1 2	Distorted Pacific-North American Teleconnection at the Last Glacial Maximum
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19 Abstract

20	The Pacific-North American (PNA) teleconnection is one of the most important climate
21	modes in the present climate condition, and it enables climate variations in the tropical Pacific to
22	exert significant impacts on North America. Here, we show climate simulations that the PNA
23	teleconnection was largely distorted or broken at the Last Glacial Maximum (LGM). The
24	distorted PNA is caused by a split of the westerly jet stream, which is ultimately forced by the
25	thick and large Laurentide ice sheet at the LGM. Changes in the jet stream greatly alter the
26	extratropical wave guide, distorting wave propagation from the North Pacific to North America.
27	The distorted PNA suggests that climate variability in the tropical Pacific, notably, El Niño and
28	Southern Oscillation (ENSO), would have little direct impact on North American climate at the
29	LGM.
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42 **1 Introduction**

The Pacific-Northern-American (PNA) teleconnection is the major atmospheric 43 teleconnection mode that links climate variations from the tropical Pacific to North America for 44 the present-day climate state (Horel and Wallace, 1981; Wallace and Gutzler, 1981). Especially, 45 climate variability associated with El Niño and Southern Oscillation (ENSO) exerts great 46 impacts on the North American climate through the PNA teleconnection (Henderson and 47 Robinson, 1994; Lau, 1997; Leathers et al., 1991; Straus and Shukla, 2002). It is well known that 48 the PNA is largely constrained by extratropical atmospheric flows, notably, the extratropical 49 wave guide (Held, 1983; Held et al., 2002; Hoskins and Karoly, 1981; Jin and Hoskins, 1995). 50 Thus, changes in extratropical atmospheric flows should alter the PNA under different climate 51 conditions. 52

It has been shown that greenhouse warming leads to a strengthening and a shift of the PNA 53 due to altered extratropical atmospheric flows (Allan et al., 2014; Chen et al., 2017). There has 54 55 also been a large body of works that demonstrated significant differences in extratropical atmospheric circulations in cold climates, notably, the Last Glacial Maximum (LGM). It was 56 shown that during the LGM the Aleutian low pressure system was enhanced in winter, the 57 Pacific high pressure system was weakened in summer (Yanase and Abe-Ouchi, 2007; Yanase 58 59 and Abe-Ouchi, 2010), the westerly jet shifted southward (Braconnot et al., 2007; Otto-Bliesner et al., 2006), and transient waves were weakened over the North Pacific and strengthened over 60 the North Atlantic (Justino and Peltier, 2005; Justino et al., 2005). These works suggest that the 61 62 PNA could be changed for different climate regimes. Therefore, a natural question is whether the PNA is also significantly altered due to atmospheric circulation changes at the LGM. 63

64	The LGM occurred between 23,000 and 19,000 years ago (Clark et al., 2009; Clark and
65	Mix, 2002). One of the most significant climatic characteristics at LGM is the maximum
66	expansion of mid-latitude ice sheets. Extensive ice sheets grew over North America and
67	northwestern Europe, with the Laurentide ice sheet over North America, in particular, of an ice
68	thickness of 3 to 4 kilometers (Marshall et al., 2002). Early simulations have shown that the thick
69	and large Laurentide ice sheet forced a split of the extratropical westerly jet stream into the
70	northern and southern branches (Cohmap, 1988; Kutzbach and Wright, 1985; Rind, 1987), and
71	that the jet split leads to regional climate changes over the globe, especially over North America.
72	Proxy records showed that there were more storms and precipitation associated with the southern
73	branch, causing high lake levels and increased woodlands in the southwestern United States
74	(Cohmap, 1988; Kutzbach and Wright, 1985).
75	Recent modeling studies showed that the Arctic Oscillation and storm tracks at LGM
76	differ significantly from the present (Justino and Peltier, 2005; La în éet al., 2009; Li and Battisti,
77	2008; Lüet al., 2010; Rivière et al., 2010), and that the Laurentide ice sheet can also influence
78	the Southern-Hemisphere atmospheric teleconnection and climate variability over West
79	Antarctic (Jones et al., 2018). Therefore, it is possible that changed atmospheric circulations at
80	LGM might also significantly alter the PNA and thus climate linkage between the tropical
81	Pacific and North America.
82	In the present paper, using climate simulation results, we show that the PNA is largely
83	distorted or even broken by the Laurentide ice sheet at LGM, and that ENSO had little direct
84	impact on North American climates. We will also address how the PNA is altered by the
85	Laurentide ice sheet.

