from  
 146 HadCM3 simulations of 21 ka BP and 140 ka BP (Figs. B1 and B2). The modelled annual average SSTs are cooler at the LGM  
 147 than at the PGM, everywhere, except in the North Atlantic due to less sea ice cover in this region (Fig. 2b). However, the  
 148 summer SSTs are warmer in the NH at the LGM compared to the PGM. The decision to use these constant SST and sea ice  
 149 fields, rather than a statistical reconstruction as in Gandy et al., (2023), was made due to the lack of both empirical and modelled  
 150 PGM SST data available to produce an equivalent reconstruction. The HadCM3 SSTs are colder on average than the  
 151 reconstruction in Gandy et al., (2023), with the largest differences, of up to 6 °C, seen in the tropics and mid-latitudes. This  
 152 may introduce another source of uncertainty in the simulations. In the ice sheet model, we use the same ice sheet domain and  
 153 initial condition for the LGM and PGM, which is the same as used in Gandy et al., (2023). The interactive ice sheet model  
 154 domain in Glimmer covers North America and Greenland, and the initial ice sheet extent, thickness and bedrock elevation is  
 155 from a previous Last Deglaciation ensemble of the NAIS, at 18.2 ka BP (Gregoire et al., 2016). This is a smaller intermediate





156 (MIS 3 like) ice sheet which is used as an approximate pre-glacial maximum extent from which to grow the ice sheet towards  
157 an equilibrium ice volume.  
158



159  
160 **Figure 2. Difference between the LGM and PGM (a) latitudinal distribution of incoming top of the atmosphere shortwave radiation**  
161 **each month, and (b) modelled sea surface temperatures.**

### 162 2.3 Ensemble design

163 The ensemble by Gandy et al., (2023) showed that uncertainty in parameters controlling SMB, ice sheet dynamics and climatic  
164 conditions over the ice sheets had a significant influence on the extent and volume of the LGM NAIS, with albedo parameters  
165 explaining the majority of the variation in model output. Since these parameters needed re-tuning from simulations of the  
166 present day Greenland ice sheet to produce an acceptable LGM NAIS configuration in FAMOUS-ice, the PGM may also show  
167 different sensitivities to the uncertain parameters. Therefore, we run new ensembles of the LGM and PGM in order to quantify  
168 uncertainties and identify combinations of climate and ice sheet parameters that perform well for both periods.

169 Following on from Gandy et al., (2023), a second wave of simulations was performed and compared to reconstructions of ice  
170 sheet extent and volume to identify NROY parameter combinations (see methodology in Appendix C), the results of which  
171 formed the basis of the ensemble design in this study. We reran the LGM ensemble to allow for slight changes in the experiment  
172 design compared to Gandy et al., (2023): we use orbital parameters for 21 ka BP rather than 23 ka BP and HadCM3 SSTs  
173 instead of a reconstruction (see Sect. 2.2). Table 2 details the 13 parameters that were varied in these simulations. Out of the  
174 176 NROY parameter combinations from the Wave 2, a representative subset of 62 were selected which provided adequate  
175 coverage of the NROY space (see Appendix C for details). Each was run for 1000 climate years (10,000 ice sheet years) for  
176 both the LGM and PGM set ups until the majority of the ice sheet has reached close to equilibrium. Despite differences in the  
177 model set up between this study and Gandy et al., (2023), we expect the 62 samples chosen from their design to be a good  
178 estimate to an optimal parameter design for our set up (Appendix C).



179

180

**Table 2. Description of parameters varied in the ensembles. Adapted from Gandy et al., (2023).**

<i>Parameter</i>	<i>Description</i>
<i>Lapse Rate</i>	Prescribed lapse rate for air temperature used to downscale FAMOUS near-surface ice sheet climate onto surface elevation tiles. Downwelling longwave radiation is also adjusted for consistency.
<i>Daice</i>	Sensitivity of bare-ice albedo to surface air temperatures once the surface is in a melt regime.
<i>Rho</i>	The threshold in surface snow density at which the FAMOUS albedo scheme switches from a scattering paradigm appropriate for a conglomeration of snow grains to one more appropriate for a solid surface.
<i>AV_GR</i>	Sensitivity of the snow albedo to variation in surface grain size.
<i>RHcrit</i>	The threshold of relative humidity for cloud formation (R. Smith, 1990). A higher value means clouds can form less easily.
<i>VFI</i>	The precipitating ice fall-out speed (Heymsfield, 1977).
<i>CT</i>	The conversion rate of cloud liquid water droplets to precipitation (R. Smith, 1990).
<i>CW</i>	The threshold values of cloud liquid water for formation of precipitation (R. Smith, 1990). Only the value for the land is varied.
<i>Entrainment Coefficient</i>	Convection Scales rate of mixing between environmental air and convective plume.
<i>Alpham</i>	The sea ice low albedo (Crossley & Roberts, 1995).
<i>Basal Sliding</i>	The basal sliding rate. A higher value allows increased ice velocity.
<i>Mantle Relaxation Time</i>	The relaxation time of the mantle, a lower value essentially making the mantle less viscous, thus allowing a quicker topographic rebound.
<i>Flow Enhancement Factor</i>	The softness of ice. Increasing the factor makes the ice softer and more deformable (Rutt et al., 2009).





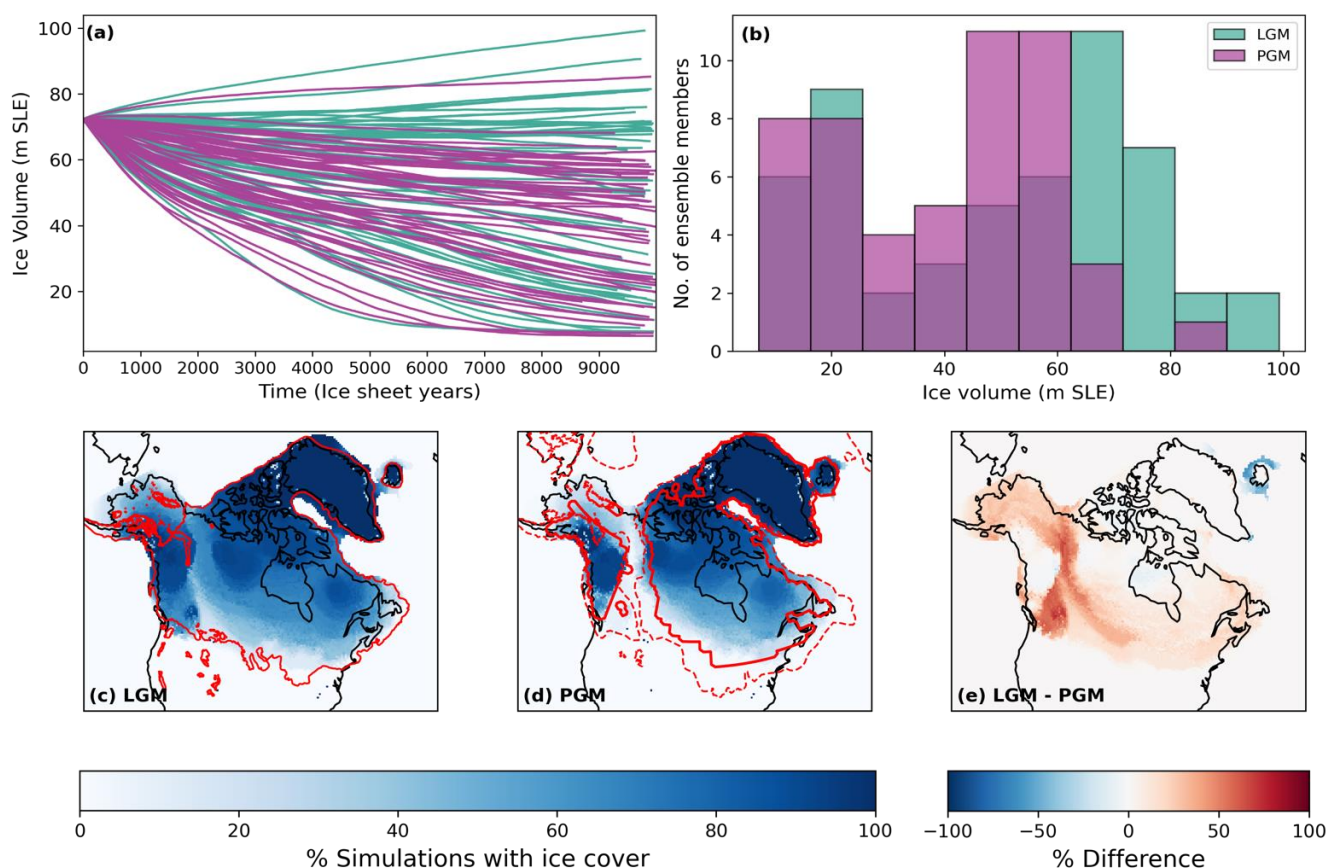
### 181 3 Results and discussion

#### 182 3.1 Ensembles

##### 183 3.1.1. Unconstrained ensembles

184 Our ensembles of 62 North American Ice Sheet configurations spans uncertainty in model parameters and reveals the wide  
185 range of possible modelled ice sheet evolutions. Over the full ensembles, we find that the set-up of the original Wave 2 meant  
186 that the albedo values were too high and so the use of more realistic albedos in these ensembles led to many of the runs  
187 deglaciating to very low volumes as shown in Fig. 3 (see Appendix C for more detail).

