

Rejuvenating the ocean: mean ocean radiocarbon, CO₂ release, and radiocarbon budget closure across the last deglaciation

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Abstract. Radiocarbon is a tracer that provides unique insights into the ocean's ability to sequester CO₂ from the atmosphere. While spatial patterns of radiocarbon in the ocean interior can indicate the vectors and timescales for carbon transport through the ocean, estimates of the global average ocean-atmosphere radiocarbon age offset (B-Atm) place constraints on the closure of the global carbon cycle. Here, we apply a Bayesian interpolation method to compiled B-Atm data to generate global interpolated fields and mean ocean B-Atm estimates for a suite of time-slices across the last deglaciation. The compiled data and interpolations confirm a stepwise and spatially heterogeneous 'rejuvenation' of the ocean, suggesting that carbon was released to the atmosphere through two swings of a 'ventilation seesaw' operating between the North Atlantic and the Southern Ocean/North Pacific. Sensitivity tests using the Bern3D model of intermediate complexity demonstrate that a portion of the reconstructed deglacial B-Atm changes may reflect 'phase-attenuation' biases that are unrelated to ocean ventilation, and that arise from independent atmospheric radiocarbon dynamics instead. A deglacial minimum in B-Atm offsets during the Bolling-Allerød could partly reflect such a bias. However, the sensitivity tests further demonstrate that when correcting for such biases, ocean 'ventilation' could still account for at least one third of deglacial atmospheric CO₂ rise. This contribution to CO₂ rise appears to have continued through the Younger Dryas, though much of the impact was likely achieved by the end of the Bolling-Allerød, indicating a key role for marine carbon cycle adjustment early in the deglacial process. Our global average B-Atm estimates place further new constraints on the long-standing mystery of global radiocarbon budget closure across the last deglaciation and suggest that glacial radiocarbon production levels are likely underestimated on average by existing reconstructions.

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Summary

45 Radiocarbon is best known as a dating tool, but it also allows us to track CO₂ exchange between the ocean and atmosphere. Using decades of data and novel mapping methods, we have charted the ocean's average radiocarbon 'age' since the last Ice Age. Combined with climate model simulations, these data quantify the ocean's role in atmospheric CO₂ rise since the last ice Age, while also revealing that Earth likely received far more cosmic radiation during the last Ice Age than hitherto believed.

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1 Introduction

50 The evolving spatial distribution of marine radiocarbon provides unique constraints on air-sea gas exchange at the sea surface, and the ocean dynamics that convey surface ocean-atmosphere pCO₂ differences into the ocean interior (Koeve et al., 2015). These transports of carbon in turn set constraints on the ability of the ocean to store CO₂, away from the atmosphere, in the form of either 'disequilibrium' or 'respired' dissolved inorganic carbon (Galbraith and Skinner, 2020). A slower overturning circulation enhances the accumulation of respired carbon in the ocean interior, in parallel with greater radiocarbon decay

55 (Menviel et al., 2017; Eggleston and Galbraith, 2018; Tschumi et al., 2011). Reduced air-sea exchange impedes the release of respired carbon from upwelled water to the atmosphere, while also impeding radiocarbon input from the atmosphere to the ocean (e.g. Khatiwala et al., 2019; Stocker and Wright, 1996). Both processes enhance carbon sequestration in the ocean (Galbraith and Skinner, 2020; Eggleston and Galbraith, 2018), while depleting the ocean's average radiocarbon activity relative to the atmosphere (Skinner and Bard, 2022). Therefore, estimates of the global average ocean-atmosphere radiocarbon age

60 offset (B-Atm), may provide powerful constraints on the evolving carbon exchange between the ocean and atmosphere (Siegenthaler et al., 1980). This in turn has implications for the closure of the radiocarbon budget of the atmosphere, given that the atmospheric radiocarbon inventory is largely set by the radiocarbon production rate, and ocean-atmosphere radiocarbon exchange (Siegenthaler, 1989).

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65 Few estimates of global average B-Atm prior to the instrumental record currently exist. This largely relates to the difficulty of estimating the 'volumetric representativity' of sparse data points, in order to weight their contribution to the global mean. Using a simple unweighted arithmetic mean of existing data from the Last Glacial Maximum (LGM, ~20 ka), Sarnthein et al. (2013) demonstrated an increase in the mean ocean-atmosphere radiocarbon age offset (B-Atm) by ~600 ¹⁴Cyr. A subsequent study instead made use of a Bayesian interpolation method, over a 4° resolution global grid, to generate a global field of B-Atm offsets, from which to derive a volume-weighted global average (Skinner et al., 2017). This alternative approach again showed that the glacial ocean was significantly more radiocarbon-depleted relative to the atmosphere, especially in the deep Atlantic and Southern Ocean (Sikes et al., 2016a; Skinner et al., 2021; Skinner et al., 2010; Gottschalk et al., 2020). This study inferred a mean ocean B-Atm increase of 689 ± 53 ¹⁴Cyr, as compared to the pre-industrial ocean. Using yet another approach, whereby individual data points were weighted according to their location (and their current offset from modern regional/basin

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85 averages), Rafter et al. (2022) have again confirmed an ‘aging’ of the deep ocean at the LGM. The study of Rafter et al. (2022) was restricted to the deep ocean > ~1000m water depth, and therefore did not yield global mean estimates. Nevertheless, the ~1050 ± 150 ¹⁴Cyr increase inferred by Rafter et al. (2022) for the deep ocean (approximately equivalent to >2km water depth), almost certainly implies a significant increase in the global mean B-Atm at the LGM. Note that Rafter et al. (2022) incorrectly identified the global mean estimate of Skinner et al. (2017) as referring to the deep ocean >3km water depth, and that consequently the estimates of Rafter et al. (2022) and Skinner et al. (2017) may be more consistent with each other than initially
90 apparent.

Overall, a consistent picture has therefore emerged of a relatively ‘aged’ dissolved inorganic carbon (DIC) pool in the LGM ocean. A recent review of the existing data has further demonstrated coherent patterns in marine radiocarbon evolution since the LGM, across the last deglaciation (Skinner and Bard, 2022). Several key observations have so far emerged from the compiled data: 1) at the onset of deglaciation, during Heinrich Stadial 1 (HS1, ~17.5-14.7ka), and to a lesser extent the Younger Dryas (YD, ~12.9ka-11.7ka), B-Atm offsets increased in the deep and intermediate North Atlantic, while they decreased in the Southern Ocean and North Pacific (particularly in the intermediate North Pacific, suggesting a ‘ventilation seesaw’ (e.g. Broecker, 1998; Skinner et al., 2013; Skinner et al., 2014; Menviel et al., 2018; Ahagon et al., 2003; Menviel et al., 2014; Freeman et al., 2015; Max et al., 2014; Okazaki et al., 2010; Walczak et al., 2020); 2) B-Atm offsets decreased significantly throughout the global ocean at the Bølling-Allerød (BA, ~14.7-12.9ka), in places suggesting an overshoot to values younger than pre-industrial (e.g. Barker et al., 2010); and 3) surface reservoir ages in the high latitudes of the North Atlantic and Southern Ocean tend to track B-Atm changes at intermediate depths (Skinner et al., 2019), suggesting limited changes in transport <2km water depth (and therefore a significant contribution from restricted gas-exchange efficiency), but with evidence for additional transport pathway and/or rate changes in the deepest ocean >2km (Skinner et al., 2021; Marchal and Zhao, 2021a). Encouragingly, similar observations have recently been reported based on a complex approach to data screening, averaging, and weighting (Rafter et al., 2022). The latter study emphasized the influence of transport changes on B-Atm offsets; enhanced mid-depth ventilation was inferred in the North Pacific at the LGM, and reduced transport rates were inferred in the deep Pacific at the LGM, and in the deep North Atlantic during HS1 and the YD.

110 Several questions arise in light of the existing body of deglacial marine radiocarbon data. First, what do the existing data imply for global average (i.e. rather than regional, or deep-ocean) B-Atm offsets across the last deglaciation? Secondly, can we determine the degree to which observed changes in the mean ocean B-Atm offset reflect changes in ‘ventilation’ specifically (i.e. circulation rates and gas exchange), as opposed to independent changes in atmospheric radiocarbon for example (Franke et al., 2008)? Thirdly, can we quantify the likely impact of past radiocarbon ventilation effects (regardless of their origin) on ocean-atmosphere carbon partitioning, and therefore atmospheric CO₂ (Skinner and Bard, 2022)? Finally, can the temporal evolution of mean ocean B-Atm offsets be reconciled with past atmospheric radiocarbon activities, and past radiocarbon

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production rates; and do they resolve the long-standing ‘mystery’ of deglacial radiocarbon budget closure (Broecker and Barker, 2007; Köhler et al., 2022; Kohler et al., 2006)?

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Here, we attempt to address each of these questions in turn. We update a previously deployed Bayesian interpolation technique (Skinner et al., 2017), which we apply to an updated compilation of radiocarbon data spanning the last deglaciation, to derive estimates of global average B-Atm offsets. These estimates are interpreted against a suite of new sensitivity tests and transient simulations, conducted using the Bern3D model of intermediate complexity (Müller et al., 2006; Ritz et al., 2011; Roth et al., 2014). The new simulations are aimed at constraining the potential magnitude of ‘ventilation’ *versus* ‘non-ventilation’ effects in global mean B-Atm offsets, as well as associated changes in ocean-atmosphere CO₂ partitioning.

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2 Methods

2.1 Radiocarbon data

Our data compilation is based on the work of Zhao et al. (2018), as presented in Skinner and Bard (2022). The compilation has been reconciled with a similar collection of data that has recently emerged (Rafter et al., 2022). A key difference in our compilation is that we include warm water (near surface) coral data, and surface reservoir age data where direct estimates exist (e.g. Skinner et al., 2019). Our compilation adopts revised chronologies and radiocarbon age offsets consistent with the latest *Intcal20* (Reimer et al., 2020) and *Marine20* (Heaton et al., 2020) calibration curves. To update the chronologies of the collected time-series, where the only age control derives from planktonic radiocarbon dates, we have used the *Marine20* calibration curve (Heaton et al., 2020), in conjunction with modern local deviations from the global average surface reservoir age (ΔR values \sim local R-age minus 550^{14}Cyr s) (Heaton et al., 2020), to generate ‘reservoir-age’ corrected calibrated calendar ages, using the R package *Bchron* (Parnell et al., 2008). While this approach has the advantage of remaining as faithful as possible to the original published age-scale for each record, it also has the notable drawback of neglecting potential changes in ΔR , which are known to have occurred in mid/high latitudes (e.g. Skinner et al., 2019), and that have been directly estimated in some of the studies included in the compilation. For records updated in this way, where the density and/or down-core variability of planktonic radiocarbon dates in a sediment core is sufficient to result in age-reversals (i.e. older dates stratigraphically above younger dates), we have generated new sediment depth-age models using the MCMC approach of *Bchron*, from which maximum likelihood calendar age estimates and 95% uncertainty intervals (~ 2 sigma) are obtained for each sample depth. For time-series that have made use of U-series dating (e.g. Robinson et al., 2005; Burke and Robinson, 2012; Chen et al., 2015; Hines et al., 2015), or stratigraphic alignments (including stratigraphic alignments that have provided time-varying surface reservoir age corrections, (e.g. Skinner et al., 2010; Gottschalk et al., 2020; Austin et al., 2011; Peck et al., 2006); as well as *a priori* assignment of down-core changes in reservoir ages (e.g. Ronge et al., 2020), we adopt the published calendar ages, as these do not depend on the atmospheric radiocarbon calibration curve.

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For all compiled data, we have recalculated radiocarbon reservoir age offsets relative to the *Intcal20* (Reimer et al., 2020) atmospheric reference curve using the *radcal* function in the R package *ResAge* (Soulet, 2015). This approach yields 95% uncertainty limits for the resulting radiocarbon age offsets based on the joint calendar age and radiocarbon age uncertainties. The resulting probabilistic B-Atm estimates differ from simple differences between marine and atmospheric radiocarbon values to some extent (Rafter et al., 2022). Estimates of the relative (radiocarbon) isotopic enrichment of the ocean *versus* the atmosphere are expressed as radiocarbon age offsets between the ocean and the contemporary atmosphere (i.e. in ^{14}C years), which are equivalent to $-8033 \times \ln(\alpha_{O-Atm})$, where α_{O-Atm} represents the ratio of marine radiocarbon activity *versus* the contemporary atmospheric radiocarbon activity at time T, i.e. Fm_O^T / Fm_{Atm}^T (Soulet et al., 2016). We do not refer to offsets between marine and atmospheric $\Delta^{14}\text{C}$ (i.e. $\Delta^{14}\text{C}_O - \Delta^{14}\text{C}_{Atm}$), since this metric does not relate to the isotopic depletion of the marine reservoir relative to the atmosphere in a predictable way without knowledge of the absolute atmospheric and marine radiocarbon activities (e.g. Cook and Keigwin, 2015).