86 2 Models and data

87 The simulation results from the Paleoclimate Modeling Intercomparison Project 2 (PMIP2) (Braconnot et al., 2012; Braconnot et al., 2007) and 3 (PMIP3) (Abe-Ouchi et al., 2015) are 88 utilized in this study. By comparing the PNA patterns in the Preindustrial condition (PIC) with 89 LGM simulations as well as our own sensitivity simulations, the changes in the PNA pattern at 90 LGM are identified. The horizonal resolution of the models we use are listed in table S1. For 91 92 comparison, we also use the NCEP/NCAR reanalysis data from 1988 to 2017 (Kistler et al., 2001), with horizontal resolution of 2.5 \times 2.5 $^{\circ}$. We shall mainly focus on the simulation results 93 from the Community Climate System Model version 3 (CCSM3) (Collins et al., 2006; Jones et 94 al., 2018; Otto-Bliesner et al., 2006; Yeager et al., 2006), since our sensitivity simulations are 95 performed with the same model. 96

To understand the impact of the topography of the Northern-Hemisphere glacial ice sheets 97 on the PNA, we performed a series of sensitivity simulations with different ice sheet thicknesses, 98 which are 0%, 20%, 40%, 60%, 80%, 100%, and 150% of the ice sheet thickness that was used 99 in PMIP2. Note that different ice sheet reconstructions were used in PMIP2 and PMIP3 100 simulations. PMIP2 simulations used the ICE-5G (VM2) reconstruction (Peltier, 2004), while 101 PMIP3 simulations used the ICE-6G reconstruction. In general, the ice sheet thickness in ICE-102 6G reconstruction is approximately equal to 80% of ICE-5G for most parts of the North 103 104 American region (Figure S1). In our sensitivity simulations, the case of 0% ice sheet thickness means that the thickness of the ice sheet is set to zero, but the surface albedo remains ice albedo. 105 All other conditions are the same as that in the LGM simulations of PMIP2. The model for our 106 107 sensitivity simulations is a lower-resolution version of CCSM3 (T31), with horizontal a resolution of 3.8 \times 3.8 \cdot It differs from the PMIP2 models (T42), with a horizontal resolution of 108 2.8 °×2.8 °. Although the horizontal resolution in CCSM3 T31 is lower, it can well reproduce the 109

present-day PNA pattern in the PIC simulation, consistent with the results in Magnusdottir and
Haynes (1999) and Löfverström et al. (2016). Therefore, the results here are not sensitive to
model resolutions.

Following Horel and Wallace (1981) and Wallace and Gutzler (1981), the PNA 113 teleconnection is characterized by the pointwise correlation method. The four base points that 114 represent the centers of action are located near Hawaii (20 N, 160 W), North Pacific (45 N, 115 116 165 W), Alberta (55 N, 115 W), and the Gulf Coast (30 N, 85 W), respectively. The four base points were objectively derived with teleconnectivity analysis (Sherriff-Tadano and Itoh, 117 2013; Wallace and Gutzler, 1981). To examine whether models can reasonably simulate the 118 PNA in PIC simulations and whether the PNA pathway is altered in LGM simulations, we 119 loosely define a circular region around each of the centers of North Pacific, North America and 120 the Gulf Coast (the base point is near Hawaii), with a radius of 10 degrees. For PIC simulations, 121 if a model cannot generate statistically significant correlations (coefficients greater than 0.35) 122 within the circular regions, the model is considered to have poor performance in simulating 123 PNA. For these models with good performance in simulating PNA their PIC simulations, if their 124 LGM simulations shows absence of significant correlations in the three circular regions, the PNA 125 pathway is considered to be distorted or broken at LGM. Because the PNA is most active in DJF, 126 127 our analysis below will mainly focus on the December-January-February (DJF) season.

In the present paper, all correlation analyses are conducted with monthly-mean model outputs of the last 30-year simulations. Correlation coefficient 0.35 corresponds to the 95% confidence level for 30-year correlations.

131 **3 Results**

Fig. 1 shows one-point correlation maps of 500 hPa geopotential heights in DJF, with the 132 base point near Hawaii. The correlation maps in Figs. 1a and 1b exhibit similar wave-train 133 patterns, with centers of positive and negative correlations extending from Hawaii to North 134 Pacific, Alberta, and finally to the Gulf Coast, respectively. Hence, the present-day PNA is 135 reproduced reasonably well in CCSM3. In contrast, this PNA pattern is altered dramatically in 136 137 the LGM simulation of CCSM3 (Fig. 1c). The negative correlation over North Pacific is reduced, and the center of positive correlation is rather weak and shifted to the Arctic. The most striking 138 feature in Fig. 1c is that the center of negative correlation near the Gulf Coast completely 139 disappears. The results in Fig. 1 indicate that the PNA teleconnection is largely distorted at 140 LGM. This is the most important point of the present paper. 141



Fig. 1. One-point correlation maps of 500 hPa geopotential heights in DJF in NCEP/NCAR
 reanalysis and PMIP2 CCSM3 simulations. (a) NCEP/NCAR, (b) PIC, and (c) LGM. The base
 point is near Hawaii. The correlation coefficient of 0.35 corresponds to the 95% confidence level
 for 30-year correlations.