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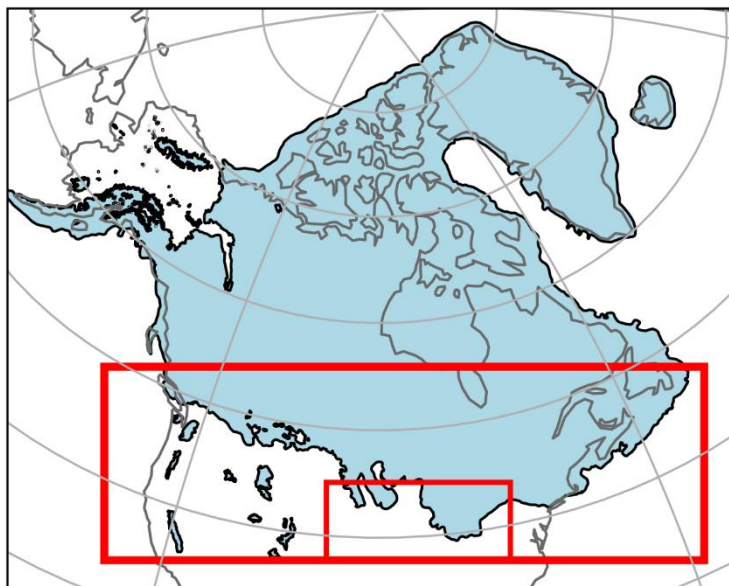
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190 **Figure 3.** (a) Ice volume evolution over modelled time, and (b) density distribution of final ice volumes for the full LGM and PGM  
191 ensembles. Percentage of simulations with ice cover for (c) LGM (with the Dalton et al., (2020) reconstructed margin shown in red);  
192 (d) PGM (with the PMIP4 PGM modelled margin shown in solid red and the Batchelor et al., (2019) reconstructed maximum MIS  
193 6 margin shown in dashed red), and (e) the difference between the LGM and PGM, at the end of the simulations.



### 194 3.1.2. Constrained ensembles

195 To filter out implausible ice sheet configurations in the results, a set of constraints, based on southern ice sheet extent and  
196 volume, were applied to the LGM ensemble. Both ensembles were filtered based on the LGM results since the extent of the  
197 NAIS is very well constrained by geological data and there are more estimates of ice volume for the LGM than the PGM. This  
198 is because there is a lack of empirical data (over both space and time) on ice sheet configuration at the PGM due to destruction  
199 by subsequent glaciations and difficulties with dating what is available (Parker et al., 2022). Thus, most of the reconstructions  
200 of NAIS extent for this period are actually the maximum extent over the whole of MIS 6 (190-132 ka BP) and are mostly based  
201 on numerical modelling combined with this scarce proxy data (e.g. Colleoni et al., 2016; Batchelor et al., 2019). The NROY  
202 LGM results can then be compared to the corresponding PGM results to advance understanding of the differences that occurred  
203 and reveal whether parameters that performed well for the LGM also give plausible PGM results. Ice extent was assessed  
204 against the reconstruction by Dalton et al. (2020). We focus our evaluation of ice extent on the southern NAIS area and chose  
205 to disregard regions of known model bias. This includes marine margins that are subject to processes not included in Glimmer  
206 and the Alaskan regions where small climate model biases lead to ice sheet overgrowth (e.g. Ganopolski et al., 2010; Gregoire  
207 et al. 2016, Sherriff-Tadano et al., 2023). Additionally, ice lobes are not well captured in many models so we do not expect  
208 our simulations to perfectly match the reconstructed Southern NAIS extent. To account for this, we applied a tolerance on the  
209 Southern ice sheet area of  $1.79 \times 10^6 \text{ km}^2$ , equivalent to 3 times the area of the lobes (Fig. 4). We thus calculate the Southern  
210 NAIS ice area as the integrated area within the large box shown in Fig. 4 at the end of each LGM simulation and selected  
211 simulations that matched the reconstructed area from Dalton et al. (2020) within plus or minus  $1.79 \times 10^6 \text{ km}^2$ . The volume of  
212 the NAIS is not as well constrained by proxy data and so estimates rely on ice sheet, glacial isostatic adjustment and sea level  
213 modelling studies. Based on a number of these studies, a minimum NAIS (including Greenland) volume of 70 m SLE ( $2.8 \times$   
214  $10^7 \text{ km}^3$ ) was applied to the ensemble (Marshall et al., 2002; Tarasov and Peltier, 2002; 2004; Tarasov et al., 2012; Lambeck  
215 et al., 2014; Peltier et al., 2015; Rohling et al., 2017; Batchelor et al., 2019; Gowan et al., 2021). The volumes in meters sea  
216 level equivalent are calculated based on present day ocean area.



217  
 218 **Figure 4. Outline of the LGM North American Ice Sheet by Dalton et al. (2020). The large red box shows the region we use to**  
 219 **calculate reconstructed and modelled Southern NAIS area. The small red box shows the region used to calculate the area of the lobes**  
 220 **for setting as the upper and lower target bounds.**

221 **Table 3. Average volumes (NAIS + Greenland) and southern NAIS areas and their standard deviations (SD) of the NROY LGM and**  
 222 **PGM simulations. Also shown are estimated values from literature for comparison.**

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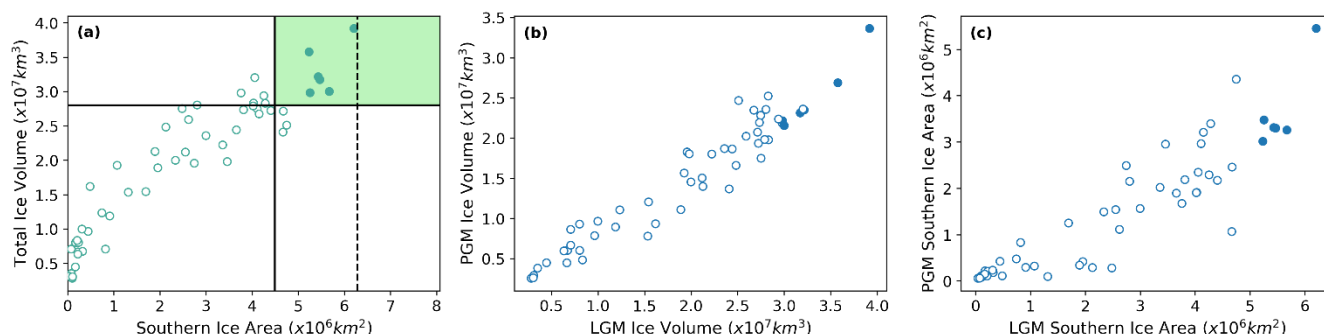
	<i>Mean Total Volume (SD), m SLE</i>	<i>Estimated Total Volume, m SLE</i>	<i>Mean Southern Area (SD), x 10<sup>6</sup> km<sup>2</sup></i>	<i>Estimated Southern Area, x 10<sup>6</sup> km<sup>2</sup></i>
<b>LGM</b>	82.1 (8.29)	61-98 (Rohling et al., 2017)	5.55 (0.33)	6.28 (Dalton et al., 2020)
<b>PGM</b>	62.3 (10.3)	49-69 (Rohling et al., 2017)	3.64 (0.82)	3.32 (Menviel et al., 2019)

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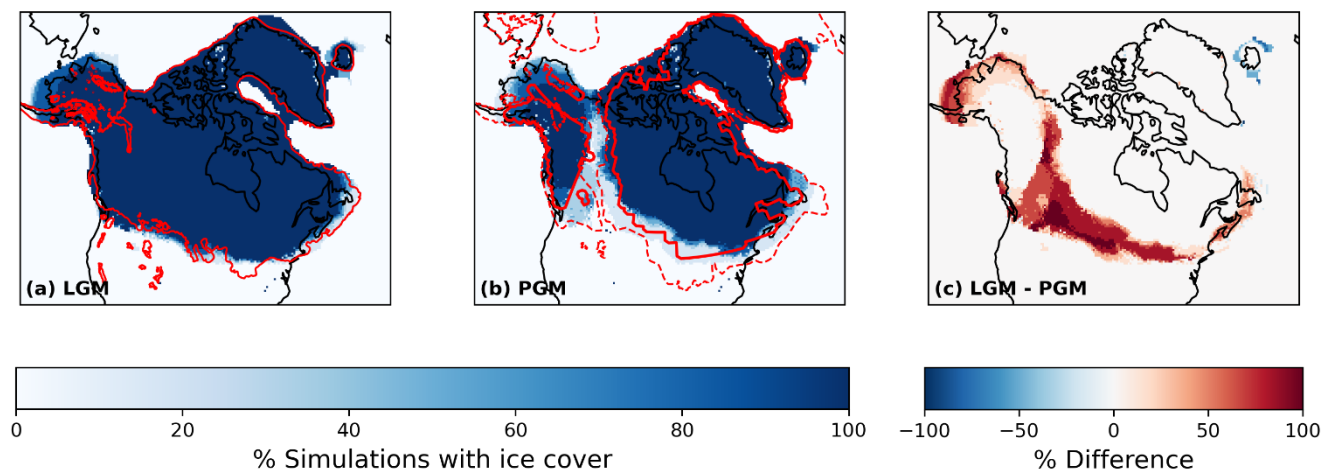
225 After applying our metric constraints, six non-implausible or NROY LGM simulations remained. Table 3 gives the average  
 226 volumes and areas of these six simulations and the corresponding six PGM ice sheets compared to estimated values from  
 227 empirical and model data. All six LGM simulations show an overgrowth of ice in Alaska of varying magnitudes, as a result of  
 228 the previously mentioned climate model bias. However, in other regions the simulations display a very similar ice extent, with  
 229 the southern area only varying by  $9.7 \times 10^5 \text{ km}^2$ . None of the simulations form ice lobes as expected but show a close match  
 230 to reconstructed ice extent in our target area and in the marine regions (Fig. 6a). There is a minimum ice volume of 73.9 m