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Data quality flags are assigned to anomalous values or datasets, including those that yield negative ventilation ages, negative reservoir ages, or negative benthic-planktonic offsets (B-P), or that exhibit significant down-core age-reversals (e.g. Rose et al., 2010). In addition, datasets that deviate significantly from regional trends, for example *due to* alternative age-models, are also flagged. The latter category includes records based on ‘plateau tuned’ (PT) chronologies (Sarnthein et al., 2015; Sarnthein et al., 2007; Sarnthein et al., 2020; Sarnthein et al., 2013; Ausín et al., 2021) that differ significantly from alternative reconstructions using the same data, or from reconstructions at proximal sites (Skinner and Bard, 2022). Such differences may relate to identifiable drawbacks of the ‘plateau tuning’ methodology (Bard and Heaton, 2021). Datasets that are interpreted as being influenced by hiatuses are also flagged (e.g. Ronge et al., 2020; Ronge et al., 2019), as are data that are interpreted as having been influenced by localized geological or sedimentary carbon sources (Rafter et al., 2019; Rafter et al., 2018; Lindsay et al., 2016; Lindsay et al., 2015; Ronge et al., 2016; Bova et al., 2018; Stott et al., 2009; Marchitto et al., 2007). A similar selection was effectively made in the recent study of Rafter et al. (2022), where all data from water depths less than ~1000m were omitted from consideration, and where individual data points were removed from *binned* regional water-depth groupings based on their deviation (by > 3 sigma) from the group’s weighted arithmetic mean.

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Although the data compiled in this study have all been published, including the vast majority in a previous compilation (Zhao et al., 2018), it remains possible that some data have been affected by as yet undetected biases arising from bioturbation (Bard et al., 1987; Dolman et al., 2021), and/or diagenesis (Wycech et al., 2016). Unless obviously *anomalous* signals are produced (Missiaen et al., 2020), such biases can be extremely difficult to positively identify, but could have profound impacts (Lougheed et al., 2020; Stott, 2023). *Although sites with very low sedimentation rates may be particularly prone to bioturbation effects (Bard et al., 1987; Dolman et al., 2021), the avoidance of all data from such locations might also introduce a sampling bias due their prevalence in the deepest ocean.* This presents a significant challenge. *In order to limit the potential impacts of anomalies caused by bioturbation (Dolman et al., 2021), we apply data flags to all sediment cores with average accumulation*

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210 rates < 2cm/kyr (rounded to the nearest digit). Of the remaining data, 75% derive from sites with accumulation rates > 10
cm/kyr and 95% derive from sites with accumulation rates > 4cm/kyr. A comparison of late Holocene data (< 6 ka BP) with
pre-industrial observations (Key et al., 2004) serves to demonstrate the general fidelity of B-Atm offsets derived from fossil
substrates, as compared to modern seawater values (see Figure 1, $R^2 = 0.81$). However, this comparison also illustrates the
significant scatter that exists in the fossil data (RMSE ~ 273 years; or an 'inverse prediction interval', which is more appropriate
if considering a 'calibration' of fossil data, ~ RMSE x 1.96 = 535 yrs (McClelland et al., 2021)). A similar result has recently
215 been demonstrated for sub-modern B-Atm offsets by Rafter et al. (2022). The observed scatter is expected to derive primarily
from sedimentary and sampling issues, rather than analytical uncertainty, and calls for caution when interpreting the subtler
features that emerge from the collected data. Accordingly, an underlying premise of our work (as for previous compilations)
is that the temporal trends and spatial patterns that emerge from a large body of independent data are far less likely to be
dominated by sedimentary or diagenetic biases that can more readily influence individual sediment cores (Ronge et al., 2019;
220 Stott, 2020).

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In order to assess the implications of applying the data flags described above, three alternative interpolations were performed:
1) the baseline interpolation, where all flagged data were excluded (Table 1, 'baseline'); 2) as for the baseline, but with low-
sedimentation rate sites (<2 cm/kyr) retained (Table 1, 'low sedimentation'); and 3) as for the baseline but retaining only high
225 sedimentation rates >10 cm/kyr (mainly affecting sites in the 2,500-4,000m water depth range). In keeping with a similar
assessment made by Rafter et al. (2022), we find that the inclusion of very low sedimentation rate sites (<2cm/kyr) has no
major effect on our findings, particularly for the global averages (Table 1). The same is true when only retaining sites with
very high sedimentation rates >10cm/kyr (Figures S2 and S3), though global averages are biased to slightly lower values
(Stott, 2023), particularly at the LGM (Table 1, 'high sedimentation'). This apparent lack of sensitivity arises here due to the
230 averaging of a relatively large number of data points (thereby reducing the influence of statistical outliers); however, it should
be borne in mind that the inclusion/exclusion of individual data points from sparsely sampled regions (e.g. for the Indian
Ocean) can influence the detail of inferred spatial patterns.

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study from the sparsely sampled deep Indian Ocean, where
sedimentation rates were slightly above the 2 cm/kyr cut-off (Table
1, 'Indian variant')

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have only a minor impact on the resulting global averages, though

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Rafter et al. (2022), we also find that the inclusion of low
sedimentation sites has no major effect on our findings, particular for
the global averages (Table 1).

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2.2 Time-slices

235 We isolate data associated with six key time-slices, based on their calendar ages and associated uncertainties: the LGM (19-
21.8 ka BP), Heinrich Stadial 1 (15-17.5 ka BP), the Bølling-Allerød (BA, 12.8-14.8 ka BP), the Younger Dryas (YD, 11.8-
12.7 ka BP), the early Holocene (EHOL, 9-11 ka BP), and the late Holocene (HOL, <6 ka BP). Note that the boundaries of
these time intervals are informed by, but do not correspond precisely to their respective ice-core chronozone definitions
(Rasmussen et al., 2014). This is due to the generally lower resolution of the marine records, and the need to be pragmatic in
240 avoiding as far as possible misattributing data to an adjacent chronozone. Where calendar age uncertainties result in an
ambiguous assignment of a single time-slice, no time-slice is assigned.

2.3 Radiocarbon ‘ventilation’ metrics

The term ‘ventilation’ is defined here as comprising the processes that influence physical and chemical property exchanges between the ocean interior and the atmosphere (Skinner and Bard, 2022). Accordingly, the term ‘(radiocarbon) ventilation age’ is not used to refer to ‘ideal ages’ or ‘transit times’ (England, 1995; Marchal and Zhao, 2021b), but rather to the degree of isotopic ($^{14}\text{C}/\text{C}$) disequilibrium between the ocean and atmosphere, as determined by gas exchange efficiency and water transport rates (Koeve et al., 2015). In this context ‘gas exchange efficiency’ is a catch-all term that refers primarily to the gas transfer piston velocity and includes the influence of the mixed layer equilibration time, which in turn depends on the mixed layer depth for example.

‘Radiocarbon ventilation’ must in turn be differentiated from measures of *relative isotopic enrichment*, such as B-Atm age offsets (Soulet et al., 2016). While ‘ventilation’ refers to a set of *processes* (principally gas exchange efficiency and water transport), observation-based metrics such as B-Atm reflect *measurements* of radiocarbon isotopic disequilibrium, typically between the ocean interior and the atmosphere or mixed layer (Skinner and Bard, 2022). This terminology mirrors a standard geological principle of separating descriptive names (e.g. diamicton) from the inference of process (e.g. glacial till).

In summary, B-Atm should not be used to directly imply transit times or ideal ages, which are not equivalent to ‘radiocarbon ventilation ages’. Furthermore, metrics such as B-Atm offsets are influenced by a wider variety of processes than just ventilation (i.e. more than just gas exchange and transport), including e.g. atmospheric radiocarbon production changes (Heaton et al., 2021). Therefore, B-Atm offsets should only be used to refer to ventilation (i.e. gas-exchange and transport) with caution, as we explore in this study.

2.4 Interpolation

We use the interpolation method described in Skinner et al. (2017), updated to weight input data according to their reported uncertainties. The method interpolates the observed ocean-atmosphere radiocarbon age offset anomalies onto the grid of an ocean model with 24 vertical levels and $2^\circ \times 2^\circ$ horizontal resolution. The anomalies are defined relative to the modern GLODAP radiocarbon ages (Key et al., 2004). The interpolating function is constructed using a weighted superposition of radial basis functions centred at the sediment-core locations. These basis functions are built here using an exponential basis function, $\varphi(r) = e^{-\varepsilon r}$, with shape parameter, ε . The distance r is defined as the time for tracer impulses to spread out from the core locations to the rest of the grid, based on the modern transport and density stratification. The basis-function weights, the shape parameter, and two hyper-parameters that scale the relative size of the precisions in the prior and likelihood functions, are inferred using a three-level Bayesian procedure (further details are given in Skinner et al. (2017)). The interpolation process is repeated independently for each time-slice so that the data from one time-slice do not influence the interpolation for the other time-slices.

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This interpolation method is somewhat conservative because it is based on the distribution of relative tracer transport/diffusion timescales in the modern ocean. This choice is premised on the need to obtain an interpolated solution in three dimensions that is physically sensible and that is informed by large-scale oceanographic requirements, including vertical stratification, basin margins/topographic boundaries, and dominant transport pathways (e.g. the Antarctic Circumpolar Current, Deep Western Boundary currents, etc.). It is important to note that the interpolation method does not impose the modern transport on the resulting interpolated fields, nor does it assume strict adherence to the modern density field. Rather, it represents a physically guided and data-constrained anomaly relative to the modern state. The inferred anomalies relax to zero in the absence of data constraints, though some locations in the ocean can have far reaching influence. This method represents a counterpoint to the approach recently taken by Rafter et al. (2022) for example, where spatial interpolations were illustrated for data projected onto two-dimensional sections. In the latter study, averages were also produced for time-series grouped and weighted according to their position within the modern density- and B-Atm distributions (i.e. locations closest to the modern mean for a given basin and density class were weighted to dominate the mean in the past). With this approach, the spatial distribution of density classes is assumed invariant over time, despite the inference of circulation changes. This mirrors to some extent our use of a modern transport field as a guide to the volumetric representativity of individual sample locations. Interpolating a sparsely sampled 3D field is a difficult problem, and a diversity of approaches is surely useful; however, it should be noted that Rafter et al. (2022) did not consider the upper ~1,000m of the ocean and did not produce or discuss global averages.

In order to assess the accuracy of the interpolation method of Skinner et al. (2017), and its ability to represent different circulation states than the modern, we apply it to modelled radiocarbon fields produced in this study using the Bern3D model (see below), and by Menviel et al. (2017) and Menviel et al. (2018) using the LOVECLIM model. We extract data from the modelled fields at the same locations as the available proxy data (e.g. for the LGM), and attempt to reproduce the modelled global field using the interpolation method. As shown in **Figure 2**, the interpolation method is generally successful in reproducing the simulated radiocarbon fields for altered circulation states. The interpolation does best for circulation states characterised by altered wind, diffusivity, and gas exchange, yielding average errors for 1000 randomly selected individual grid cell estimates of RMSE ~ 276 years (INT_ALL60, this study), RMSE ~ 244 years (INT_ALL80, this study), RMSE ~ 398 years (V3LNAw, (Menviel et al., 2017)), and RMSE ~ 511 years (V3LNAwSowSHWw, (Menviel et al., 2017)) (see **Figure 2** and **Figure 3**). The interpolation performs slightly less well for extremely different circulation states, e.g. yielding RMSE ~ 457 years for the HS1 simulation of Menviel et al. (2018), and RMSE ~ 577 years for the collapsed AMOC simulation of Menviel et al. (2017), both of which yielded reversed radiocarbon gradients in the Atlantic as compared to modern (**Figure 3**). Encouragingly, the interpolation method reproduces the modelled fields more accurately than model simulations of the LGM are typically able to match observations (Muglia et al., 2018; Menviel et al., 2017).