This distorted PNA at LGM can also be seen from correlation maps for the other three base points. When the base point is located over North Pacific (Fig. S2c), the center of positive correlation over North America is shifted to northern Canada. For the base point over North

America (Fig. S2f), the negative correlations over North Pacific and the Gulf Coast are all largely reduced, and the center of positive correlation near Hawaii disappears. This result indicates a disconnection between North America and the tropical Pacific. For the base point near the Gulf Coast (Fig. S2i), a wave train is established from North Pacific to the Gulf Coast, while the center of positive correlation over North America is largely reduced, and the center of positive correlation near Hawaii is absent.

The PNA teleconnection at LGM is even completely broken in other PMIP2 models. There 156 are seven PMIP2 models that have simulations available online. According to our definition, 157 CCSM3, ECBILTCLIO, HadCM3M2, and CNRM-CM33 can reasonably reproduce the PNA in 158 their PIC simulations (Fig. 1b and Figs. S3a-c), whereas IPSL-CM4-V1-MR, FGOALS-1.0g, 159 and MIROC3.2 have poor performance. In LGM simulations, the center of negative correlation 160 over North Pacific still exists in ECBILTCLIO, HadCM3M2, and CNRM-CM33 (Figs. S3d-f), 161 although they all shift away from the North Pacific base point and are largely reduced. However, 162 163 the center of positive correlation over North America completely disappears in these plots. Moreover, the center of negative correlation near the Gulf Coast also disappears in the three 164 models. 165

PMIP3 simulations are also used to demonstrate the changes in the PNA teleconnection at LGM. There are eight PMIP3 models that have LGM simulations available online. Again, according to our definition, CCSM4, MRI-CGCM3, and MIROC-ESM can reasonably reproduce the PNA in their PIC simulations (Figs. S4a-c). The LGM simulations of CCSM4 and MRI-CGCM3 show the absence of the center of positive correlation over North America (Figs. S4d and e). The center of positive correlation in MIROC-ESM is weak and biased toward the Arctic (Fig. S4f). The center of negative correlation near the Gulf Coast is absent in MRI-CGCM3 and

MIROC-ESM. Although there is a negative center in CCSM4 (Fig. S4d), it is more like a result
 of the subtropical wave train, rather than a part of PNA. Thus, the LGM simulations in PMIP3
 models demonstrate that the PNA is either distorted or completely broken.

We have also done Empirical Orthogonal Function (EOF) and rotated EOF (REOF) analysis to examine the PNA pattern for both LGM and PIC simulations (figures not shown here). It is found that the second REOF modes in both the NCEP reanalysis and the CCSM3 PIC simulation all well represent the loading pattern of the present-day PNA. However, the second REOF in the CCSM3 LGM simulation does not show the PNA pattern. The third and fourth REOFs in the LGM simulation show teleconnections between North Pacific and Arctic as well as between North Pacific and the southern part of North America.

183 Fig. 2 illustrates PNA responses to different ice sheet thicknesses in sensitivity simulations. The PNA pattern remains for ice sheet thicknesses no more than 60% of that in PMIP2 (Figs. 2a-184 185 d). In contrast, the PNA is distorted as ice sheet thickness is increased to 80%. The center of positive correlation is shifted to the Arctic, and the center of negative correlation near the Gulf 186 Coast disappears (Fig. 2e). As ice sheet thickness is further increased to 100 % and 150% (Figs. 187 2f-g), the center of positive correlation over North America disappears. Again, the center of 188 negative correlation is more like a part of the subtropical wave train. These results of sensitivity 189 simulations suggest that the PNA is distorted or even broken as the Laurentide ice sheet is 190 sufficiently thick. 191



Fig. 2. One-point correlation maps of 500 hPa geopotential heights in DJF in sensitivity
simulations, with different ice sheet thicknesses. The base point is near Hawaii. (a) 0%, (b) 20%.
(c) 40%, (d) 60%, (e) 80%, (f) 100%, (g) 150%, and (h) PIC. The correlation coefficient of 0.35
corresponds to the 95% confidence level for 30-year correlations.