231 SLE and a maximum of 97.1 m SLE. The maximum ice thickness varies by around 300 m but the overall shapes of the ice  
232 sheets remain the same, with the thickest ice towards the east of the ice sheet over Hudson Bay.  
233 All the PGM ice sheets were smaller in volume than their LGM counterpart (Figs. 5 and 6) and displayed a smaller extent in  
234 the southern margin and the saddle region between the western Cordilleran Ice Sheet and eastern Laurentide Ice Sheet.  
235 However, the PGM simulations also displayed more variability in their ice extent and volumes. The ice volumes range from  
236 53.4 m SLE to 83.37 m SLE and the southern extent varies by  $2.44 \times 10^6 \text{ km}^2$ . The range in maximum ice thickness is also  
237 over double the LGM, varying by around 613 m. These PGM configurations also look plausible compared to the less well  
238 constrained extent data available, including previous empirical and modelled reconstructions of the PGM/MIS 6 extent  
239 (Menviel et al., 2019; Batchelor et al., 2019; Fig. 6b). For example, all the simulations maintain an ice-free corridor between  
240 the Laurentide and Cordilleran ice sheets which is a common feature in these PGM reconstructions. In addition, the excess  
241 Alaskan ice seen in LGM simulations is also present at the PGM, however the growth is not as excessive. We therefore  
242 conclude that in our model, based on the available empirical constraints, parameters that produce a good LGM NAIS also  
243 produce a plausible PGM NAIS using PGM boundary and initial conditions (orbital parameters, SSTs and orography). Our  
244 simulations can thus be compared and analysed to understand the causes of the different configurations between the two  
245 periods.  
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248 **Figure 5. (a) The relationship between final ice volume and southern area for the LGM ensemble, and the relationship between the**  
249 **LGM and PGM; (b) final ice volume, and (c) final southern areas. The filled in blue dots represent the six NROY LGM simulations**  
250 **and the solid lines on panel (a) show the minimum volume and area constraints applied to the ensemble.**



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**Figure 6.** Percentage of simulations with ice cover for (a) LGM with the Dalton et al., (2020) reconstructed margin shown in red; (b) PGM with the PMIP PGM modelled margin shown in solid red and the Batchelor et al., (2019) reconstructed maximum MIS 6 margin shown in dashed red, and (c) the difference between the LGM and PGM, at the end of the simulations for the six NROY ensemble members.

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### 3.2 Uncertainty due to model parameters

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Most of the uncertainty in the results for both the LGM and PGM can be explained by parameters that affect the surface albedo of the ice sheet (*Daice*, *Rho* and *AV\_GR*) and *basal sliding* (Fig. 7). Similar conclusions were drawn by Gandy et al., (2023), on which this study is based, as well as other ensemble based studies exploring the sensitivity of the LGM NAIS to model parameters (e.g. Sherriff-Tadano et al., 2023). The similar behaviour between the LGM and PGM across the parameter ranges (Figs. D1 and D2) further implies that similar model parameter values are appropriate for use when modelling both periods and within the bounds of available model and data constraints, our results show that retuning the model would not lead to significant changes in predicted ice sheet configurations between the LGM and PGM. However, since the ice volume is most sensitive to surface albedo and most simulations deglacierate under low values of *Daice*, this suggests that the value of bare ice albedo in the model may need to be increased for future work.

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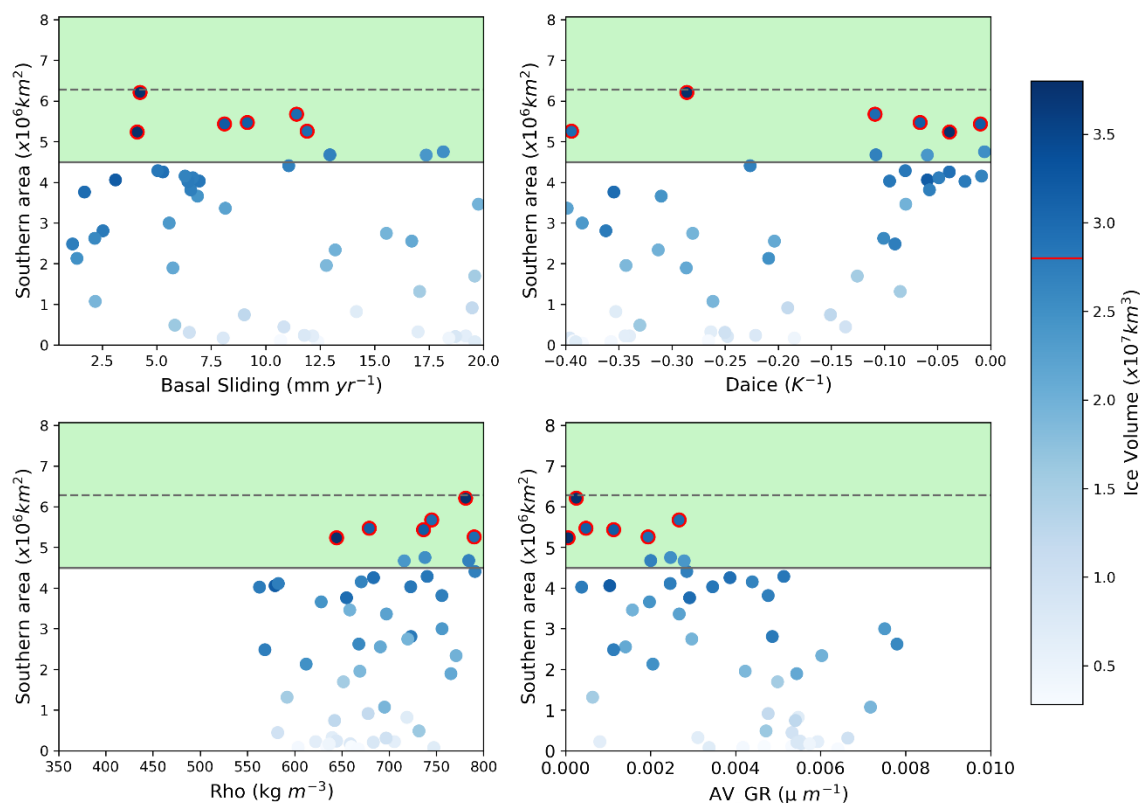
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Additionally, the difference in ice volume and area between the LGM and PGM is most influenced by the parameters *AV\_GR*, *Daice* and *basal sliding*, however the effect of these parameters on the differences seen is minor (Fig. D3). This suggests that the higher the albedo and lower the ice sheet velocity, the more sensitive the ice sheet is to changes in radiative forcings from the orbital boundary conditions. Due to the sampling strategy, this ensemble is not the best design to analyse the sensitivity of the ice sheets during the two time periods to the different parameters and would require a larger ensemble and a sensitivity analysis with Gaussian Process emulation (e.g. Pollard et al., 2023), as is presented in Gandy et al. (2023) and Sherriff-Tadano et al. (2023).



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**Figure 7. Relationship between LGM southern area and the four most influential parameters. The green shaded region shows the southern area constraint applied with the dotted line showing the exact area of the reconstruction and the solid line the minimum bound applied. The colour scale represents ice volume and the dots outlined in red are the six NROY LGM simulations with the red line on the colour bar showing the volume constraint.**

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### 3.3 Climate-ice sheet interactions

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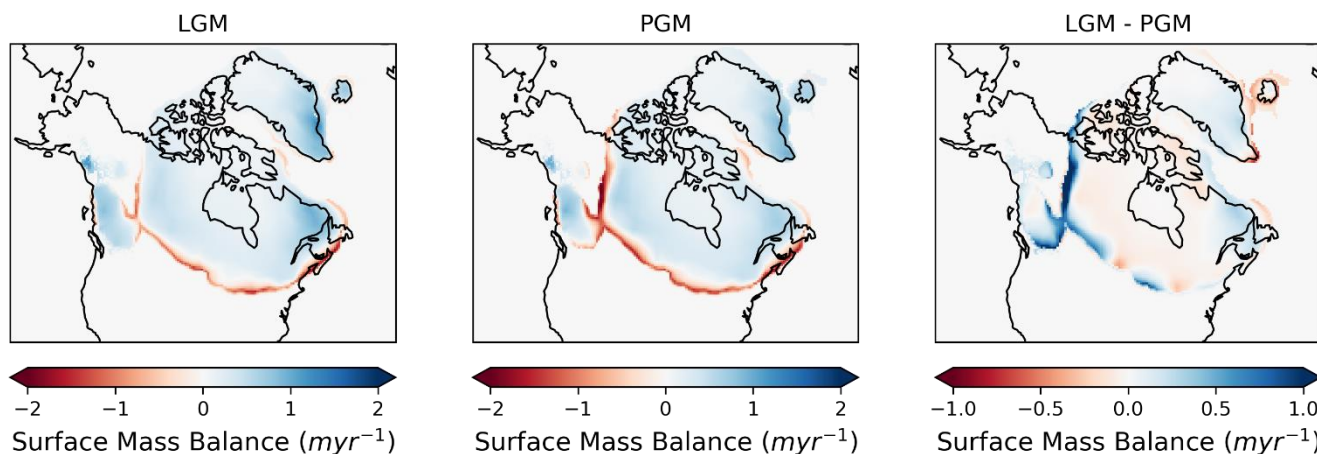
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The main cause of the difference in configurations between the LGM and PGM in this study is the less negative SMB at the LGM in the saddle region (Fig. 8). This is mostly a result of much lower ablation rates (runoff) in the summer months (JJA) at the LGM compared to the PGM, and to a lesser extent in spring and autumn (MAM and SON), and an increase in sublimation. The accumulation (snowfall) is similar between the two periods and does not contribute much to the SMB difference.