In the present study, the interpolation is primarily intended for calculating global average B-Atm offsets, using a relatively sparse dataset (N = 124 to 476). It is therefore of particular importance that the interpolation is found to reproduce modelled

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global averages even more accurately than the spatial distributions, with RMSE ~ 53 years (**Figure 4**, circles). In contrast, simple arithmetic means of B-Atm values drawn from a sub-set of locations in the modelled fields consistently underestimate the true global mean, by ~274 years on average (**Figure 4**, triangles). This comparison underlines the importance of applying a spatially resolved volumetric weighting to the observations for the derivation of accurate global averages. While our interpolation approach leaves room for improvement, it complements alternative data-constrained modelling approaches that have been applied to the LGM and the YD for example (Pöppelmeier et al., 2023), and it represents a first step towards addressing the problem of deriving 3D global fields, and appropriately weighted global averages, from sparsely sampled proxy data, without (as yet) the use of forward or inverse models.

2.5 Modelling

The intermediate complexity Bern3D v2.0 Earth System model (Ritz et al., 2011; Roth et al., 2014) has been used to explore a series of idealised scenarios that aim to probe the parallel sensitivities of atmospheric CO₂ and marine radiocarbon, subject to changes in global vertical diffusivity, Southern Ocean winds, and/or Southern Ocean gas-exchange efficiency. The implementation of radiocarbon in the Bern3D model is described elsewhere (Müller et al., 2006; Muller et al., 2008; Dinauer et al., 2020), and the simulations presented here extend those performed by Jeltsch-Thömmes et al. (2019) and Dinauer et al. (2020), as well as those performed using an earlier version of the Bern3D model by Tschumi et al. (2011). The applied version of the Bern3D model comprises a single-layer energy-moisture balance atmosphere with a thermodynamic sea-ice component (Ritz et al., 2011), coupled to a 3D geostrophic-frictional balance ocean (Edwards et al., 1998; Müller et al., 2006), and a 4-box representation of the land-biosphere that simulates the dilution of an atmospheric isotopic perturbation by the land biosphere, but here does not address changes in land carbon stocks (Siegenthaler and Oeschger, 1987). Furthermore, marine sediments were not implemented in the model set-up used here, as we seek to isolate the impacts of ventilation processes alone (i.e. ocean gas exchange and transport).

A first set of sensitivity tests (PI-), performed under 'pre-industrial' conditions, and run out for 10000 years with a fully interactive climate-carbon cycle system (minus sediments), consisted of step changes in: 1) global vertical diffusivity in the ocean (Kv); 2) Southern Ocean wind stress (SW); 3) Southern Ocean gas-exchange efficiency, or piston velocity (SG); or 4) a combination of Kv, SW and SG (ALL). In each case the relevant parameter was reduced by 20%, 40%, 60% or 80% relative to its baseline value, and impacts were evaluated after 2000 years (impacts after 500 and 5000 years are illustrated in **Figure S1**).

In an additional set of simulations, we employed a range of glacial/deglacial scenarios, similar to Jeltsch-Thömmes et al. (2019), where each is again defined by idealised adjustments to vertical diffusivity, wind stress, Southern Ocean gas exchange or a combination of these, as described for the above PI- sensitivity tests. For these simulations, the model was 'spun up' into equilibrium over 35,000 model years under pre-industrial (1700 CE) boundary conditions. The model was then ramped linearly

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over 5,000 years from the resulting PI equilibrium state into ‘glacial’ boundary conditions (i.e. representative of 50 ka BP) for ice sheet albedo, greenhouse gas radiative forcing, and insolation, as well as prescribed values for diffusivity, wind stress, and Southern Ocean gas-exchange efficiency as described above. After this, ice-sheet albedo, greenhouse gas radiative forcing, and insolation were varied based on observations since 50 ka BP, with diffusivity, Southern Ocean wind stress, and Southern Ocean gas-exchange efficiency remaining constant for 32,000 years and then relaxing linearly back to PI control values from 18 ka BP to 14.6 ka BP, after which PI control values were maintained.

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In one set of ‘glacial/deglacial’ model simulations (INT-), atmospheric radiocarbon was prescribed from 50 ka BP based on the *Intcal20* reference curve (Reimer et al., 2020). Radiocarbon concentrations in the model were therefore required to be consistent with the observed atmospheric radiocarbon concentration, via changes in the global radiocarbon inventory (these are equivalent to *ad hoc* changes in radiocarbon production). In a second set of simulations (FIX-), atmospheric radiocarbon activity was held constant at 140 permil, while in a third set of simulations (CONST-) atmospheric radiocarbon production rates were held constant at the PI value. In each case a control simulation was performed with evolving boundary conditions as described above. A series of additional simulations were then performed with altered global vertical diffusivity (Kv), Southern Ocean wind stress (SW), Southern Ocean gas-exchange efficiency (SG), or a combination of these (see Table S1). The goal of these transient simulations is not to reproduce the last deglaciation, but to assess the sensitivity of both marine radiocarbon and atmospheric CO₂ to a variety of changes in ocean ventilation (in terms of both the type of forcing, and its magnitude), specifically where concurrent radiocarbon production rate changes are either minimised (CONST-/FIX-) or maximised (INT-).

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3 Results

Figure 5 illustrates zonally averaged radiocarbon age offsets (B-Atm) for the LGM, HS1, BA, and YD, in the Atlantic and Pacific basins, based on global interpolations of the compiled data (using the ‘baseline’ data flags). Differences between each successive time-slice interpolation are shown in Figure 6. Zonal averages and time-slice offsets for the Indian Ocean are shown in Figure 7. Zonal averages for the EHOL and HOL, which differ only slightly from the modern field (Key et al., 2004), are shown in Figure 8, compared with the interpolated fields for the BA and LGM. As noted elsewhere (Skinner and Bard, 2022; Rafter et al., 2022), the current paucity of data from the Indian Ocean makes interpolations for this basin somewhat tentative, and more strongly dependent on individual data points (Figure 7). This highlights the Indian basin as an important target for future work reconstructing past B-Atm changes. Across all the time-slice interpolations, correlations between observed and interpolated values for the same locations yield R² values in the range 0.59-0.86. Because the interpolation is not performed on a 2D zonal plane, local B-Atm estimates may deviate from the zonal average where zonal gradients exist. Furthermore, the interpolations necessarily provide an average ‘snapshot’ for an entire time-slice, and therefore will mask variability within the time-slice. These aspects appear to be particularly relevant for HS1 in the Atlantic, as discussed below.

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Deleted: Figure 7 illustrates the impact of including a few data points from the LGM and HS1 (Bharti et al., 2022), where sedimentation rates were just above the 2cm/kyr cut-off threshold. While the intensity of the LGM and HS1 anomalies in the Indian Ocean are significantly affected by these data, the associated global average B-Atm offsets remain similar (Table 1), albeit with the most significant impact on the reconstructed anomaly between HS1 and the LGM. ...

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The global interpolations illustrated in **Figure 5** capture the main features of deglacial B-Atm evolution that have previously been identified in individual time-series, and in compiled time-series that were grouped by basin and region/water depth (Skinner and Bard, 2022). Broadly similar patterns have also been identified in [averaged time-series and 2D zonal projection contour plots of compiled data](#) (Rafter et al., 2022). The main features include:

- 1) increased B-Atm offsets throughout the global ocean at the LGM, as compared to the [HOL/modern state](#) (**Figure 8a, b and Figure 8g, h**), in particular >2000m water depth, as demonstrated previously (Skinner et al., 2017), but with a slightly larger anomaly ~ 760 ^{14}C years;
- 2) a step-wise ‘rejuvenation’ of the ocean interior across the last deglaciation;
- 3) evidence for positive B-Atm anomalies in the deep North Atlantic, from the LGM to HS1, and again from the BA to the YD, occurring in parallel with changes of broadly opposite sign in the Southern Ocean and the intermediate depth North Pacific (boxed areas in **Figure 6g, h** and **Figure 6c, d**);
- 4) a marked ‘rebound’ to lower B-Atm offsets throughout the ocean, from HS1 to the BA (**Figure 6e, f**); and
- 5) a further rebound to lower B-Atm offsets in the deep North Atlantic, from the YD to the EHOL, again with evidence for changes of opposite sign in the Southern Ocean and intermediate depth North Pacific (boxed areas in **Figure 6a, b**).

Perhaps the most striking aspect of the successive time-slice reconstructions is the marked drop in B-Atm offsets that coincides with the transition from HS1 to the BA (**Figure 6e, f**), which involved positively correlated changes in B-Atm in the deep Southern Ocean and deep North Atlantic (Skinner et al., 2013), and which has been linked to an ‘overshoot’ in B-Atm offsets at some locations (Barker et al., 2010; Hines et al., 2015). Notably, [as shown in Figure 8](#), the step change at the BA resulted in a global B-Atm distribution very similar to the [late Holocene and modern](#), despite representing the approximate temporal mid-point of deglaciation.

The features identified above can also be discerned in regional time-series averages (Skinner and Bard, 2022; Rafter et al., 2022). [We show this here](#), using cubic splines for North Atlantic, Southern Ocean and North Pacific data grouped according to the boxed areas highlighted in **Figure 6**. [The resulting regional average splines are illustrated in Figure 9b, c, d, compared](#) with global average B-Atm values derived for each successive time-slice (**Table 1**, filled circles in **Figure 9d**). During HS1 and the YD, the regional splines exhibit broadly anti-phased trends in B-Atm between the deep North Atlantic and the deep Southern Ocean and intermediate North Pacific (shaded vertical bars, **Figure 9**). The collected time-series thus support the antiphase patterns apparent in [the zonal averages shown in Figure 6](#) (boxed areas). [The averaged time-series only tentatively support the suggestion of a ‘flipped’ vertical radiocarbon gradient between the intermediate- and deep North Pacific at the LGM](#) (Rafter et al., 2022), and more emphatically demonstrate this for the deglacial period that includes the YD, BA and HS1. [Indeed, during the BA, the North Pacific, North Atlantic and Southern Ocean all exhibit B-Atm offsets at least as low as](#)

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480 modern, resulting in global average B-Atm that also approaches the modern, and slightly ‘overshoots’ relative to the subsequent
YD (see **Figure 9d**). Consistent with the fact that the **Pacific accounts for ~50% of the global ocean volume and the Southern**
Ocean ventilates ~58% of the ocean interior (Primeau, 2005), **Figure 9** also shows that global average B-Atm estimates
generally track the **deep North Pacific and the deep Southern Ocean**. Different patterns of variability are expressed in the North
Atlantic, and the intermediate North Pacific, where more rapid and regionally important fluctuations are apparent (Freeman et
485 al., 2015).

4 Discussion

In principle, evolving large-scale patterns of ocean-atmosphere radiocarbon age offsets (B-Atm) will primarily reflect the
combined influences of 1) radiocarbon production, 2) ocean transports, 3) air-sea gas exchange efficiency (especially in the
490 regions of deep-water export), and 4) the changing contributions of different source regions to locations in the ocean interior.
While the influence of atmospheric radiocarbon production changes on evolving B-Atm offsets is often ignored, it may in
principle influence deep ocean B-Atm offsets through relatively rapid (i.e. sub-millennial) changes in atmospheric radiocarbon
activity that are only conveyed to the deep ocean on the millennial time-scale of ocean turn over (Adkins and Boyle, 1997).
This can produce a convergence or divergence of marine and atmospheric radiocarbon ages, due to atmospheric radiocarbon
495 changing quickly and the deep ocean remaining relatively invariant (Franke et al., 2008; Heaton et al., 2021).

Below, we discuss how all these processes have influenced marine radiocarbon cycling across the last deglaciation.
Accordingly, we address: 1) evidence for the operation of a ‘ventilation seesaw’ that we emphasize was linked to *both* gas-
exchange and transport anomalies emanating from the main regions of deep- and intermediate water formation (i.e. the North
500 Atlantic, Southern Ocean, and North Pacific); and 2) the potential for atmospheric radiocarbon dynamics (independent of
ocean ventilation) to bias B-Atm offsets by ~ hundreds of ¹⁴C years, in particular during the apparent BA ‘overshoot’. We
further seek to quantify the likely carbon cycle impacts associated with the observed global average B-Atm changes, and to
reconcile these with the evolution of the global radiocarbon budget since the last glacial period.