Fig. 3 summarizes correlation coefficients around the four base points for PMIP2, 197 PMIP3, and our sensitivity simulations, according to our definition above. In Fig. 3a, both 198 CCSM3 and CCSM4 show statistically significant correlations at all the four points in the PIC 199 simulations. In contrast, they all demonstrate insignificant correlations near Alberta in LGM 200 simulations. The significant correlation of CCSM4 LGM simulation near the Gulf coast is a 201 result of subtropical wave train (Fig. S4d), as mentioned above. In Fig. 3b, the correlation 202 coefficient near Alberta becomes less significant as ice sheet thickness reaches 80%. Correlation 203 coefficients at the Gulf coast are insignificant for 80% and 150% ice sheet thickness. The 204 significant correlation for 100% ice sheet thickness is a result of subtropical wave train, as shown 205 in Fig. 2f. 206



Fig. 3. Correlation coefficients at the four PNA action centers in PIC and LGM simulations for PMIP2 and PMIP3 models, with the base point near Hawaii. The negative values over North Pacific and the Gulf Coast are reversed to positive. The dashed lines correspond to 0.35, which represent the 95% confidence level. (a) CCSM3 and CCSM4, (b) sensitivity simulations, (c) PIC simulations of PMIP2 models, (d) PIC simulations of PMIP3 models, (e) LGM and PIC

simulations for well-performing PMIP2 models, and (f) LGM and PIC simulations for well-performing PMIP3 models.

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216	Figs. 3c and d shows that most PMIP2 and PMIP3 models are able to reproduce the
217	center of negative correlations over the North Pacific in their PIC simulations, except for
218	FGOALS-1.0g. FGOALS-1.0g that generates insignificant correlations at either North Pacific or
219	Alberta. CNRM-CM33 and MIROC3.2 cannot generate significant correlations near the Gulf
220	coast. Fig. 3d shows that CCSM4, MRI-CGCM3, and MIROC-ESM are able to reproduce
221	significant correlations at all four points in their PIC simulations, whereas the other 5 models
222	have insignificant correlations at either Alberta or the Gulf Coast. Figs. 3e and f show that
223	PMIP2 and PMIP3 models, which have good performance in simulating the PNA teleconnection
224	in PIC simulations, all cannot reproduce significant positive correlations at Alberta or even
225	negative correlations near the Gulf coast in the LGM simulations. These results all suggest that
226	the PNA was distorted or broken at LGM.

Because the PNA pattern is characterized by a quasi-stationary wave train from the 227 tropical Pacific to North America, the above simulation results suggest that the PNA wave-train 228 propagation is largely altered at LGM. This can be confirmed by activity fluxes of stationary 229 waves at 500 hPa calculated, using equation 7.1 in Plumb (1985) (Fig. 4), which represents the 230 propagation direction of stationary waves (Plumb, 1985). At present, the wave activity fluxes 231 have two branches for wave propagation from the North Pacific toward North America (Fig. 4a). 232 The major branch propagates northeastward, forming the PNA teleconnection, while the minor 233 branch propagates southeastward. At LGM, however, wave propagation is altered drastically. 234 Wave propagation is deflected toward the subtropics (Figs. 4b and c). This is consistent with the 235 correlation map in Fig. S2i that shows a wave train from North Pacific to the Gulf Coast. 236

237 Therefore, the distorted or broken PNA at LGM is mainly due to the deflection of wave





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Fig. 4. Stationary wave activity fluxes in PMIP2 CCSM3 simulations at 500 hPa. (a) PIC, (b) LGM, and (c) LGM – PIC. Length scales of wave activity vectors are marked in plots. Wave activity vectors are plotted as their length scales are greater than $12 \text{ m}^2 \text{ s}^{-2}$ in plots (a) and (b) and $6.5 \text{ m}^2 \text{ s}^{-2}$ in plot (c). Here, stationary wave activity fluxes are calculated with monthly-mean data.

Wave propagation is oriented by the extratropical wave guide, which in turn is 245 determined by extratropical zonal flows (Hoskins and Karoly, 1981; Jin and Hoskins, 1995). 246 Therefore, the deflection of stationary wave propagation at LGM is caused due to changes in 247 extratropical zonal flows. A comparison of zonal winds between PIC and LGM simulations 248 shows several major differences (Figs. 5a vs. 5b). First, the zonal jet stream is much stronger at 249 LGM than at present. Second, the jet is shifted equatorward at LGM, and the jet is turned 250 southeastward as it approaches the North American continent, in contrast to the northeast 251 orientation at present. Third, similar to that in early studies (Cohmap, 1988; Kutzbach and 252 Wright, 1985; Rind, 1987), the jet splits over North America with the much stronger branch 253 located in the subtropics, leaving the much weaker branch over northern Canada. These features 254 can be seen more clearly in differences of zonal winds between LGM and PIC simulations (Fig. 255 256 5c).



Fig. 5. Maps of 500 hPa zonal winds in DJF in PMIP2 CCSM3 simulations. (a) PIC, (b) LGM, and (c) LGM – PIC. Color interval: 5 m s^{-1} .