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**Figure 8. Mean surface mass balance of the constrained LGM and PGM ensembles averaged over model years 10 - 20 of the simulations and the difference between them.**

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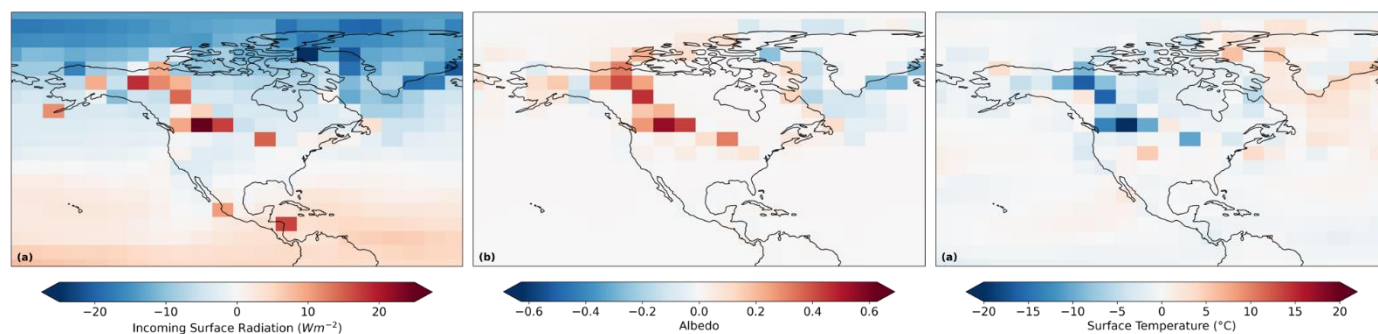
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This reduction in runoff occurs despite the LGM receiving more incoming top of the atmosphere shortwave radiation in early summer and more incoming surface radiation over North America at this time (Fig. 9a). Therefore, the positive SMB anomaly is a result of much more of this shortwave radiation being reflected back off the surface causing lower surface temperatures than at the PGM, allowing ice to build up and be maintained. In other words, the LGM has a higher albedo in this saddle region resulting in a more positive ice-albedo feedback (Figs. 9b and 9c). In contrast, the PGM has much lower albedo over the southern margin and saddle region, preventing ice growth in these regions.

During the analysis of these runs, we found that the coupling in the model was not passing all reductions in ice sheet area from Glimmer to FAMOUS in certain regions, particularly where entire FAMOUS gridboxes were initially covered in ice at all elevations (i.e. the saddle region at the LGM). This would have tended to reinforce the initial high albedo, positive mass balance surface conditions in the saddle region for the LGM configuration but would not have prevented the PGM simulations from growing ice in that area. To assess the role played by the initial conditions in our model simulations, we present an additional sensitivity analysis in the following section.

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302



303 **Figure 9. Mean difference between the NROY LGM and PGM simulations of mean summer (a) incoming surface shortwave**  
 304 **radiation; (b) albedo, and (c) surface temperature. All plots show the June-July-August average over model years 10-20 of the**  
 305 **simulations.**

### 306 3.4 Sensitivity analysis

307 To investigate the sensitivity of the final ice volumes to these differences in the initial conditions versus the differences in the  
 308 climate, a sensitivity analysis was carried out along with factorisation based on the method used in Lunt et al., (2012), also  
 309 used in Gregoire et al. (2015). We divided the differences in inputs between LGM and PGM into two factors; the initial ice  
 310 sheet configuration used in FAMOUS and the climate boundary conditions (orbital parameters, greenhouse gases and SSTs/sea  
 311 ice). Thus, the total difference in final ice volume ( $\Delta V$ ) between the LGM and the PGM can be written as Eq. (1):

$$312 \Delta V = dV_{ice} + dV_{climate} , \quad (1)$$

313 where  $dV_{ice}$  is the difference in final ice volume due to the different initial ice sheet configurations and  $dV_{climate}$  is the  
 314 difference due to the difference climate boundary conditions used.

315 The factorisation method requires  $2^N$  simulations (where N is the number of different components) to determine the  
 316 contribution of each component to ice volume difference, therefore  $2^2 = 4$  experiments are needed that systematically change  
 317 one variable. These experiments are listed in Table 4. We chose one of the NROY pairs of simulations (xpken and xpkyn  
 318 respectively) to carry out this analysis, these thus correspond to the full LGM and PGM simulations ( $E$  and  $E_{ci}$  respectively)  
 319 in the factorial decomposition. We further performed the two additional simulations needed for the decomposition (Table 4).  
 320 The relative contributions of the initial conditions and climate can be calculated by Eqs. (2) and (3):

$$321 dV_{ice} = \frac{1}{2}((V_i - V) + (V_{ci} - V_c)), \quad (2)$$

$$322 dV_{climate} = \frac{1}{2}((V_c - V) + (V_{ci} - V_i)), \quad (3)$$

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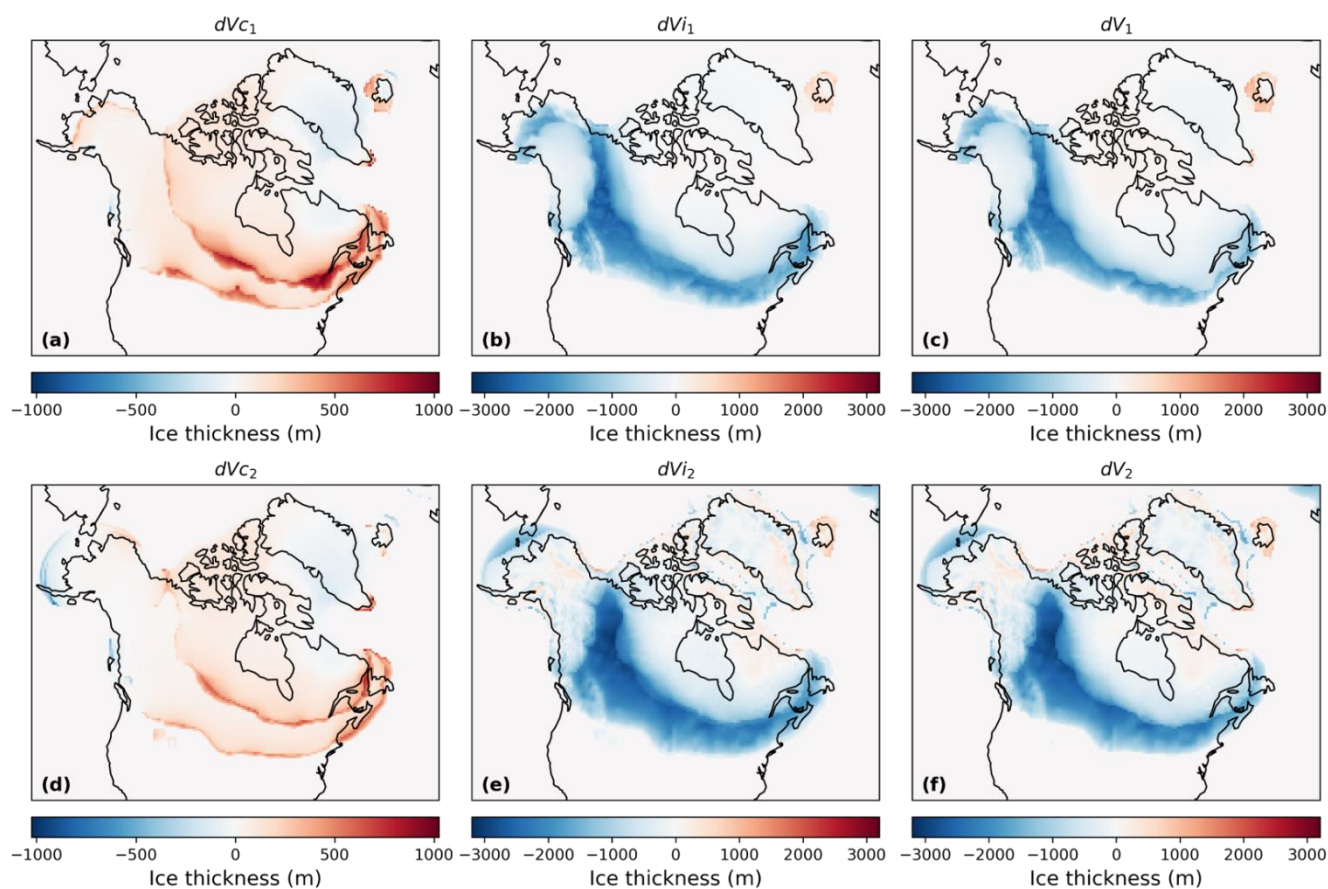
324 **Table 4. Ice sheet and climate conditions in each of the four experiments used in the sensitivity analysis**

<i>Experiment (final volume)</i>	<i>FAMOUS initial ice sheet</i>	<i>Climate (PMIP4)</i>
$E (V)$	LGM GLAC1D	LGM
$E_c (V_c)$	LGM GLAC1D	PGM
$E_i (V_i)$	PGM PMIP4	LGM
$E_{ci} (V_{ci})$	PGM PMIP4	PGM

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326 The results show that between the LGM and the PGM there was a total volume decrease ( $dV_1$ ) of  $8.89 \times 10^6 \text{ km}^3$ . The initial  
327 ice sheet configuration ( $dV_{i1}$ ) alone caused a decrease of  $1.12 \times 10^7 \text{ km}^3$  (125% contribution) but this was offset by the climatic  
328 conditions ( $dV_{c1}$ ) which resulted in an increase in volume of  $2.27 \times 10^6 \text{ km}^3$  (25% contribution) (Figs. 10a-c). To further  
329 understand the effect of initial conditions, we performed further simulations in which the initial conditions in the ice sheet  
330 component Glimmer was set to closely match the initial ice sheet extent and topography set in the climate component  
331 FAMOUS. Specifically, for the LGM, the Glimmer initial bedrock topography and ice surface elevation was prescribed from  
332 the GLAC-1D reconstruction used in the FAMOUS LGM boundary condition. For the PGM, the data needed for PMIP4  
333 reconstruction to be converted to the Glimmer initial condition were not available. Instead, both Glimmer and FAMOUS were  
334 initialised with the final timestep of the PGM experiment (xpkyn) since it closely resembles the PMIP4 reconstruction. This  
335 produced a similar result to the original factorial decomposition, with the initial ice conditions ( $dV_{i2}$ ) resulting in a 35%  
336 decrease in volume which was offset by the climate ( $dV_{c2}$ ) by a 4% increase (Figs. 10d-f).