505 4.1 Ventilation ‘seesaws’: gas-exchange and transport effects

A survey of the modern ocean’s transport pathways indicates that ~86% of the ocean interior is sourced by water that last made
contact with the atmosphere in three key regions: the Southern Ocean (contributing ~58%), the high latitude North Atlantic
(contributing ~21%), and the North Pacific (contributing ~7%) (Primeau, 2005). These three regions account for ~31% of the
ocean’s surface, and ‘ventilate’ ~86% of the ocean’s interior (Primeau, 2005). Although the long equilibration time for
510 dissolved $\Delta^{14}\text{C}(\text{DIC})$ means that a water parcel’s point of last contact with the atmosphere will not necessarily be equivalent
to its point of ‘radiocarbon equilibration’ (high latitude sources are typically characterised by distinct levels of disequilibrium
(Bard, 1988; Matsumoto, 2007)), changes in the three main regions of deep-water export will exert a strong influence on
temporal variations in the ocean interior’s radiocarbon distribution.

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This expectation is clearly borne out in the global interpolations shown in **Figure 5** and **6**, and in the regional stacks illustrated in **Figure 9**. Deglacial changes in marine radiocarbon distribution were apparently dominated by anomalies extending from the North Atlantic (affecting the deep Atlantic in particular), coordinated with anomalies of broadly opposite sign originating in the Southern Ocean and North Pacific (**Figure 6** and **Figure 9**). The time-slice interpolations and regional averages therefore cohere with numerous previous proposals for ‘ventilation seesaws’ operating between the North Atlantic and the Southern Ocean (Broecker, 1998; Skinner et al., 2013; Skinner et al., 2014; Menviel et al., 2018), and between the North Atlantic and North Pacific (Menviel et al., 2014; Freeman et al., 2015; Max et al., 2014; Okazaki et al., 2010; Walczak et al., 2020).

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As suggested previously (Skinner et al., 2019; Skinner and Bard, 2022), **Figure 9** also shows that B-Atm offsets in the deep (>2km) North Atlantic (**Figure 9b**, blue line and shaded area) exhibit a similar pattern of variability to the upper ocean (0.2-2km) and surface ‘reservoir ages’ in the region (**Figure 9b** grey lines with shaded area, and dashed red lines). The same is apparent in the Southern Ocean (**Figure 9d**). These relationships, and their further link to polar climate variability (**Figure 9a**, e), suggest a mechanistic link between the observed deglacial B-Atm variability and the ‘thermal bipolar seesaw’ (Stocker and Johnsen, 2003; Epica Community Members, 2006). This association most likely operated via coordinated changes in North Atlantic and Southern Ocean convection and advection (Broecker, 1998; Menviel et al., 2015; Skinner et al., 2014; Skinner et al., 2020), associated with regional changes in sea-ice (Skinner et al., 2019; Rae et al., 2018), winds (Sikes et al., 2016b; Menviel et al., 2018), and/or buoyancy forcing (Ferrari et al., 2014; Hines et al., 2019; Watson et al., 2015). The further coupling between the Southern Ocean and North Pacific has been proposed to relate to freshwater balance in the Pacific basin (Menviel et al., 2014), possibly influenced by changing moisture transports across the Isthmus of Panama (Leduc et al., 2007), and/or Cordilleran ice mass balance (Walczak et al., 2020).

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While the ‘ventilation seesaws’ noted above may reflect changes in ocean circulation to some extent, it is important to note that they also reflect changes in gas-exchange efficiency. This is demonstrated by the fact that the amplitude of B-Atm variability in the shallow Atlantic (0.2-2 km water depth (e.g. Freeman et al., 2015; Chen et al., 2015) differs very little from that of surface ‘reservoir ages’ in the Northeast Atlantic (Skinner et al., 2019) (**Figure 9b**, grey line and dashed red line, respectively). This implies a muted contribution from flow speed changes in the upper Atlantic < 2km (Bradtmiller et al., 2014), and a dominant influence on radiocarbon signatures from gas-exchange (i.e. ‘pre-formed ages’) instead. Again, the same is true for the Southern Ocean, where B-Atm offsets from the shallow ocean (0.2-2km) generally overlap with surface reservoir age reconstructions (**Figure 9d**, grey line and dashed red line, respectively). Accordingly, B-Atm values at the LGM in mid-depths may indeed have been consistent with the modern transport (Marchal and Zhao, 2021a).

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In contrast however, ‘pre-formed ages’ cannot account for the amplitude of B-Atm changes observed in the deep Atlantic (**Figure 9b**, blue line), and the deep Southern Ocean (**Figure 9d**, blue line), indicating a more significant contribution from

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570 water sourcing and/or transport changes >2 km, and perhaps between 2 and 3km water depth in particular (Skinner et al., 2021;
Lund et al., 2015; Rafter et al., 2022). The patchy response seen in the zonal average anomaly for the North Atlantic between
HS1 and the LGM (**Figure 6h**) contrasts with the clearer signal seen in individual and collected time-series (**Figure 9b**). In
part, this likely reflects a spatially heterogeneous hydrographic response, both in the depth domain (Skinner et al., 2021; Lund
et al., 2015), and in the eastern *versus* western North Atlantic (Gherardi et al., 2005; Ng et al., 2018). This interpretation is
supported by the less ambiguous positive B-Atm anomaly seen in the deep North Atlantic regional time-series spline during
575 HS1 (**Figure 9b**, blue line and shaded area; see also Rafter et al. (2022)).

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The patchy spatial interpolation outcome for HS1 may also reflect a more complex pattern of ventilation during HS1 than is
captured by a ‘collapsed AMOC’ scenario that lasted until the onset of the BA, as indicated by a recent multi-proxy and
modelling study (Pöppelmeier et al., 2023). Furthermore, in addition to tentative evidence for a ‘mid-HS1 B-Atm minimum’
580 (Skinner and Bard, 2022), there is also evidence for a drop in B-Atm offsets at intermediate depths (< 2.5km) in the western
Atlantic late in HS1 (Robinson et al., 2005; Thiagarajan et al., 2014), and evidence for declining B-Atm offsets in the Nordic
Seas ~400 years prior to the onset of the BA (Muschiello et al., 2019). Any such changes during HS1 will have been averaged
out in our time-slice interpolation for 15-17.8 ka BP. These observations underline the need for further detailed reconstruction
of the spatial expression and temporal evolution of ocean-atmosphere radiocarbon offsets across the North Atlantic during
585 HS1, particularly with a view to disentangling transport and gas-exchange impacts.

Overall, the message that emerges from the collected data is that deglacial marine B-Atm changes throughout the global ocean
were significantly influenced by *both* air-sea gas exchange effects and mass transport effects (Koeve et al., 2015), with the
latter primarily affecting parts of the deep ocean >2km, but with evidence for a complex response during HS1. The inferred
590 spatial expression and coordination of both gas-exchange and transport effects in the Atlantic may have been related to changes
in the spatial extent of different water masses in the deep ocean. This would be consistent with previous work where B-Atm
offsets have been found to place stronger constraints on water mass ‘geometry’ than e.g. Atlantic Meridional Overturning
Circulation (AMOC) strength *per se* (Muglia and Schmittner, 2021; Pöppelmeier et al., 2023). The influence of further non-
ventilation effects at each time-slice is taken up in the following section.

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4.2 ‘Attenuation biases’ in B-Atm offsets at the BA

As noted above, global average B-Atm changes can occur independently of ventilation effects, due to rapid atmospheric
radiocarbon variability that the deep ocean is too slow to respond to. We refer to such effects as ‘attenuation biases’, as they
reflect the phase and attenuation response of a slowly adjusting reservoir (the global ocean interior), subject to continuous
600 exchange with a more rapidly changing reservoir (the atmosphere) (Maier-Reimer and Hasselmann, 1987). The primary
external (i.e. non-marine) driver for rapid atmospheric radiocarbon variability is likely to be changes in radiocarbon production

605 (Köhler et al., 2022), though transient terrestrial carbon sources (e.g. from permafrost) might also be hypothesised (Köhler et al., 2014; Wu et al., 2022).

For attenuation biases in mean ocean B-Atm offsets to occur, atmospheric radiocarbon activity would have to change rapidly, on a timescale that is shorter than the mean mixing time of the ocean (i.e. < 1000 yrs). Accordingly, gradual long-term trends in radiocarbon production are unlikely to produce significant attenuation biases, as is confirmed by model simulations where atmospheric radiocarbon production is based on relatively smooth trends in mean relative (geomagnetic) palaeointensity (RPI) (Dinauer et al., 2020). However, given the scatter amongst existing reconstructions of past radiocarbon production (e.g. Laj et al., 2004; Adolphi et al., 2018; Nowaczyk et al., 2013; Channell et al., 2018), particularly on (sub-) millennial time scales, it remains unclear to what extent rapid (centennial/millennial) radiocarbon production variability, and/or other non-marine carbon sources, might have biased B-Atm offsets via their impact on the atmosphere. Therefore, to explore the maximum possible contribution of externally driven atmospheric radiocarbon variability to mean ocean B-Atm changes, we compare transient simulations using the Bern3D model where radiocarbon production rates or atmospheric radiocarbon are held constant (i.e. FIX- and CONST-, respectively) with simulations where nearly all variability in atmospheric radiocarbon is assumed to have occurred independently of ocean ventilation change (i.e. INT-, where atmospheric radiocarbon is prescribed according to *Intcal20*). The difference between the INT- simulations and their CONST/FIX- counterparts yields the *maximum* amplitude of B-Atm changes that could be produced independently of ocean ventilation change, for each idealised scenario. Indeed, as discussed below, these estimates are likely over-estimates, as they are premised on the assumption that all atmospheric radiocarbon variability occurred independently of ocean-atmosphere radiocarbon exchange, which is unlikely.

625 As illustrated in **Figure J0** (and summarised in **Table S2**), our idealised scenarios demonstrate three key points regarding the emergence of ‘attenuation biases’ in B-Atm offsets:

- 1) these biases depend on the occurrence of rapid atmospheric variability that is not forced by ocean ventilation change, but they can in principle result in persistent long term biases *via* an accumulation of centennial/millennial perturbations (e.g. as for a ‘random walk’ process);
- 630 2) such biases are time-varying, and may be positive relative to a given reference time (e.g. at the Laschamps event, **Figure J0a**), negative (e.g. notably, at the BA onset, **Figure J0c**), or nil (e.g. at the LGM, **Figure J0b**), thus enhancing, diminishing or not affecting the B-Atm changes that are expressed between time periods (**Figure J0d**); and
- 635 3) the relative contribution of such biases will be diminished and potentially eliminated to the extent that they coincide with large abrupt ocean ventilation changes (**not included in our idealised model simulations**), though **it is apparent that** their magnitude depends mainly on the amplitude of atmospheric radiocarbon production changes, **as they yield** a relatively invariant offset between the INT- and CONST/FIX- simulations for any given time-slice, **Figure J0a-d**).

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It is important to stress that the true magnitude of such attenuation biases cannot be determined without prior detailed knowledge of the history and magnitude of both ocean ventilation changes and non-marine carbon/radiocarbon inputs to the atmosphere. Nevertheless, two observations are worth underlining. The first is that maximum estimates of the attenuation

650 biases that could hypothetically affect deglacial B-Atm offsets would only result in relatively minor biases in the incremental B-Atm changes between time-slices (ranging from -227 to +25 ¹⁴C-years). Even if maximal attenuation biases are hypothesised, the observed deglacial mean ocean B-Atm trends must include a significant ventilation contribution. Indeed, such changes are directly attested to by proxy evidence for ocean transport and sea-ice change (McManus et al., 2004; Schüpbach et al., 2018), as well as the spatially heterogeneous patterns in marine B-Atm offsets that we present here (e.g.

655 **Figures 6 and 9).**

A second key observation is that correcting for such biases would tend to diminish the apparent ‘ventilation surge’ that might be inferred from the change in B-Atm between HS and the BA (and more so between the LGM and the BA). Thus, if the rapid atmospheric radiocarbon decline across HS1 and at the onset of the BA was entirely driven by radiocarbon production changes

660 and/or non-marine carbon inputs to the atmosphere, then the change in global average B-Atm between the LGM and the BA would be biased by at most ~ 227 ¹⁴C yrs. In this case, radiocarbon ventilation ages during the BA would be ~ 227 ¹⁴C yrs *older* than inferred from observed B-Atm offsets, resulting in a smaller radiocarbon ventilation change between the LGM and the BA. Non-ventilation biases affecting global average B-Atm differences between the LGM and the BA could therefore imply a smaller contribution from gas-exchange and transport rate changes to the apparent radiocarbon ‘ventilation surge’

665 suggested **for the BA in Figure 6e, f**. In this case, the ‘rejuvenation’ of the marine radiocarbon pool may not in fact have been completed mid-way through deglaciation as initially apparent (Rafter et al., 2022). This inference would be more consistent with stable carbon isotope (¹³C/¹²C) evidence (Sikes et al., 2016b), which suggests an ongoing contribution to ocean-atmosphere carbon exchange well beyond the BA. Indeed, a significant portion of the convergence between marine and atmospheric radiocarbon observed at the BA may have been driven by ‘old’ carbon release to the atmosphere, e.g. from melting

670 permafrost (Köhler et al., 2014; Wu et al., 2022).