Differences of zonal winds over North America can also be illustrated with the vertical cross-261 sections along 100 W (Fig. 6). The single subtropical westerly jet in the PIC simulation (Fig. 262 6b) is split into two jets at LGM (Fig. 6c): a subtropical jet at 30 °N and 200 hPa, and a subpolar 263 jet at 63 N and between 400 and 300 hPa. The subtropical jet is intensified to a maximum wind 264 speed of 40 m s⁻¹ and is located at a lower latitude, and it is much stronger than that in the PIC 265 simulation (~ 30 m s⁻¹). The subpolar jet is much weaker, with a maximum speed of about 12 m $\frac{1}{2}$ 266 s⁻¹. The differences in zonal winds are associated with different thermal structures between LGM 267 and PIC simulations. Comparison of Figs. 6f with 6e shows that latitudinal temperature gradients 268 in the subtropics are sharper at LGM than at present. Thus, the stronger subtropical jet is 269 associated with the sharper temperature gradient. 270



Fig. 6. Vertical cross sections of DJF zonal winds and meridional temperature gradients along the
longitude of 100 °W in the NCEP/NCAR reanalysis and PMIP2 CCSM3 simulations. Top panels:
zonal winds, and bottom panels: temperature gradients. Left panels: NCEP/NCAR, middle panels:
PIC, and right panels: LGM. Zonal-wind unit is ms-1, and temperature gradient unit is K/(1000 km).

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The jet split and the equatorward shift of the major jet branch are caused by the orographic 278 forcing of the large and thick Laurentide ice sheet. Fig. S5 shows how the westerly jet responds 279 to the ice sheet thickness in the sensitivity simulations. In the case with 0% ice sheet thickness, 280 there is only a single jet in the subtropics (Fig. S5a), almost the same as that in the PIC 281 simulation. As ice sheet thickness is increased, the jet is strengthened associated with the sharper 282 meridional temperature gradient (Fig. S6), and the core of the jet becomes narrower. Significant 283 jet splitting occurs as ice sheet thickness reaches 80% (Fig. S5e). It is the reason why the 284 distortion of the PNA occurs as ice sheet thickness reaches 80%. As the ice sheet thickness is 285 increased to 100% and 150%, the jet split becomes more significant, and easterly winds begin to 286 develop over the ice sheet. 287

Note that the orographic forcing is further reinforced by the thermal forcing of the large ice sheet (Liakka, 2012). The high albedo of the ice sheet causes cold air aloft, resulting in sharper latitudinal temperature gradients in the subtropics at LGM. Thus, this enhanced temperature gradient causes a stronger subtropical jet through the thermal wind relation. Our sensitivity simulations also show that subtropical temperature gradients become sharper with increasing ice sheet thicknesses.

The split of the westerly jet acts as wave guides to orient wave propagation, as shown in Fig. 294 4. The major path of wave propagation is associated with the major jet branch. Both Figs. S2c 295 and S2i all show that a southern wave train is established along the southern jet branch from 296 North Pacific sweeping across the southern US. This wave train would lead to more storms and 297 precipitation in the American Southwest, consistent with proxy records and previous modeling 298 studies (Cohmap, 1988). The minor path of wave propagation toward the Arctic is along with the 299 northern branch (Fig. 1c), but of a much reduced strength. As such, a southern wave guide is 300 301 established along the subtropical jet, while the northern wave guide is either distorted toward the Arctic or completely broken. 302

Our sensitivity simulations demonstrate dramatic changes in the PNA wave train between 303 80% and 100% ice sheet thicknesses (Fig. 2e vs. Fig. 2f). The dramatic changes are associated 304 with the occurrence of easterly winds over the Laurentide ice sheet (Figs. 7a-c). For the case of 305 80% ice sheet thickness, westerly winds remain between the two jet streams (Fig. 7b). In 306 contrast, easterly winds appear over the ice sheet as the ice sheet thickness is increased to 100% 307 (Fig. 7c). The zero-wind line between easterly and westerly winds acts as the critical layer to 308 309 reflect stationary waves (Held, 1983). This can be addressed with calculations of critical stationary wavenumbers (Fig. 7 d-f) (eq. 6.29 in Held (1983)). The orange-red shading indicates 310

the areas where stationary waves can propagate, while the shallow-blue shading indicates the 311 areas with imaginary wavenumbers, in which propagation of stationary waves is prohibited. 312 These shallow-blue areas are associated with the easterly winds. When the ice sheet thickness is 313 60% (Fig. 7d), North Pacific and North America are dominated with positive wavenumbers, and 314 the PNA remains. For 80% ice sheet thickness, imaginary wavenumbers occur in Northeast 315 Pacific and North America (Fig. 7e), and it forces the PNA wave train distorted toward the 316 Arctic. For 100% ice sheet thickness, the subpolar region is dominated with imaginary 317 wavenumbers (Fig. 7f). It causes stationary waves reflected southeastward, leading to the 318 establishment of the southern wave train and the breaking up of the northern wave train. 319



Fig. 7. Distributions of zonal winds and stationary wavenumbers for different ice sheet thicknesses in sensitivity simulations in DJF. Top panels: zonal winds, and bottom panels: stationary wavenumbers. (a, d) 60%, (b, e) 80%, and (c, f) 100%. Zonal-wind unit is m s⁻¹, and stationary wavenumber unit is m⁻¹. The shallow-blue areas in the bottom panels have imaginary wavenumbers.