337  
338 **Figure 10. Difference in final ice thickness between the PGM and LGM due to (a) climate parameters; (b) initial ice sheet conditions,**  
339 **and (c) the total difference. (d-f) show the same but for the corrected Glimmer ice sheets.**

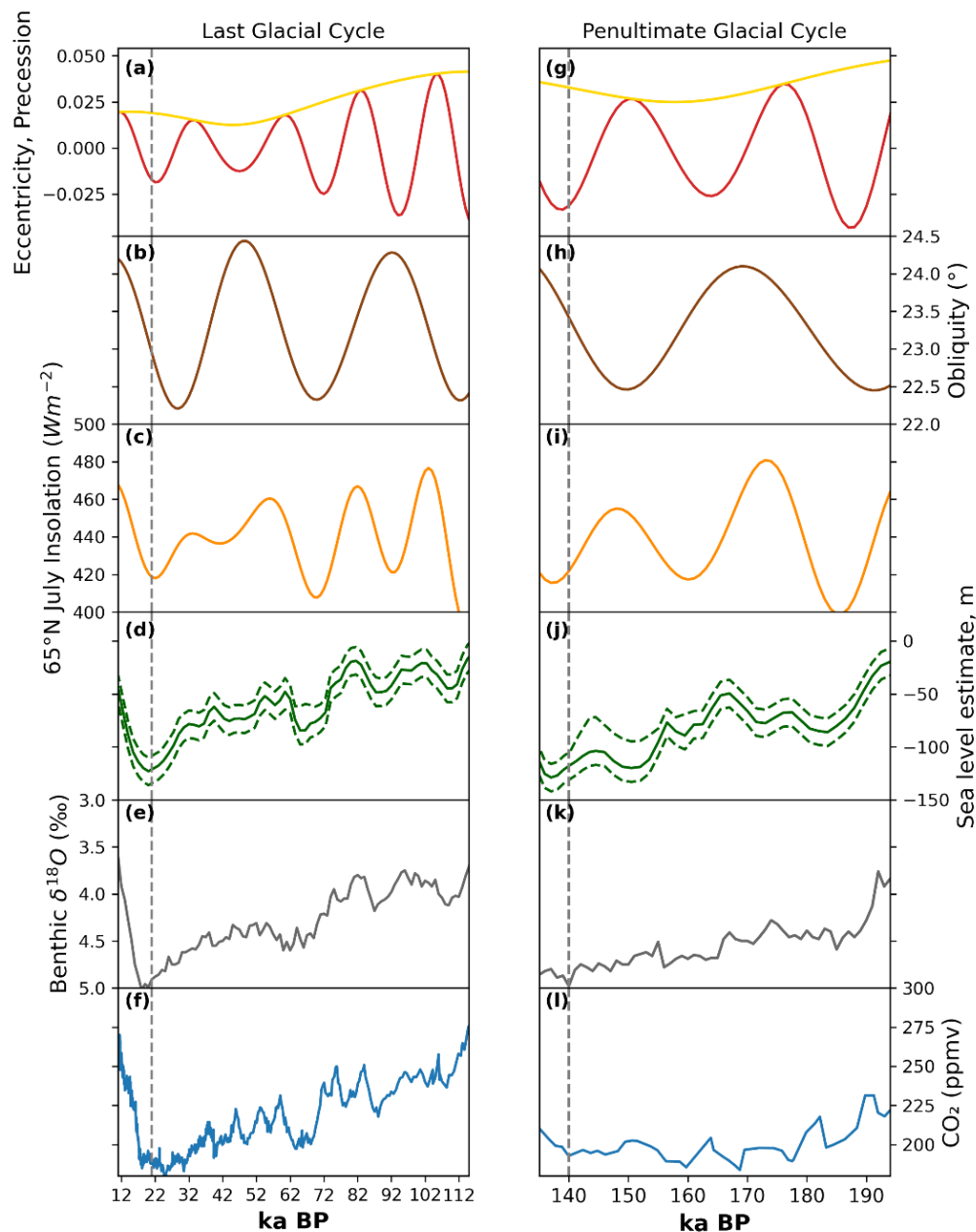


340 Based on these results, it is clear that the difference in initial ice cover, and resulting ice-albedo feedback, overwhelmingly  
341 determined the difference in final ice volume between the LGM and PGM in these simulations. Additionally, the vegetation  
342 fraction in the non-ice covered areas in the PGM initial conditions will have a compounding effect on this difference by  
343 introducing a vegetation-albedo feedback that limits the ice growth at the PGM. A similar conclusion was obtained by Abe-  
344 Ouchi et al., (2007) who studied the relative contribution to climate over ice sheets from the ice sheet itself and the orbital  
345 parameters and CO<sub>2</sub> concentration. They found the cooling caused by the ice sheet themselves was the dominant effect, mostly  
346 due to albedo feedbacks, which increase with ice sheet area. Kageyama et al., (2004) also highlighted in their study the  
347 importance of the albedo feedback on the maximum modelled North American ice volume. They show that changes in  
348 vegetation are needed to initiate glaciation over North America which is then accelerated by the ice-albedo feedback.  
349 Interestingly, the difference in orbital parameters, GHGs and SSTs (climate) between the LGM and PGM caused a larger North  
350 American ice sheet at the PGM in contrast to what evidence suggests. This is mostly a result of the orbital configuration which  
351 resulted in the Northern Hemisphere receiving less incoming solar radiation in spring and early summer (Table 1; Fig 2). This  
352 reduces the melting of snow that has accumulated in winter. The winter snow accumulation is also higher at the PGM than the  
353 LGM due to warmer air temperatures in autumn and winter, as a result of the orbital configuration, leading to a wetter climate.  
354 Summer SSTs are also cooler at the PGM due to lower spring insolation, further contributing to reduced runoff. In contrast,  
355 the Greenland ice sheet decreases in size due to PGM climate conditions likely due to higher sea ice concentration south of  
356 Greenland reducing the moisture source available for precipitation (Fig. 2b). This result highlights the importance of the  
357 evolution of these climate factors and the ice sheets during the preceding glacial cycles in determining the glacial maxima  
358 configurations.

359 For example, during the start of the Last Glacial Cycle (MIS 5; ~115-80 ka BP), the variation in 65° N summer insolation was  
360 relatively large as a result of changes in orbital parameters (Fig 11a-c), which resulted in multiple cycles of growth and  
361 recession of the North American Ice Sheets during this period, but total ice volume remained low (Bonelli et al., 2009;  
362 Ganopolski et al., 2010; Dalton et al., 2022). Insolation then reaches a minimum at ~70 ka BP (Fig 11c) which, combined with  
363 decreasing concentrations of CO<sub>2</sub> (~190 ppm at ~65 ka BP; Fig. 11f), lead to a significant increase in ice sheet volume to  
364 almost LGM extent (Fig. 11d) and a switch to more widespread glacial conditions at the MIS 5/MIS 4 transition (Bonelli et  
365 al., 2009; Dalton et al., 2022). The size of the NAIS at this time was large enough to induce positive feedbacks, such as the  
366 ice-albedo feedback, allowing its maintenance throughout MIS 4 and MIS 3 (~70-30 ka BP) despite an increase in insolation  
367 from ~50-30 ka BP (Fig. 11c). This was also supported by a continued decrease in CO<sub>2</sub> (Bonelli et al., 2009). Growth of the  
368 ice sheet could then continue to its glacial maximum extent following a further insolation and CO<sub>2</sub> decrease during MIS 2  
369 (~30-21 ka BP) (Fig. 11c-f). In contrast, prior to the PGM there were peaks in insolation at ~172 and ~148 ka BP that reached  
370 higher levels than were reached during MIS 4 and MIS 2, respectively, which were significant periods of growth at the LGM  
371 (Fig. 11i; Berger; 1978). This may have inhibited an initial significant build-up of ice over North America, as during MIS 4,  
372 preventing the initiation of an ice-albedo feedback strong enough to enable the continued growth towards a larger LGM



373 configuration and/or maintain its volume through the second insolation peak. In addition, there was more time between the  
374 LGM and the insolation maximum at ~50-30 ka BP compared to the PGM and the maximum at ~147 ka BP. Therefore, the  
375 PGM NAIS may have not had enough time to regrow before insolation started to increase again.  
376



377





378 **Figure 11. Evolution of climate proxies over the last two glacials: (a,g) precession index (red) with eccentricity as an envelope**  
379 **(yellow); (b,h) obliquity (Berger, 1978); (c,i) July insolation at 65° N (Berger and Loutre, 1999); (d,j) reconstruction of global mean**  
380 **sea level and uncertainty estimate (dotted lines) (Waelbroeck et al., 2002); (e,k) benthic  $\delta^{18}\text{O}$  global stack record (Lisiecki and**  
381 **Raymo, 2005), and (f,l) EPICA Dome C carbon dioxide ice core records (Luthi et al., 2008; Bereiter et al., 2015). The PGM and**  
382 **LGM are indicated by the dotted line.**

#### 383 4 Conclusions

384 We have performed and compared ensemble simulations of the LGM and the PGM using a coupled climate-ice sheet model  
385 (FAMOUS-ice) with an interactive North American Ice Sheet. The model was able to successfully simulate the ice sheet at  
386 both periods, compared to empirical evidence and other modelling studies, under different LGM and PGM climate boundary  
387 conditions and initial ice sheets. Overall, this study has shown that the underlying surface conditions, ice and snow cover and  
388 vegetation, used as boundary conditions in coupled climate-ice sheet simulations are extremely important in the resulting ice  
389 sheet volumes and extents because of the strong influence of the ice-albedo and vegetation-albedo feedbacks on the expansion  
390 of ice. In this study, the climate of each glacial maxima period has only a negligible effect on the simulated ice volume. Thus,  
391 investigation of the processes and interactions that took place prior to the glacial maxima will be needed to fully understand  
392 why the LGM and PGM NAIS configuration differed.