Although they remain difficult to accurately quantify, ‘attenuation biases’ unrelated to ventilation are important to acknowledge, as these may exert a subtle yet potentially significant influence on our interpretation of transient changes in B-Atm offsets and their quantification, as illustrated above for the BA. These effects will need to be explored in greater detail in

675 the future; a task that will inevitably require the use of numerical models, and that will also require accurate knowledge of past radiocarbon production changes (Köhler et al., 2022; Dinauer et al., 2020). The accuracy of our existing radiocarbon production records is discussed further in section 4.4.

4.3 Ventilation-related atmospheric CO₂ change

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685 In theory, a broadly linear relationship between atmospheric CO₂ change and global average marine radiocarbon age anomalies
is to be expected, when these are driven by gas-exchange and/or transport rates (Skinner and Bard, 2022; Skinner et al., 2017).
This is because: 1) longer residence times in the ocean interior result in greater respired carbon accumulation along with greater
radiocarbon decay; and 2) restricted gas exchange in regions of ‘upwelling’ and/or deep mixing impede the conversion of
690 respired and/or disequilibrium carbon to equilibrium carbon (Eggleston and Galbraith, 2018), while also impeding the invasion
of radiocarbon into the ocean.

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Our model sensitivity tests, involving shifts in vertical diffusivity, Southern Ocean winds, and/or gas exchange, cohere with a
number of existing simulations using box-models and Earth System models of intermediate complexity (e.g. Tschumi et al.,
2011; Kwon et al., 2011; Skinner and Bard, 2022), confirming a broad relationship between B-Atm anomalies and associated
695 atmospheric CO₂ change (Table 2 and Figure 11a, symbols), despite large differences in model experiment set-up. These
experiments include the effects of pCO₂ changes on air-sea ¹⁴C exchange (Galbraith et al., 2015; Bard, 1988). Our sensitivity
tests indicate consistent sensitivities for individual suites of experiments, e.g. for varying Southern Ocean winds, vertical
diffusivity, or Southern Ocean gas exchange rates. However, depending on the processes responsible for altering deep ocean
ventilation, modelled sensitivities span a range of approximately ±50% for a given magnitude of B-Atm change. Furthermore,
700 the sensitivities may vary with perturbation timescale, with different impacts on millennial timescales than on glacial-
interglacial or longer timescales (Jeltsch-Thömmes and Joos, 2020). For example, a reduction in Southern Ocean winds draws
down atmospheric CO₂ and increases global average B-Atm on a timescale of ~2,000 years, via impacts on gas-exchange rates
and upper ocean turnover (Figure 11a). However, after ~5000 years (see Figure S1) the mean ocean B-Atm response is
diminished, primarily due to the spin up of North Atlantic deep water, which replaces southern sourced water in the deep
705 Atlantic. A muted B-Atm response can therefore be seen for stronger Southern Ocean wind forcing in the 50ka simulations
illustrated in Figure 10a (purple symbols). Despite their potential role in transient millennial scale pulses, Southern Ocean
winds are therefore unlikely to account for long-term glacial-interglacial B-Atm changes, at least in the absence of further
forcing to suppress North Atlantic ventilation (Meniel et al., 2018).

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710 The broadly linear scaling apparent in the suite of PI model sensitivity tests is consistent with basic theory (Skinner et al.,
2017; Skinner and Bard, 2022), based on a two-box ocean connected to an atmosphere to form a closed system (Figure 11a,
solid line). This simple inventory theory predicts a higher sensitivity for higher global export productivity and/or a higher
‘Revelle buffer factor’ (i.e. higher background atmospheric pCO₂), all else being equal. Clearly, the marine carbon cycle
response to the variety of ventilation processes that can affect radiocarbon cannot be reduced to a single linear scaling.
715 However, the degree of consistency in the parallel sensitivities of atmospheric CO₂ and marine B-Atm offsets implies that the
wide range of modelled sensitivities may be approximated by a theoretical prediction using an arbitrary ± 50 % range in export
productivity as a tuning parameter in the 2-box ocean model (Figure 11a, broken lines). This theoretical scaling, arbitrarily
tuned to the range of more complex model outputs, would suggest -6.3 ± 3.2 ppm CO₂ change per 100 ¹⁴Cyrs of global mean

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730 B-Atm change. Such a sensitivity would tentatively imply a drawdown of atmospheric CO₂ by $\sim 50 \pm 27$ ppm associated with an increase in global average B-Atm by $\sim 789 \pm 32$ ¹⁴C yrs, as reconstructed for our 'baseline' data flag scenario at the LGM and assuming a maximum 'attenuation bias correction' (Table 1). Estimates based on uncorrected B-Atm offsets, and/or on alternative data flagging scenarios (Table 1), differ by up to ~ 8 ppm. These estimates should be interpreted as indicating a non-negligible contribution to deglacial CO₂ rise, perhaps equivalent to over a third of the total glacial-interglacial atmospheric CO₂ change.

735 Interestingly, observed global average B-Atm estimates also loosely correlate with observed atmospheric CO₂ changes across the last deglaciation (open symbols, Figure 11b). The correlation with observed pCO₂ anomalies, particularly for the BA, is improved when global average B-Atm is 'corrected' for maximal possible attenuation biases as discussed above (filled symbols in Figure 11b). The similarity of the observed correlation and the modelled CO₂-sensitivity (red lines, Figure 11b) again merely suggests that a significant portion of the incremental changes in atmospheric CO₂, stepping through the deglaciation from the LGM to the early Holocene, could in principle be accounted for by ocean ventilation changes that influenced global average B-Atm offsets. Although the steeper dashed line in Figure 11b would imply a maximal ~ 80 ppm contribution to deglacial atmospheric CO₂ rise due to ocean ventilation alone, we believe this is unlikely. Rather, the observed correlation between atmospheric CO₂ and global average B-Atm offsets more likely implies a mixture of direct and indirect causal connections that may have been coordinated by the thermal bipolar seesaw. Direct impacts of ocean ventilation on atmospheric CO₂ (e.g. Figure 11a) would therefore have coincided with contributions from linked processes, such as ocean temperature change, export productivity anomalies, etc. (Marchal et al., 1998; Menviel et al., 2008; Jochum et al., 2022; Gottschalk et al., 2019; Menviel et al., 2012), or indeed long-term Southern Ocean winds changes that have had little effect on mean ocean B-Atm.

750 In any event, if mean ocean B-Atm changes scale in a consistent manner with marine CO₂ sequestration changes (as suggested by simple inventory theory and the relationships in Figure 11a), then the ocean ventilation contribution to deglacial atmospheric CO₂ rise would have been primarily associated with HS1, the BA and the YD. The magnitude of ocean ventilation impacts on CO₂ after ~ 15 ka BP appear to hinge crucially on attenuation biases that may have affected B-Atm offsets during the BA. Nevertheless, the increases in atmospheric CO₂ observed across the Holocene coincide with rather muted changes in global average B-Atm, and low potential for significant attenuation biases. This suggests a minor role for ocean ventilation in CO₂ rise from the onset of the Holocene, which would be consistent with the interference that CO₂ rise from ~ 6 ka BP was primarily linked to changes in ocean temperature, the terrestrial biosphere, coral reef formation, and/or the solid Earth (i.e. volcanism, ocean alkalinity) (Joos et al., 2004; Broecker and Clark, 2007). This observation might also resonate with the speculative proposal of a 'natural tendency' for atmospheric decline across an interglacial due to ocean ventilation processes (Barker et al., 2019).

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790 4.4 Towards a closure of the global carbon and radiocarbon cycles since the LGM

The reconciliation of past radiocarbon production changes with records of atmospheric pCO₂ and $\Delta^{14}\text{C}$ (hereafter $\Delta^{14}\text{C}_{\text{atm}}$) across the last deglaciation represents a long-standing puzzle that remains unresolved (Bard, 1998). This puzzle has a direct bearing on our understanding of past atmospheric pCO₂ change, as well as our understanding of geomagnetic and solar variability (Heaton et al., 2021). The record of atmospheric radiocarbon variability indicates a significant decrease across the last deglaciation, equivalent to a change in $\Delta^{14}\text{C}_{\text{atm}}$ of just over ~400 permil (Reimer et al., 2020) (Figure J2b, blue line). If there was no change in the steady state global radiocarbon inventory, the 90 ppm increase in atmospheric pCO₂ that occurred since the last glacial period would alone account for only ~25 permil of this change (Bard, 1998; Siegenthaler et al., 1980). This implies a dominant role for changes in the global radiocarbon budget (i.e. radiocarbon production) and/or the distribution of radiocarbon between atmosphere and other carbon reservoirs (i.e. the carbon cycle).

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800 Model simulations of deglacial radiocarbon production and atmospheric radiocarbon and pCO₂ consistently indicate that past $\Delta^{14}\text{C}_{\text{atm}}$ cannot be accounted for by existing reconstructions of changing radiocarbon production alone (e.g. Hain et al., 2014; Kohler et al., 2006; Dinauer et al., 2020; Köhler et al., 2022). Earth System model simulations applying mean relative palaeomagnetic intensity (RPI) based radiocarbon production rate changes (Dinauer et al., 2020), yield only a small increase (~150 permil) in $\Delta^{14}\text{C}_{\text{atm}}$ at the LGM relative to the late Holocene (Figure J2b, black line). Similar $\Delta^{14}\text{C}_{\text{atm}}$ changes are obtained using alternative radiocarbon production histories (Hain et al., 2014; Kohler et al., 2006; Dinauer et al., 2020; Köhler et al., 2022). The widely recognised implication of these results is that additional carbon cycle changes (i.e. altered rates of carbon exchange with other reservoirs) and/or different radiocarbon production changes are required in order to account for the observed $\Delta^{14}\text{C}_{\text{atm}}$ amplitude.

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810 Turning to carbon cycle changes first: given the tight coupling of the marine and atmospheric carbon pools, altered exchange rates between the ocean and atmosphere are likely to have played a leading role in deglacial $\Delta^{14}\text{C}_{\text{atm}}$ variability (Muscheler et al., 2004). Arguably for the first time, our global average B-Atm estimates confirm such a role. Indeed, our mean ocean B-Atm estimates indicate a lower average exchange rate of radiocarbon (and likely CO₂) between the ocean and atmosphere during the last glacial period, and an increase in this exchange rate across the deglaciation (Figure J2c, black line and circles). All else being equal, an increase in ocean-atmosphere radiocarbon exchange would result in a drop in $\Delta^{14}\text{C}_{\text{atm}}$ during deglaciation, in parallel with a decrease in marine B-Atm offsets, as observed in Figure J2b, c. Such changes would have also contributed to deglacial atmospheric CO₂ rise (e.g. Muglia et al., 2018; Khaliwala et al., 2019; Brovkin et al., 2012; Ganopolski and Brovkin, 2017). As discussed above, a tentative quantification of this contribution to atmospheric CO₂ rise can be derived from the observed mean ocean B-Atm (Table 1) and modelled sensitivities (Figure J1a), as illustrated in Figure J2d (open circles and shaded region). This tentative contribution compares well with simulated CO₂ effects in the Bern3D model (INT_ALL60), bearing in mind that the simulated aging of the global ocean is only ~50% of that observed (Figure J2c, blue line and grey dashed line), and that ~50% of the simulated CO₂ signal derives from ocean ventilation impacts alone.