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The occurrence of easterly winds can be further illustrated with the geopotential heights at 327 500 hPa (Fig. 8). In both NCEP/NCAR reanalysis and the PIC simulation, there is only a weak 328 ridge along the west coast of North America (Figs. 8a and b). In contrast, the ridge at LGM is 329 largely enhanced and shows northwestern tilting (Fig. 8c). It is this strong ridge that leads to 330 altered zonal flows. The major branch moves equatorward, and the minor branch flows around 331 the ridge northward, resulting in the formation of easterly winds over the ice sheet and North 332 Pacific. It also can be seen in the sensitivity simulations that the west-coast ridge increases with 333 334 increasing ice sheet thickness (Fig. S7).





The distorted or broken PNA teleconnection at LGM suggests a disconnection of climate variability from the tropical Pacific to the North American continent, such that ENSO would have little direct influence on North American climates. Fig. 9 shows regression maps of surface
air temperatures (SATs) on the Nino3.4 index in DJF. At present, the remote ENSO impacts on
North American SATs through the PNA teleconnection can be identified clearly (Figs. 9a and
9b), which is characterized by an anomalously warm climate over the northwestern North
America and an anomalously cold climate over the southeastern United State. However, there are
no significant regressions of SATs over North America at LGM (Fig. 9c), except for the positive
values near the east coast.



Fig. 9. DJF SAT regressions on the Nino3.4 index in NCEP/NCAR reanalysis and PMIP2 349 CCSM3 simulations. (a) NCEP/NCAR reanalysis, (b) PIC, and (c) LGM. The regression value 350 of 0.21 corresponds to the 95% confidence level for 30-year regressions. 351 352 At present, ENSO also has important influences on North American precipitation. Similar 353 354 features can also be seen from regression maps of precipitation (Fig. 10). Fig. 10a shows precipitation regression on the Nino3.4 index in the PIC simulation. The wave train pattern of 355 precipitation is clearly shown in the plot. However, the wave train of precipitation is absent in 356 the LGM simulations (Fig. 10b). 357





Fig. 10. Precipitation regressions on the Nino3.4 index in the CCSM4 PMIP3 simulations. (a) PIC,
 and (b) LGM. Dotted areas indicate significant regressions for the 95% confidence level for 30 year regressions.

363 4 Conclusions and Discussion

We have shown in climate simulations that the large and thick Laurentide ice sheet at 364 LGM forced jet splitting and the formation of easterly winds over North America. It 365 consequently causes altered wave guides and distorted or broken PNA. It appears that the PNA 366 was separated into two teleconnections at LGM. One is from North Pacific to Arctic, and the 367 other one is from North Pacific to the southern part of North America. 368 This result suggests that ENSO would have little direct influence on North American 369 climates at LGM. Our study provides a dynamic framework to understand the PNA 370 teleconnection not only at LGM but also in other glacial periods. This understanding may help us 371 interpreting proxy records in the past. For example, a previous study on varve record in New 372 England linked the change of the intensity of interannual variability in the northeastern US 373 during the early glacial period to the change of ENSO intensity (Rittenour et al., 2000). Our 374 375 study suggests that this interannual variability is unlikely to be caused by the climate variability

from the tropical Pacific, because of the distorted or broken PNA teleconnection; instead, it 376 reflects mainly the change of local climate variability (Liu et al., 2014). Much further work is 377 needed in developing proxy records of high temporal resolutions to identify the PNA change in 378 paleoclimate records. 379 Previous works have shown weaker ENSO variability at LGM (Merkel et al., 2010; Zhu 380 et al., 2017), and that the weaker ENSO leads to altered teleconnection patterns. How the weaker 381 tropical variability would impact climates over extratropics and high-latitudes, through the 382 altered atmospheric teleconnections, deserves future studies. 383 384 Acknowledgements 385 We thank the international modeling groups of the PMIP2 and PMIP3 projects who make 386 the simulation data available. We also thank modeling groups of CMIP3 and CMIP5 whose pre-387 industry simulation data are used here. NCEP Reanalysis data are provided by the 388 NOAA/OAR/ESRL PSD, Boulder, Colorado, USA, from their Web site at 389 http://www.cdc.noaa.gov/. Y. Hu and Y. Xia are supported by the National Natural Science 390 Foundation of China (NSFC) under grants 41888101 and 41761144072, and Z. Liu is supported 391 392 by NSFC under grant 41630527. We thank the Editor and two anonymous reviewers for their insightful comments on the paper. 393