393 Additional feedbacks that played a role in the development of glacials into either an LGM-like or PGM-like mode are also  
394 missing in these simulations due to computational constraints. For example, the low resolution of the atmospheric component  
395 of FAMOUS means that it is capable of performing ensembles and long paleo runs while directly coupled to an ice sheet  
396 model, however it also means that many small-scale atmospheric processes (e.g. stationary wave response) caused by and  
397 affecting the ice sheet topography are not captured (Kageyama and Valdes, 2000; Liakka and Nilsson, 2010; Beghin et al.,  
398 2014; 2015; Liakka et al., 2012; 2016). Additionally, the shallow ice approximation used in Glimmer means that the ice sheet  
399 will not be able to simulate marine instabilities of advance and retreat (Pattyn et al., 2012). This effect will be minimal for the  
400 NAIS, but a more advanced ice sheet model would be required to simulate a marine ice sheet like the EIS.

401 As previously mentioned, the vegetation was kept fixed at present day distributions but the vegetation prior to and next to the  
402 ice cover has been shown to be very important for determining ice sheet expansion in models through the vegetation-albedo  
403 feedback (Kageyama et al., 2004; Colleoni et al., 2009b; Horton et al., 2010; Stone and Lunt, 2013). Therefore, implementing  
404 glacial maxima distributions or dynamical vegetation may affect the results since the reduction in forest and expansion of  
405 tundra/shrubs compared to present day would increase the albedo of the surface next to the ice and affect the climate (Meissner  
406 et al., 2003). Similarly, the fixed SSTs and sea ice concentrations used introduce uncertainty due to lack of constraint data and  
407 neglect any effects changes in ocean conditions and ice sheets have on each other (e.g. Timmerman et al., 2010; Colleoni et  
408 al., 2011; Ullman et al., 2014; Sherriff-Tadano et al., 2018; 2021). We recommend the use of a fully coupled atmosphere-  
409 ocean-vegetation-ice sheet model to further investigate these feedbacks. The effect of dust deposition and ice dammed lakes





410 have also been shown to have a large influence on the build-up of ice (e.g. Krinner et al., 2004; 2006; Naafs et al., 2012;  
411 Colleoni et al., 2009a) however further model developments would be needed to investigate these effects.  
412 Finally, the Eurasian ice sheet also displayed important differences between the LGM and PGM and had a large influence on  
413 the climate. It is likely that some of the differences in the configurations of the NAIS and EIS between the two glacial maxima  
414 resulted from their interactions with each other (Beghin et al., 2014; 2015; Liakka et al., 2016). To investigate the EIS at the  
415 PGM, we recommend the use of an efficient marine ice sheet model such as BISICLES that uses Adaptive Mesh Refinement  
416 to refine the processes occurring at marine margins that are more important for the marine based Eurasian ice sheet (Cornford  
417 et al., 2013; Gandy et al., 2019).

#### 418 **Appendix A: Eccentricity equation correction**

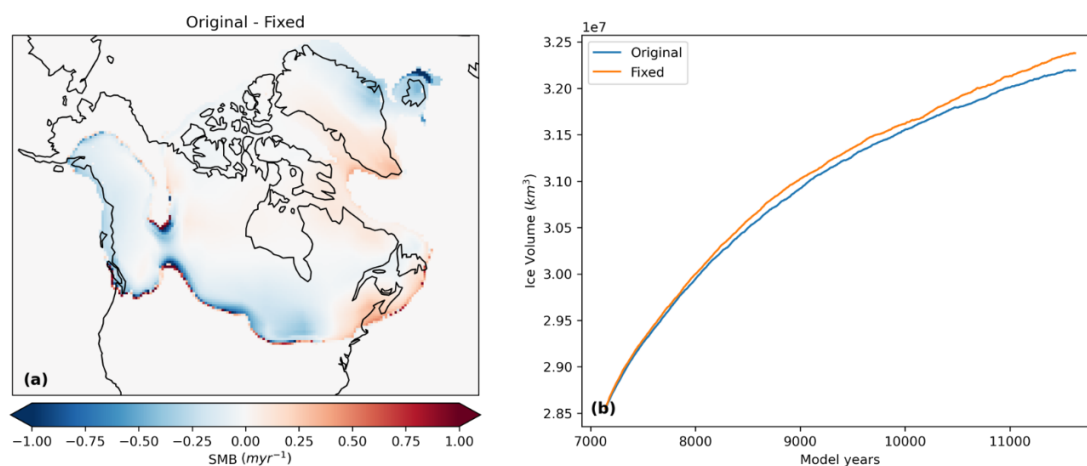
419 The equation for the role of eccentricity on solar insolation used in the simulations in this paper was:

$$420 S(t) = S_o \left(1 + \frac{e^2}{2}\right) (1 + e \cos v) / (1 - e^2)^2 \quad (4)$$

421  
422 However, this is incorrect and has now been corrected in the model to:

$$423 S(t) = S_o ((1 + e \cos v) / (1 - e^2))^2 \quad (5)$$

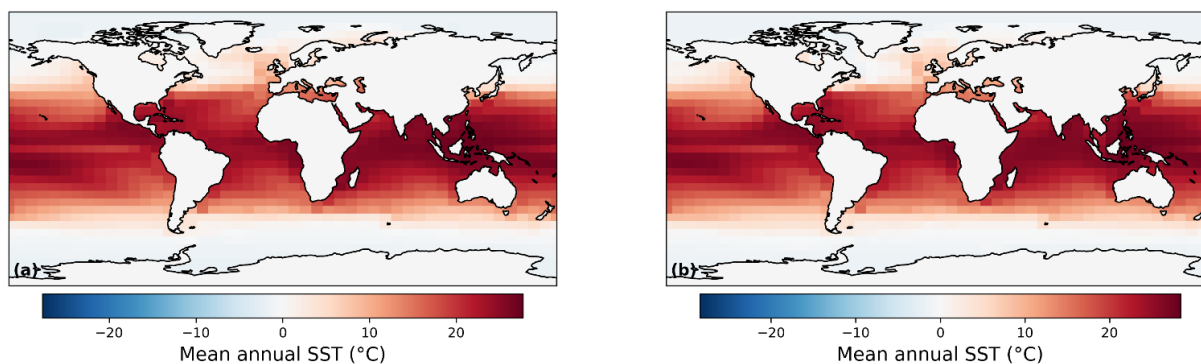
424  
425 The PGM experiment ‘xpky0’ was re-run with the correct equation and shows that on average the SMB was slightly more  
426 negative in our simulations than it should have been, leading to slightly smaller ice sheets (Fig. A1). However, the impact is  
427 small (and would be even smaller for the LGM given the lower eccentricity) and does not affect our overall conclusions.



428  
429 **Figure A1. (a) Difference between the SMB after 450 model years between the original simulation and the simulation using the**  
430 **corrected eccentricity equation and (b) the evolution of ice sheet volume for both experiments.**



## 431 Appendix B: Sea surface temperatures



432  
433 **Figure B1.** Mean annual SSTs used in this study from HadCM3 for (a) LGM and (b) PGM.

## 434 Appendix C: Wave 2 methodology

435 The ensemble design in this study was based on the ‘Not Ruled Out Yet’ (NROY) parameter combinations from a second  
436 wave of ensemble members that followed on from the 280 member ensemble performed in Gandy et al., (2023). From the first  
437 wave of simulations, only 18 out of these 280 members produced a large enough LGM North American Ice sheet to meet the  
438 volume and extent criteria they imposed (see details in reference). Further work was thus performed to augment the ensemble  
439 of simulations that met the NROY criteria. We used statistical emulation to identify plausible regions in the parameter space.  
440 As there was limited information to constrain the domain of plausibility in the parameter space, we instead implemented an  
441 early-stopping criteria that allowed us to prevent the full execution of model runs that were not expected to produce good ice  
442 sheets. To do this we first modelled, from Wave 1, the predicted equilibrium area of the ice sheet from the value of the initial  
443 surface mass balance. Mathematically, we specified;

$$444$$
$$445 A = f(b) + \epsilon, \tag{6}$$
$$446$$

447 where  $A$  is the ‘equilibrium’ ice sheet area after 10,000 ice sheet years,  $b$  is the 20 year averaged SMB value and  $f(\cdot)$   
448 may be any function. We considered  $f$  to be either linear or sampled from a Gaussian Process (GP) and found the linear model  
449 gave more conservative uncertainty estimates which was desired since the Wave 2 runs needed to bound the NROY space.  
450 The predictive interval for the model is  $P(b) = [f(b) + 3\sqrt{\text{var}(\epsilon)}, f(b) - 3\sqrt{\text{var}(\epsilon)}]$  and we targeted equilibrium ice sheet areas in  
451 the interval  $T = [1.5 \times 10^7 \text{ km}^2, 2 \times 10^7 \text{ km}^2]$ . The interval  $T$  is analogous to the target interval defined using Pukelsheim’s 3-  
452 sigma rule in standard history matching (Pukelsheim, 1994). Plausible values of  $b$  satisfy the condition that  $P(b) \cap T$  is non-  
453 zero, that is, for  $b$  to be plausible, the predictive bound  $P(b)$  and the plausible equilibrium ice sheet area  $T$  must intersect. It  
454 was found that the 20-year averaged SMB had to be at least positive to produce a plausible ice sheet.