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835 While the observed mean ocean B-Atm estimates suggest a significant impact on the carbon cycle, the observed changes are still too small to account for the observed ~400 permil drop in $\Delta^{14}\text{C}_{\text{atm}}$ across the deglaciation. This is demonstrated by model simulations that produce a mean ocean aging of ~500 ^{14}C yrs (~63% of the observed value, **Table 1**), and yield a $\Delta^{14}\text{C}_{\text{atm}}$ increase of only ~56 permil, or ~14% of observed (**Figure 12b**, grey dashed line). Similar results have been obtained using the BICYCLE box-model (Kohler et al., 2006), which produced a $\Delta^{14}\text{C}_{\text{atm}}$ increase of only ~200 permil at the LGM (~50% of observed), despite yielding mean ocean B-Atm offsets that match our observed values (**Figure 12b,c**, dashed orange line). A mismatch was also obtained for the LGM using the CLIMBER intermediate complexity model (Ganopolski and Brovkin, 840 2017), where a 20% higher radiocarbon production rate, combined with reduced ocean ventilation causing a global mean B-Atm increase of ~800 years (again in line with our LGM estimate), coincided with a $\Delta^{14}\text{C}_{\text{atm}}$ increase of only ~280 permil (~70% of observed).

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845 Therefore, existing radiocarbon production records cannot account for past atmospheric radiocarbon variability, either alone or in conjunction with ocean ventilation changes that are consistent with global mean B-Atm estimates (**Figure 12b, c**). While there remain uncertainties in atmospheric radiocarbon reconstructions, these are likely on the order of ~1% (Reimer et al., 2020), and it seems highly unlikely that $\Delta^{14}\text{C}_{\text{atm}}$ variability over the last glacial cycle has been significantly overestimated. Similarly, it seems implausible (though clearly not impossible) that existing marine radiocarbon data significantly underestimate the magnitude of mean ocean B-Atm change since the last glacial period. The problem of closing the global radiocarbon budget since the last glacial becomes even more difficult if volcanic/metamorphic CO_2 inputs are invoked as a 850 significant contributor to the global carbon pool during the last glacial period (Stott et al., 2019; Stott et al., 2009). or if no change in ocean ventilation is inferred (Stott et al., 2021).

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855 Given the wide range of existing radiocarbon production estimates (e.g. as compiled by Dinauer et al., 2020) (**Figure 12a**, shaded area), it seems reasonable to postulate that existing production reconstructions might, on average, underestimate the amplitude of radiocarbon production rate change between the last glacial and the late Holocene, as suggested by a recent box-model study (Köhler et al., 2022). A ‘polar bias’ in radiocarbon production records derived from ice-core ^{10}Be fluxes has indeed been recently quantified separately for the geomagnetic and heliomagnetic modulations of cosmogenic production (Adolphi et al., 2023). Nevertheless, it has also been noted that correcting for the identified long-term geomagnetic bias in 860 ^{10}Be would not eliminate the mismatch between observed and modelled $\Delta^{14}\text{C}_{\text{atm}}$ at the LGM (Adolphi et al., 2023). Our idealised simulations with prescribed $\Delta^{14}\text{C}_{\text{atm}}$ allow us to infer the radiocarbon production changes that would be needed to reconcile imposed ventilation/carbon cycle changes with observed atmospheric $\Delta^{14}\text{C}_{\text{atm}}$ (**Figure 11a**, blue line). The rapid deglacial fluctuations in production that are inferred in this way (e.g. across HS1, the BA and YD) almost certainly reflect biases due to transient carbon cycle and ocean ventilation changes that have not been implemented in our idealised simulations (see section 2.5). However, the longer-term trend in inferred production rates indicates levels at the LGM that are close to the

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high end of the existing range of estimates (Figure 12a, shaded area). A similar result has recently been obtained using the BICYCLE box-model (Köhler et al., 2022). The higher radiocarbon production rates that are inferred at the LGM would have had a significant impact, possibly accounting for the majority of the deglacial atmospheric radiocarbon signal (Figure 12b, dashed blue line). The implication of these results is that a parallel closure of the radiocarbon and carbon cycles since the last glacial period might yet be obtained by exploiting the plausible range of reconstructed radiocarbon production rates (Köhler et al., 2022). Our global average B-Atm estimates provide a useful new constraint for achieving this goal.

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5 Conclusions

We present spatial interpolations of compiled radiocarbon data for a suite of time-slices spanning the last deglaciation. The primary purpose of these interpolations is to derive global average B-Atm estimates. A clear trend in global average B-Atm offsets is apparent since the LGM (when B-Atm was ~760 ¹⁴C yrs higher than modern on average), demonstrating unambiguous changes in the partitioning of radiocarbon between the ocean and atmosphere across the last deglaciation.

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The spatial interpolations cohere with previous studies in indicating a stepwise and spatially heterogenous rejuvenation of the ocean interior across the last deglaciation, and in suggesting the operation of a ‘ventilation seesaw’ between the North Atlantic and the North Pacific/Southern Ocean, especially during HS1, the YD,

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A comparison of surface-, shallow- and deep-water B-Atm trends indicates that transport changes in the upper ocean across the last deglaciation were likely modest (Marchal and Zhao, 2021a), and that B-Atm changes in upper ocean (< 2 km) were more strongly influenced by gas-exchange efficiency changes at high latitudes. In contrast, a more significant contribution from transport and/or water mass geometry is apparent in the deeper ocean, > 2km.

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The time-slice reconstructions emphasize a widespread drop in B-Atm at the onset of the BA, resulting in a global average B-Atm within ~100 ¹⁴C yrs of modern. However, model sensitivity tests indicate that a portion of this B-Atm drop could have resulted from atmospheric radiocarbon dynamics that were independent of ocean ventilation (e.g. radiocarbon production, terrestrial carbon release, etc.). The exact magnitude of this effect cannot yet be quantified, but a maximum bias of ~-190 ¹⁴C yrs relative to the LGM is estimated. Such a bias would imply that mean ocean B-Atm at the BA underestimates the true ‘ventilation age’ of the ocean.

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Model sensitivity tests further suggest a direct relationship between global average B-Atm anomalies and atmospheric CO₂ change, with a tentative average sensitivity of ~ -6.3 ppm CO₂ per 100 ¹⁴C yrs. On this basis, our global average B-Atm estimates would imply a non-negligible contribution to atmospheric CO₂ change across the last deglaciation (perhaps as much as >30% of the total observed). Global average B-Atm estimates also suggest that any ventilation contribution to atmospheric CO₂ change was concentrated during HS1, the BA and the YD, and was largely exhausted by the onset of the Holocene.

While our results serve to underline, and tentatively to quantify, the ocean's role in deglacial carbon cycle change, they also demonstrate that a complete closure and reconciliation of the radiocarbon and carbon cycles since the last glacial remains to be achieved. Our results point to the possibility that, on average, existing reconstructions may tend to underestimate radiocarbon production rates during the last glacial period. Further work to improve the accuracy of past radiocarbon production rate estimates therefore emerges as a priority.

Data availability

Data presented in this study are lodged with the PANGAEA database at: <https://www.pangaea.de>.

Author contributions

LCS designed the study, compiled and processed the radiocarbon data, and performed the interpolations with the assistance of FP. FP developed the interpolation code, and LCS and FP analysed the interpolation outputs. Numerical model runs using Bern3D were performed by AJ-T, and analysed by AJ-T, FJ and LCS. LCS wrote the manuscript with input from all co-authors.

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Competing interests

The authors declare that there are no competing financial and/or non-financial interests in relation to the work described.

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Table 1. Time-slice global average B-Atm values and anomalies *versus* modern, based on 3D interpolations, for three different data flagging approaches, as described in the main text: ‘baseline’ (all flagged data are excluded); ‘low sedimentation’ (as for baseline, but with low sedimentation rate sites included, <2cm/kyr); ‘high sedimentation’ (as for baseline, but retaining only high sedimentation rate sites, >10cm/kyr). Hypothetical corrections for maximum ‘attenuation biases’ and inferred atmospheric CO₂ impacts are shown compared with observed mean atmospheric CO₂ levels. For each data flag scenario, the total number of deglacial data points and average recorded water depth is shown as a rough indication of the impact of the data flags: as more stringent flags are applied, the data are biased to fewer points at shallower depths.

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Time slice	Baseline (n = 2579, ave. z = 1407m)			Low sedimentation (n = 2688, ave. z = 1472m)			High sedimentation (n = 2036, ave. z = 1244m)			Max. attenuation bias in model (CONST)	"Corrected" D(B-Atm), baseline	Estimated pCO ₂ impact, baseline	Obs pCO ₂ ppmv	Obs D(pCO ₂) vs PI ppmv	err ppmv
	Mean ocean B-Atm	D(B-Atm) vs PI	err	Mean ocean B-Atm	D(B-Atm) vs PI	err	Mean ocean B-Atm	D(B-Atm) vs PI	err						
	14C yrs	14C yrs	14C yrs	14C yrs	14C yrs	14C yrs	14C yrs	14C yrs	14C yrs						
HOL	1335	-24	22	1328	-31	24	1340	-19	19	25	-49	3	281.9	-3.6	0.5
EHOL	1359	0	26	1360	1	30	1320	-39	34	-52	52	-9	268.2	-17.3	0.4
YD	1484	125	37	1515	156	37	1494	135	40	28	97	-6	253.7	-31.8	0.7
BA	1435	76	32	1487	128	37	1430	71	35	-217	293	-18	241.5	-43.9	0.5
HSI	1877	518	40	1981	621	55	1795	436	58	-107	625	-39	219.7	-65.8	1.5
LGM	2123	764	37	2185	825	36	1990	631	46	-26	789	-50	194.7	-90.8	0.6

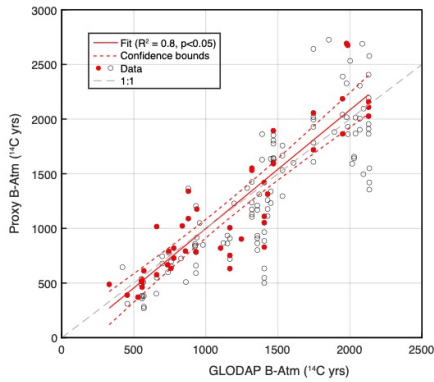
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Table 2 B-Atm and atmospheric CO₂ results for sensitivity experiments carried out using the Bern3D model with varying Southern Ocean wind, gas exchange, or global vertical diffusivity (illustrated in Figure 10(a) of the main text). Results are evaluated 2000 years after forcing is applied (see text).

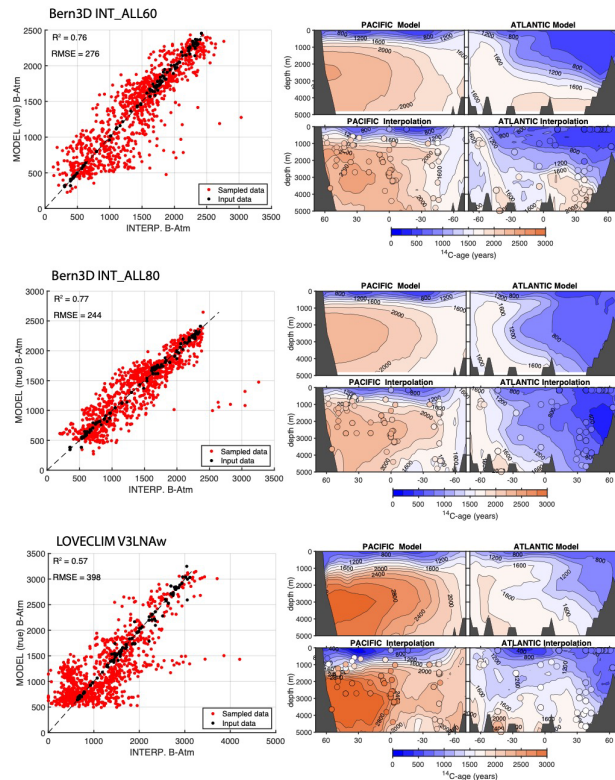
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Scenario	Ventilation parameter changed	Reduction (%)	B-Atm	pCO ₂	D(B-Atm)	D(pCO ₂)
PI (ctrl)	none	0	1428	276.5857	0	0
PI	S.O. wind	20	1495	270.2083	67	-6
		40	1548	264.4817	120	-12
		60	1577	260.3277	150	-16
		80	1588	259.9954	160	-17
PI	S.O gas exchange	20	1472.1	275.7578	44	-1
		40	1523.3	274.6609	96	-2
		60	1585.5	272.8658	158	-4
		80	1662.4	269.6315	235	-7
PI	vertical diffusivity	20	1480.3	273.7565	53	-3
		40	1523.6	271.5337	96	-5
		60	1564.6	269.6152	137	-7
		80	1597.5	268.211	170	-8