References

395	Abe-Ouchi, A., Saito, F., Kageyama, M., Braconnot, P., Harrison, S. P., Lambeck, K., Otto-
396	Bliesner, B. L., Peltier, W. R., Tarasov, L., Peterschmitt, J. Y., and Takahashi, K.: Ice-sheet
397	configuration in the CMIP5/PMIP3 Last Glacial Maximum experiments, Geosci. Model
398	Dev., 8, 3621-3637, 2015.
399	Allan, A. M., Hostetler, S. W., and Alder, J. R.: Analysis of the present and future winter Pacific-
400	North American teleconnection in the ECHAM5 global and RegCM3 regional climate
401	models, Climate Dynamics, 42, 1671-1682, 2014.
402	Braconnot, P., Harrison, S. P., Kageyama, M., Bartlein, P. J., Masson-Delmotte, V., Abe-Ouchi,
403	A., Otto-Bliesner, B., and Zhao, Y.: Evaluation of climate models using palaeoclimatic data,
404	Nature Climate Change, 2, 417, 2012.
405	Braconnot, P., Otto-Bliesner, B., Harrison, S., Joussaume, S., Peterchmitt, J. Y., Abe-Ouchi, A.,
406	Crucifix, M., Driesschaert, E., Fichefet, T., Hewitt, C. D., Kageyama, M., Kitoh, A., La în é,
407	A., Loutre, M. F., Marti, O., Merkel, U., Ramstein, G., Valdes, P., Weber, S. L., Yu, Y., and
408	Zhao, Y.: Results of PMIP2 coupled simulations of the Mid-Holocene and Last Glacial
409	Maximum – Part 1: experiments and large-scale features, Clim. Past, 3, 261-277, 2007.
410	Chen, Z., Gan, B., Wu, L., and Jia, F.: Pacific-North American teleconnection and North Pacific
411	Oscillation: historical simulation and future projection in CMIP5 models, Climate
412	Dynamics, doi: 10.1007/s00382-017-3881-9, 2017. 2017.
413	Clark, P. U., Dyke, A. S., Shakun, J. D., Carlson, A. E., Clark, J., Wohlfarth, B., Mitrovica, J. X.,
414	Hostetler, S. W., and McCabe, A. M.: The Last Glacial Maximum, Science, 325, 710-714,
415	2009.

- Clark, P. U. and Mix, A. C.: Ice sheets and sea level of the Last Glacial Maximum, Quaternary
 Science Reviews, 21, 1-7, 2002.
- Cohmap, M.: Climatic Changes of the Last 18,000 Years: Observations and Model Simulations,
 Science, 241, 1043-1052, 1988.
- 420 Collins, W. D., Bitz, C. M., Blackmon, M. L., Bonan, G. B., Bretherton, C. S., Carton, J. A.,
- 421 Chang, P., Doney, S. C., Hack, J. J., Henderson, T. B., Kiehl, J. T., Large, W. G., McKenna,
- 422 D. S., Santer, B. D., and Smith, R. D.: The Community Climate System Model Version 3
- 423 (CCSM3), Journal of Climate, 19, 2122-2143, 2006.
- Held, I. M.: Stationary and Quasi-stationary Eddies in the Extratropical Troposphere: Theory. In:
- Large-scale Dynamical Processes in the Atmosphere, B. J. Hoskins and Pearce, R. P. (Eds.),
 Academic Press, 1983.
- Held, I. M., Ting, M., and Wang, H.: Northern Winter Stationary Waves: Theory and Modeling,
 Journal of Climate, 15, 2125-2144, 2002.
- 429 Henderson, K. G. and Robinson, P. J.: Relationships between the pacific/north american
- 430 teleconnection patterns and precipitation events in the south eastern USA, International
- 431 Journal of Climatology, 14, 307-323, 1994.
- Horel, J. D. and Wallace, J. M.: Planetary-Scale Atmospheric Phenomena Associated with the
 Southern Oscillation, Monthly Weather Review, 109, 813-829, 1981.

Hoskins, B. J. and Karoly, D. J.: The Steady Linear Response of a Spherical Atmosphere to

Thermal and Orographic Forcing, Journal of the Atmospheric Sciences, 38, 1179-1196,

436 1981.

434

- 437 Jin, F. and Hoskins, B. J.: The Direct Response to Tropical Heating in a Baroclinic Atmosphere,
- Journal of the Atmospheric Sciences, 52, 307-319, 1995.