455 To further improve efficiency, we used statistical emulation to produce plausible values of  $b$  (and hence equilibrium ice sheet  
456 areas); iterating the training data of the emulator with each wave of simulator runs. Define by  $\mathbf{x}$  the multivariate vector of  
457 parameters that they build the emulator over: here  $\mathbf{x}$  comprised of the 4 most influential parameters  $F_{snow}$ ,  $AV_{GR}$ ,  $Daice$ ,  
458 and  $Flow\ Factor$ . We model  $b$  with a random error process,  $b \sim GP(\mathbf{x}) + \eta$ , where the effects of the parameters not explicitly  
459 represented in  $\mathbf{x}$  are handled by the stochasticity of the process represented by  $\eta$ . Values of  $b$  were sampled using a stratified  
460 k-extended Latin Hypercube design (Williamson, 2015) and three sub-waves were executed, from which, a candidate set for  
461 the Wave 2 ensemble was extracted.

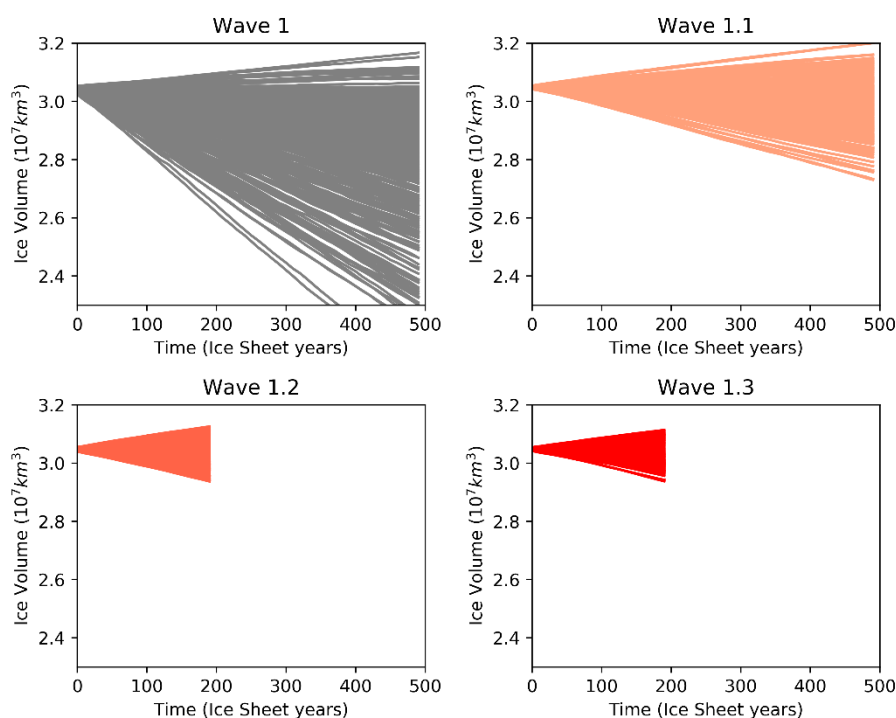
462 The first sub-wave (Wave 1.1) samples 200 ensemble members, which are predicted from the emulator to have non-negligible  
463 probability of positive SMB. This results in around 50% of simulations in this sub-wave having a positive SMB, an increase  
464 from 15% in the original wave (Fig. C1, Wave 1.1). We attempt to refine the predictive bounds on the GP model twice more  
465 (Fig. C1, Wave 1.2 and 1.3), with no improvement. This is likely due to the inherent stochasticity of the climate model and  
466 cumulative effects of the parameters that they absorb into the predictive error term. At the end of this process of iterative short  
467 waves, the candidate set contains over 1000 20-year long simulations that have a positive SMB over the North American ice  
468 sheet. From this candidate set, and again using stratified k-extended Latin Hypercubes, we select an optimal (with respect to  
469 space-filling and accounting for the previous Wave 1 runs) design of 200 ensemble members to continue for a full 10,000  
470 years to an equilibrium North American Ice Sheet. These 200 simulations make up the Wave 2. For context, this workflow of  
471 GP model sub-waves saved around 230,000 core hours (or about two months of real time) compared to running a full second  
472 ensemble wave.

473 Out of these 200 Wave 2 simulations, 176 members were identified to be NROY based on the original volume and extent  
474 thresholds. It is based on these results that we sub-sampled 62 parameter combinations for our simulations. This number of  
475 simulations was selected to enable us to run long equilibrium LGM and PGM simulations over a full ensemble within  
476 reasonable computational requirements. From the 176 NROY parameter combinations we randomly generated  $10^7$  candidate  
477 designs of size 62 from which we selected an approximate maximin design. This is obtained by: first linearly transforming  
478 each parameter onto the same range of  $[0, 1]$  to aid comparability; before computing the minimum distance between a  
479 parameter vector and its nearest neighbour; and then selecting the candidate design that maximised this distance. The resulting  
480 design possesses parameter vectors which are well-spaced and thus adequately cover the NROY space.

481 Our simulations use slightly different orbital parameter values and sea surface conditions to that of Gandy et al., (2023) (see  
482 Sect. 2). Thus, we do not expect the sample of 62 parameter combinations to provide full coverage of the NROY space but, as  
483 seen in section S2 of the supplementary information in Gandy et al., (2023), the output trends are sufficiently similar that we  
484 expect this to be close enough to an optimal sample. Whilst we may have also sampled some parameter combinations outside  
485 of the NROY space, we feel these will still provide valuable information about uncertainty in outputs at the LGM and PGM.  
486 Our detailed comparison to observations (see Sect. 3.1) identified six parameter combinations that match our criteria for LGM



487 and PGM ice extent and volume, thus demonstrating the success of this approach. Further exploration of the parameter space  
488 may produce NROY simulations in a different part of the parameter space but would not change the conclusion of this paper.  
489 Upon analysing the results, we found a technical error in the original Wave 2 ensemble which resulted in the values of the  
490 parameter *Daice* being shifted from its intended range of  $-0.4-0\text{ K}^{-1}$  to  $0-0.4\text{ K}^{-1}$ , this means that the albedo of the bare ice was  
491 increasing with melting, which is likely not the case. This produced larger values of surface albedo and thus larger ice sheets  
492 in these Wave 2 simulations (not shown here). In the ensemble of simulations presented here, we corrected the *Daice* values  
493 to match the intended parameter range. In some simulations, the switch of *Daice* value from a large positive number to a large  
494 negative number would have resulted in a decrease in surface albedo and resulting ice sheet volume. This effect is negligible  
495 for values of *Daice* closer to zero.  
496

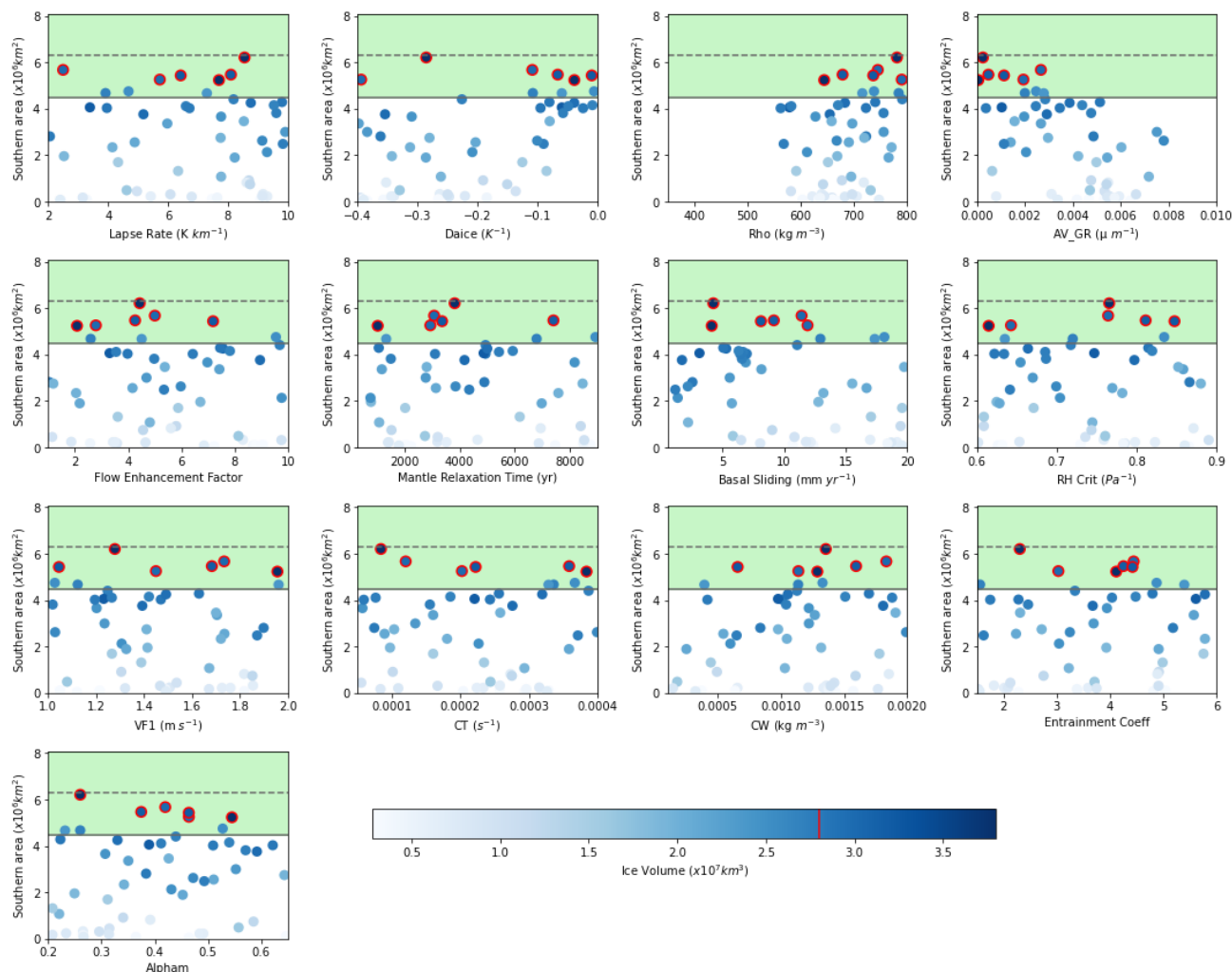


497  
498 **Figure C1. Ice volumes simulated in the successive ensemble sub-waves of simulations sampled to have a positive initial surface mass**  
499 **balance using the Gaussian Process emulator**