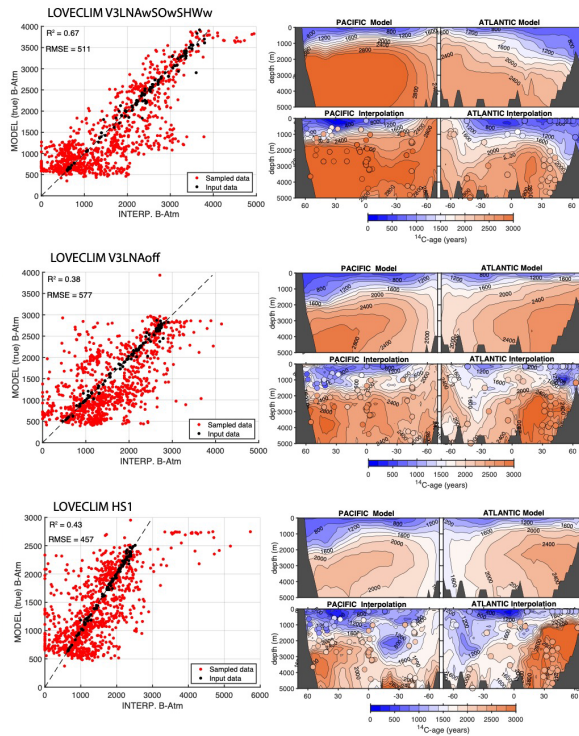


1340 Figure 1. Comparison of modern seawater (bomb-corrected, or ‘background’) B-Atm radiocarbon age offsets (Key et
 al., 2004) versus proxy-based B-Atm reconstructions using material deposited during the last ~6,000 years (black open
 circles) and the last ~1,500 years (red filled circles). Dashed grey line indicates the 1:1 trend. The linear fit to the data
 is indicated by the red line (dashed red lines show 95% confidence limits), with equation: $y = (1.1 \pm 0.1)x + (89 \pm$
 98), $R^2 = 0.81$, $p = 1.77 \times 10^{-17}$.

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1350 **Figure 2. Testing the interpolation method using numerical model outputs. Interpolations for the Bern3D and**
LOVECLIM models, from top to bottom: INT_ALL60 (this study), INT_ALL80 (this study), and V3LNAW (Men
et al., 2017). Cross plots at left show ‘true’ versus interpolated B-Atm offsets (black circles for input data, red circles
for 1000 randomly sampled locations in the ocean interior), with R^2 and RMSE shown. Contoured panels show zonally
averaged B-Atm offsets (i.e. ^{14}C -age relative to the atmosphere), for the Pacific (left) and Atlantic (right), for both the
model output (upper panels) and the interpolated reconstructions (lower panels), with input data used for the
 1355 **interpolations indicated by the filled circles.**



1360 **Figure 3. Testing the interpolation method using numerical model outputs, as for Figure 2. Interpolations for**
simulations using the LOVECLIM model, from top to bottom: V3LNAwSOWSHWw (Menviel et al., 2017), V3LNAoff
(Menviel et al., 2017), and HS1 (Menviel et al., 2018). Cross plots at left show ‘true’ versus interpolated B-Atm offsets
(black circles for input data, red circles for 1000 randomly sampled locations in the ocean interior), with R^2 and RMSE
shown. Contoured panels show zonally averaged B-Atm offsets (i.e. ^{14}C -age relative to the atmosphere), for the Pacific
(left) and Atlantic (right), for both the model output (upper panels) and the interpolated reconstructions (lower panels),
1365 **with input data used for the interpolations indicated by the filled circles.**

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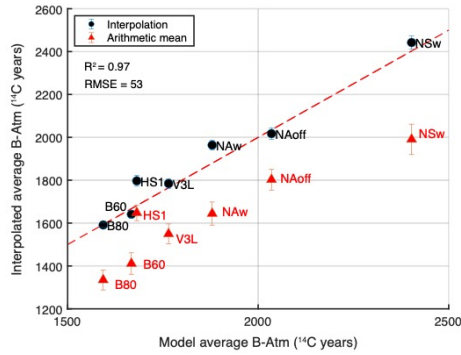


Figure 4. Comparison of global average B-Atm offsets based on interpolated fields (circles, with estimated interpolation error), and geometric averages of the input data used for the interpolations (triangles, with standard error), versus the true values for the model simulations shown in Figure 1 and an additional simulation (NAw) from (Menviel et al., 2017) (V3LNAw). The R^2 correlation coefficient, RMSE, and 1:1 line (red dashed line) are indicated.

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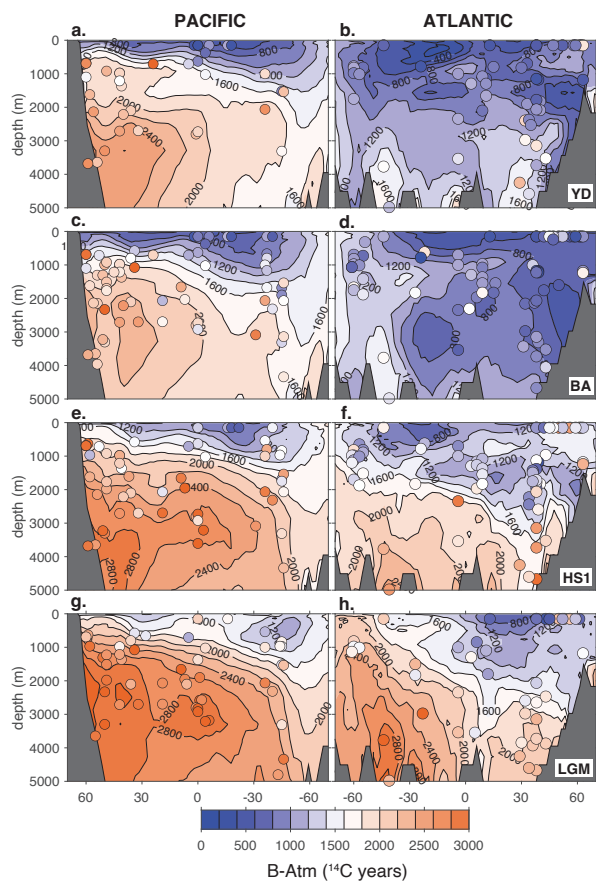


Figure 5. Zonally averaged interpolated B-Atm radiocarbon age offsets for the LGM, HS1, BA, and YD (Pacific zonal averages at left, Atlantic at right). Filled circles and shading indicate input data and values.

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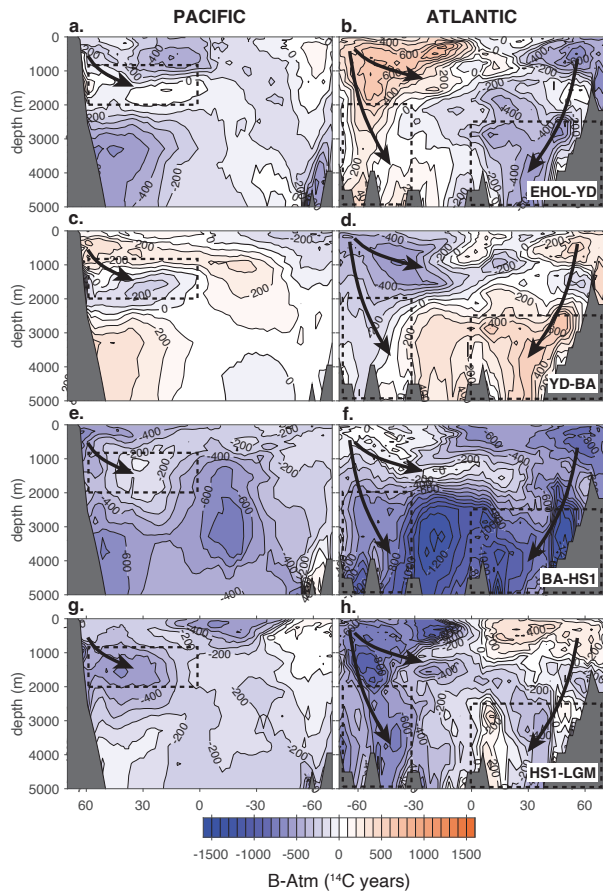


Figure 6. Offsets between spatial B-Atm interpolations for successive time-slice reconstructions. Boxed areas highlight regions of the deep North Atlantic ($>0^{\circ}\text{N}$, $>2.5\text{km}$), deep Southern Ocean ($>30^{\circ}\text{S}$, $>2\text{km}$), and intermediate North Pacific ($>0^{\circ}\text{N}$, $0.9\text{-}2\text{km}$) and deep North Pacific ($>0^{\circ}\text{N}$, $>2\text{km}$), for which regional time-series splines are illustrated in Figure 2. Broadly antiphased anomalies are indicated between the North Atlantic and Southern Ocean (especially the Atlantic sector), and between the North Atlantic and intermediate North Pacific.

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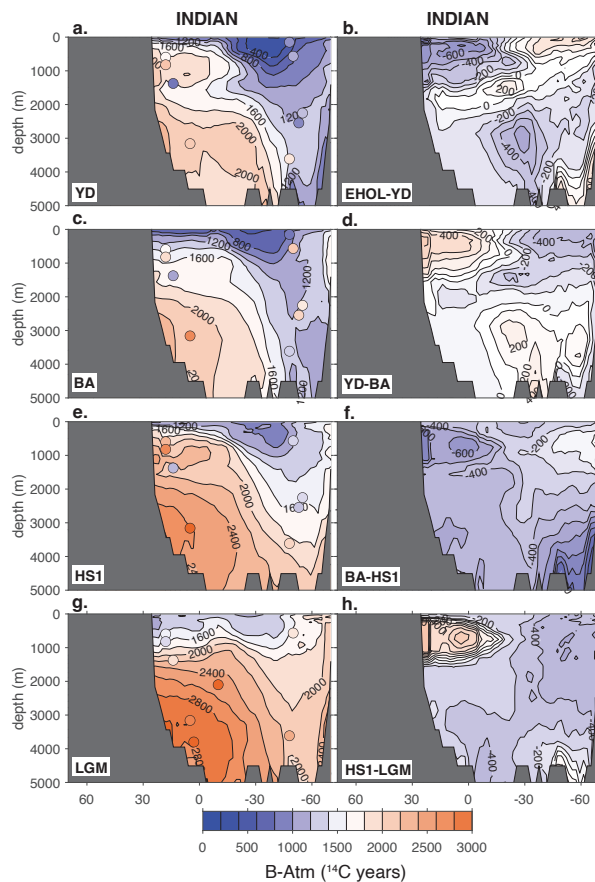


Figure 7. Zonally averaged interpolated B-Atm radiocarbon age offsets in the Indian Ocean, for the LGM, HS1, BA, and YD. Sparse data coverage is notable. Left: interpolated time-slice reconstructions (as for Figure 5), Right: Offsets between spatial B-Atm interpolations for successive time-slice reconstructions (as for Figure 6).

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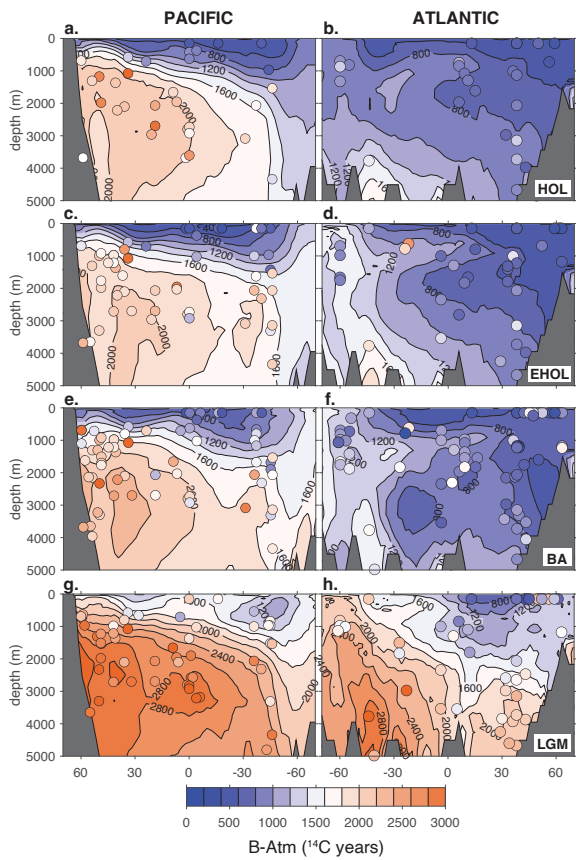


Figure 8. Zonally averaged interpolated B-Atm offsets for the LGM, BA, EHOL and HOL (Pacific zonal averages at left, Atlantic at right). Filled circles and shading indicate input data and values.

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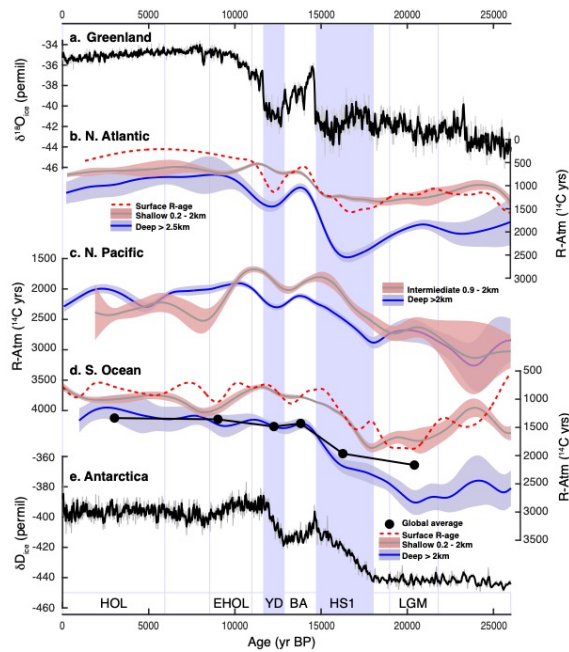


Figure 9. ‘Ventilation seesaws’ across the last deglaciation, based on cubic spline fits to compiled ocean-atmosphere radiocarbon age offsets (Skinner and Bard, 2022). Cubic splines with a stiffness factor of 10^{-9} are fit to all existing data, using updated age-models and taking into account B-Atm uncertainties, and applying the same ‘baseline’ data flags as for the time-slice interpolations. (a) Greenland temperature proxy (Svensson et al., 2008). (b) NE Atlantic shallow sub-surface reservoir ages (dashed red line; Skinner et al. (2019)); B-Atm from the North Atlantic $0.2-2\text{km}$ (grey line, red shaded area); B-Atm from the deep North Atlantic $>2.5\text{km}$ (blue line, blue shaded area). (c) B-Atm from the intermediate North Pacific $0.9-2\text{km}$ (grey line, red shaded area); B-Atm from the deep North Pacific $>2\text{km}$ (blue line, blue shaded area). (d) Mean ocean B-Atm estimates (black line and circles, Table 1 ‘baseline’ data flags), compiled shallow sub-surface reservoir ages from the Southern Ocean (dashed red line; Skinner et al. (2019)); B-Atm from the Southern Ocean $<2\text{km}$ (grey line and red shaded area); B-Atm from the deep Southern Ocean $>2\text{km}$ (blue line and blue shaded area). (e) Antarctic temperature proxy (Epica Community Members, 2004; Lemieux-Dudon et al., 2010). Vertical lines and shaded bars indicate the timing of the LGM, HSI, BA, YD, EHOL and HOL time-slice.

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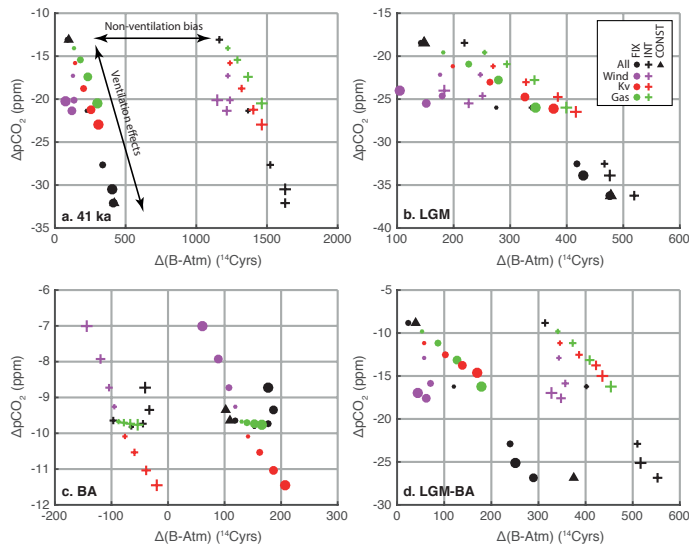
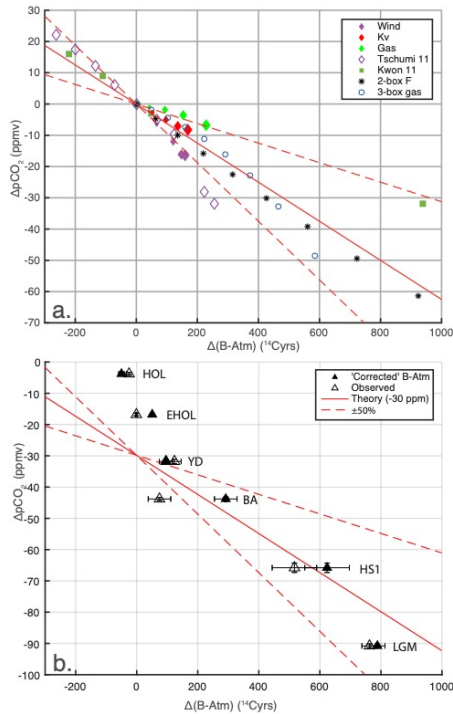


Figure 10. Model outputs from idealised 'deglacial' scenarios, illustrating hypothetical attenuation biases that are unrelated to ocean 'ventilation' effects (indicated by horizontal offsets between similar symbol colours; horizontal arrow in panel (a)). (a) 41 ka BP (coinciding with the Laschamps geomagnetic excursion); (b) the LGM; (c) the BA; (d) the offset between the LGM and the BA. Data are expressed as 1000-yr average anomalies relative to the pre-industrial period 4-5 ka BP, when atmospheric radiocarbon was relatively stable and therefore plausibly approached a pseudo-equilibrium state. In all panels, crosses indicate model experiments carried out with prescribed atmospheric radiocarbon as given by Intcal20 (Reimer et al., 2020) (INT); circles are for atmospheric radiocarbon fixed at 140 permil (FIX); and triangles are for constant pre-industrial radiocarbon production rates (CONST). Symbol colours distinguish experiments carried out with altered Southern Ocean winds (purple), vertical diffusivity (red), Southern Ocean gas-exchange efficiency (green), or all combined (black). Four symbol sizes indicate the extent of parameter reduction: 0% (control, smallest symbols), 20%, 40%, 60% and 80% (largest symbols). For CONST, only two sets of experiments were run: for control conditions, and for all tuning variables reduced by 60%. Note that the impact on B-Atm due to reduced Southern Ocean winds (purple circles) is attenuated in well equilibrated states (e.g. 41ka and LGM), because of e.g. slow adjustments in the distribution of North Atlantic sourced deep water (see text).

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455 Figure 11. (a) Theoretical and modelled sensitivity of atmospheric CO₂ anomalies and mean ocean radiocarbon
 ventilation age (B-Atm) anomalies, to air-sea gas exchange and transport changes: filled diamonds indicate model
 sensitivity tests from this study, based on step-changes under PI conditions and 2,000 years of equilibration time
 (purple, Southern Ocean wind; red, vertical diffusivity; green, Southern Ocean gas-exchange efficiency); open
 diamonds are Southern Ocean wind experiments (Tschumi et al., 2011); filled squares are wind and/or diffusivity
 1460 experiments using an idealised radiocarbon-like tracer (Kwon et al., 2011); asterisks and open circles represent 2- and
 3-box model experiments, with varying overturning rates (F) and ‘high latitude’ gas-exchange rates (gas), respectively.
 A median theoretical sensitivity is indicated by the solid and broken red lines, derived for the 2-box model results
 (Skinner and Bard, 2022), equivalent to -6.3 ± 3.2 ppm CO₂ change per 100 ¹⁴Cyrs of global radiocarbon ventilation

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age change. (b) Observed values of coincident global average B-Atm and atmospheric pCO₂, expressed as time-slice anomalies *versus* pre-industrial values: open triangles show paired observations of mean ocean B-Atm (this study) and atmospheric CO₂ (Monnin et al., 2001; Marcott et al., 2014); filled triangles show with observed B-Atm age offsets corrected for hypothetical ‘non-ventilation’ biases (see text). Solid and broken red lines in panel b show the same theoretical sensitivities as for panel (a), offset by -30 ppm to aid comparison with observed trends.

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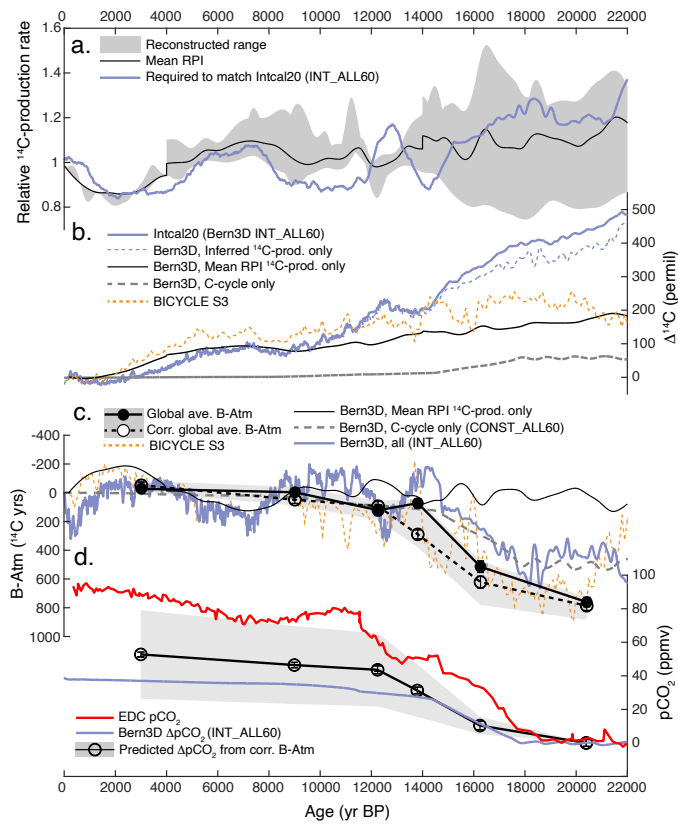


Figure 12. Observed and modelled atmospheric and marine radiocarbon, compared with radiocarbon production rates, and atmospheric CO₂. (a) Relative radiocarbon production rates: shaded area, full range of reconstructed values (Laj et al., 2000; Laj et al., 2004; Adolphi et al., 2018; Channell et al., 2018; Nowaczyk et al., 2013); solid black line, mean RPI (Laj et al., 2000; Laj et al., 2004; Channell et al., 2018; Nowaczyk et al., 2013); heavy blue line, as inferred from idealised model scenario to match Intcal20, INT_ALL60 (1kyr moving average, this study). (b) Atmospheric $\Delta^{14}\text{C}$ anomalies, normalised to modern: prescribed in the INT_ALL60 model scenario (i.e. Intcal20 (Reimer et al., 2020); solid dark blue line, this study); driven only by inferred production rates that match Intcal20 (INT_ALL60-

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1490 CONST_ALL60; dashed blue line, this study); driven only by mean RPI production rates (Dinauer et al., 2020) (solid
black line); driven only by carbon cycle/ventilation changes in the CONST-ALL60 model scenario (dashed grey line,
this study); simulated with the BICYCLE box-model (Kohler et al., 2006) using ¹⁰Be-based (Muscheler et al., 2005)
radiocarbon production estimates and a full carbon cycle scenario (dashed orange line). (c) B-Atm radiocarbon age
offset anomalies relative to modern: Bern3D and BICYCLE model outputs as for panel (b); and inferred from time-
1495 slice interpolations (black filled circles and line; this study), including a maximal correction for ‘attenuation biases’
(open black circles and dashed line), with the full range of reconstructed values (shaded area) due to uncertainties,
corrections, and alternative data flagging scenarios (Table 1). (d) Atmospheric CO₂ normalised to LGM values: from
EDC (Monnin et al., 2001) (red line); inferred from observed mean ‘corrected’ ocean B-Atm using a sensitivity of ~ -
6.3 ± 3.2 ppm per 100 ¹⁴Cyrs (open black circles and line with shaded range, this study); and simulated in the Bern3D
1500 model for the INT_ALL60 scenario that is also illustrated in panels (a)-(c).