500 **Appendix D: Metrics vs parameters plots**

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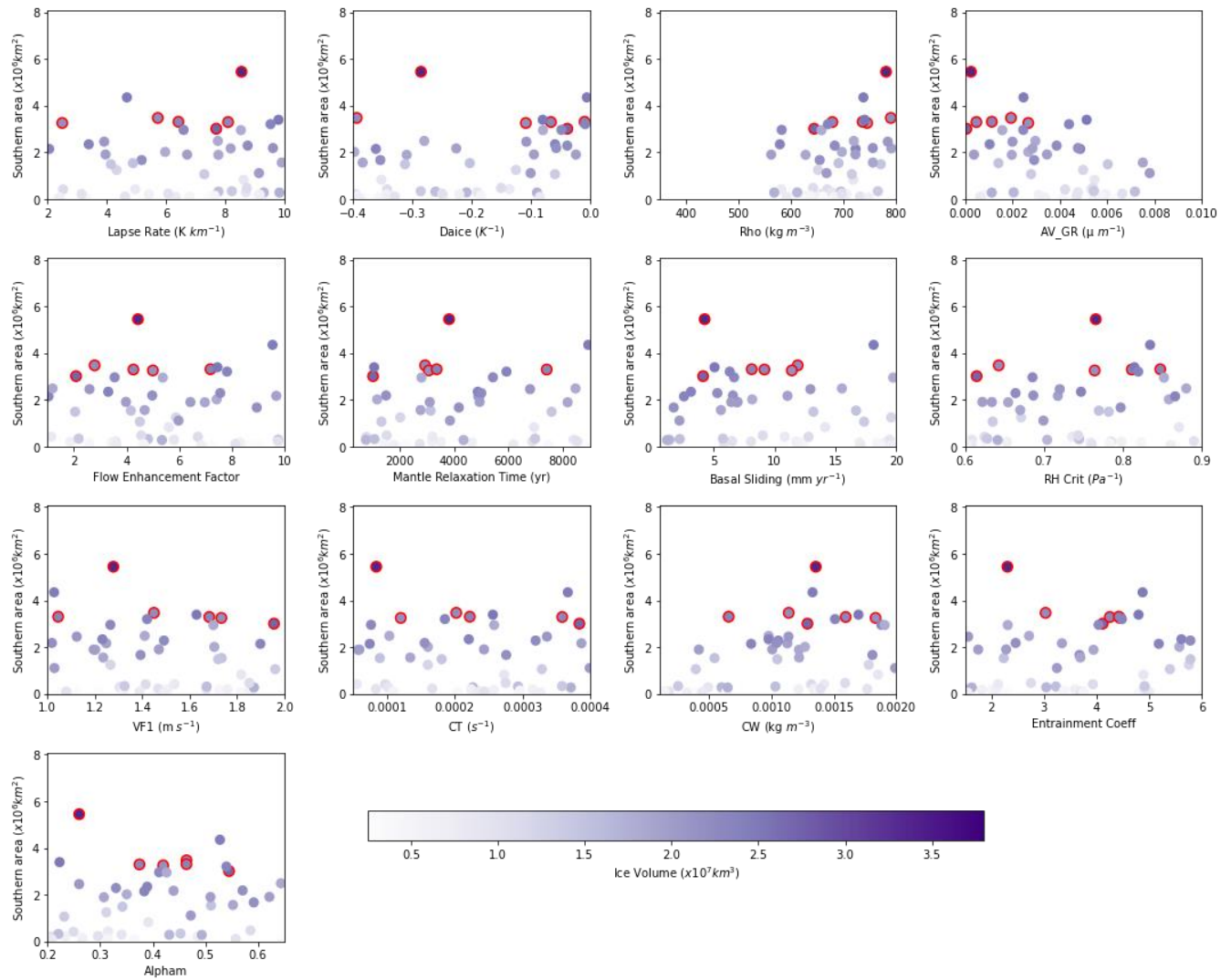
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**Figure D1. Southern area versus each of the 13 parameters varied for the LGM ensemble. The green shaded region shows the southern area constraint applied with the dotted line showing the exact area of the reconstruction and the solid line the solid line the minimum bound applied. The colour scale represents ice volume and the dots outlined in red are the six NROY LGM simulations with the red line on the colour bar showing the volume constraint.**



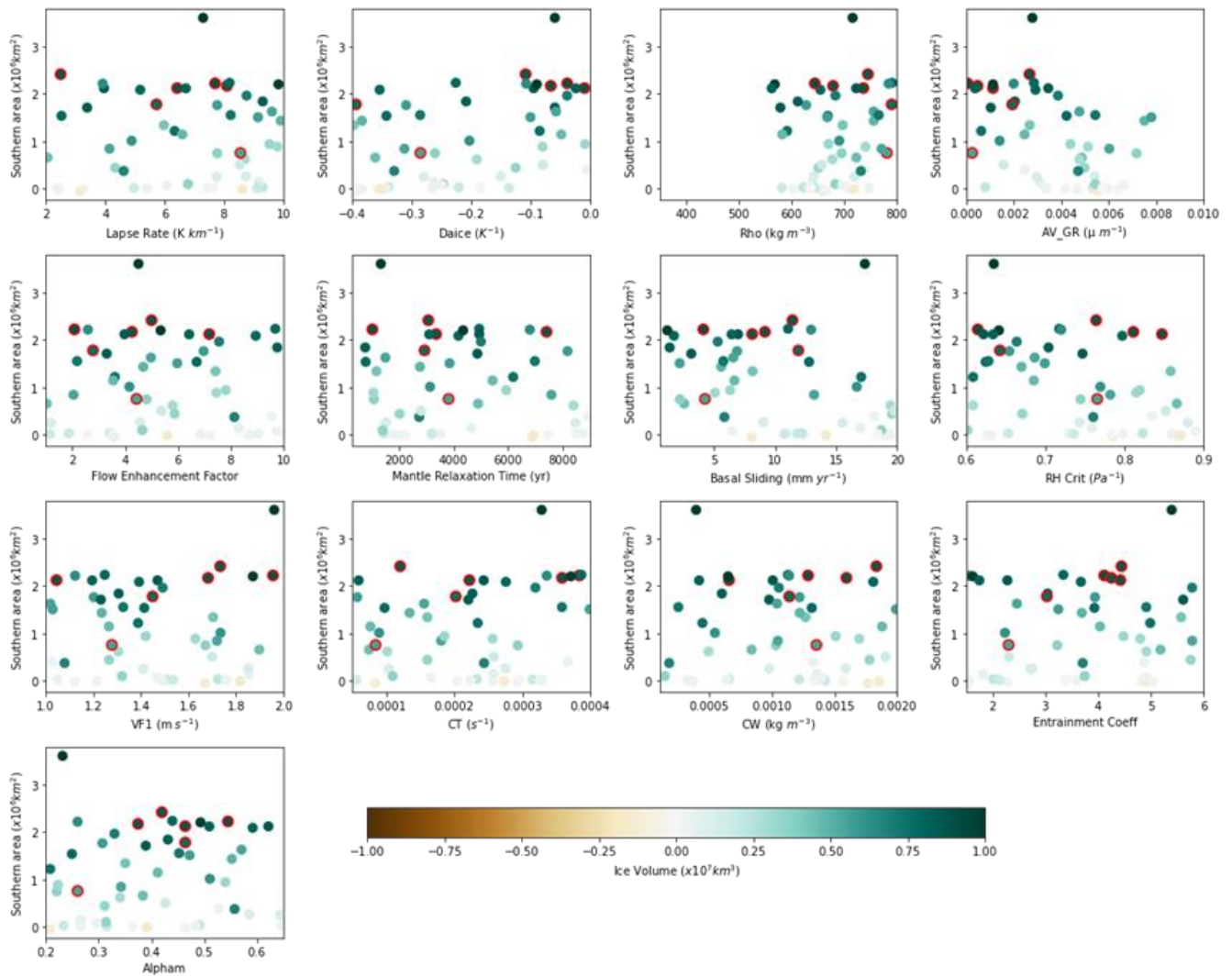
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**Figure D2.** Southern area versus each of the 13 parameters varied for the PGM ensemble. The colour scale represents ice volume and the dots outlined in red are the corresponding six NROY PGM simulations.





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**Figure D3. Difference in southern area versus each of the 13 parameters varied between the LGM and PGM ensemble members. The colour scale represents difference in ice volume and the dots outlined in red are the six NROY simulations.**



513 **Data availability**

514 For this pre-print, the boundary conditions used in this study as well as the full ensemble ice sheet model output and volume  
515 and extent metrics, climate timeseries for the NROY simulations and final ice volume data from the sensitivity tests have been  
516 made available to reviewers. All other model output data are available on request.

517 **Author contributions**

518 VLP lead the project and performed the majority of the work. VLP, LJJ, RFI and NG designed the simulations and VLP and  
519 NG prepared the initial and boundary conditions with support from OGP. VLP ran the simulations and analysed the results.  
520 LCA and NG designed and performed the Wave 2 simulations the ensembles were sampled from, and JO did the sampling.  
521 RSS provided technical support and updates for FAMOUS-Glimmer. VLP wrote the manuscript with comments and  
522 contributions from all co-authors. LJJ, RFI and NG supervised the project and LJJ acquired the funding.

523 **Competing interests**

524 The authors declare they have no conflict of interest.

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