439	Jones, T. R., Roberts, W. H. G., Steig, E. J., Cuffey, K. M., Markle, B. R., and White, J. W. C.:
440	Southern Hemisphere climate variability forced by Northern Hemisphere ice-sheet
441	topography, Nature, 554, 351, 2018.
442	Justino, F. and Peltier, W. R.: The glacial North Atlantic Oscillation, Geophysical Research
443	Letters, 32, 2005.
444	Justino, F., Timmermann, A., Merkel, U., and Souza, E. P.: Synoptic Reorganization of
445	Atmospheric Flow during the Last Glacial Maximum, Journal of Climate, 18, 2826-2846,
446	2005.
447	Kistler, R., Kalnay, E., Collins, W., Saha, S., White, G., Woollen, J., Chelliah, M., Ebisuzaki,
448	W., Kanamitsu, M., Kousky, V., van den Dool, H., Jenne, R., and Fiorino, M.: The NCEP-
449	NCAR 50-year reanalysis: Monthly means CD-ROM and documentation, B Am Meteorol
450	Soc, 82, 247-267, 2001.
451	Kutzbach, J. E. and Wright, H. E.: Simulation of the climate of 18,000 years BP: Results for the
452	North American/North Atlantic/European sector and comparison with the geologic record of
453	North America, Quaternary Science Reviews, 4, 147-187, 1985.
454	La în é, A., Kageyama, M., Salas-M dia, D., Voldoire, A., Rivière, G., Ramstein, G., Planton, S.,
455	Tyteca, S., and Peterschmitt, J. Y.: Northern hemisphere storm tracks during the last glacial
456	maximum in the PMIP2 ocean-atmosphere coupled models: energetic study, seasonal cycle,
457	precipitation, Climate Dynamics, 32, 593-614, 2009.
458	Lau, NC.: Interactions between Global SST Anomalies and the Midlatitude Atmospheric
459	Circulation, B Am Meteorol Soc, 78, 21-34, 1997.

460	Leathers, D. J., Yarnal, B., and Palecki, M. A.: The Pacific/North American Teleconnection
461	Pattern and United States Climate. Part I: Regional Temperature and Precipitation
462	Associations, Journal of Climate, 4, 517-528, 1991.
463	Li, C. and Battisti, D. S.: Reduced Atlantic Storminess during Last Glacial Maximum: Evidence
464	from a Coupled Climate Model, Journal of Climate, 21, 3561-3579, 2008.
465	Liakka, J.: Interactions between topographically and thermally forced stationary waves:
466	implications for ice-sheet evolution, Tellus A: Dynamic Meteorology and Oceanography,
467	64, 11088, 2012.
468	Liu, Z., Lu, Z., Wen, X., Otto-Bliesner, B. L., Timmermann, A., and Cobb, K. M.: Evolution and
469	forcing mechanisms of El Niño over the past 21,000 years, Nature, 515, 550, 2014.
470	Löfverström, M., Caballero, R., Nilsson, J., and Messori, G.: Stationary Wave Reflection as a
471	Mechanism for Zonalizing the Atlantic Winter Jet at the LGM, Journal of the Atmospheric
472	Sciences, 73, 3329-3342, 2016.
473	Lü, JM., Kim, SJ., Abe-Ouchi, A., Yu, Y., and Ohgaito, R.: Arctic Oscillation during the Mid-
474	Holocene and Last Glacial Maximum from PMIP2 Coupled Model Simulations, Journal of
475	Climate, 23, 3792-3813, 2010.
476	Magnusdottir, G. and Haynes, P. H.: Reflection of Planetary Waves in Three-Dimensional
477	Tropospheric Flows, Journal of the Atmospheric Sciences, 56, 652-670, 1999.
478	Marshall, S. J., James, T. S., and Clarke, G. K.: North American ice sheet reconstructions at the
479	Last Glacial Maximum, Quaternary Science Reviews, 21, 175-192, 2002.
480	Merkel, U., Prange, M., and Schulz, M.: ENSO variability and teleconnections during glacial
481	climates, Quaternary Science Reviews, 29, 86-100, 2010.

482	Otto-Bliesner, B. L., Brady, E. C., Clauzet, G., Tomas, R., Levis, S., and Kothavala, Z.: Last
483	Glacial Maximum and Holocene Climate in CCSM3, Journal of Climate, 19, 2526-2544,
484	2006.
485	Peltier, W. R.: Global glacial isostasy and the surface of the ice-age Earth: The ICE-5G (VM2)
486	model and GRACE, Annual Review of Earth and Planetary Sciences, 32, 111-149, 2004.
487	Plumb, R. A.: On the Three-Dimensional Propagation of Stationary Waves, Journal of the
488	Atmospheric Sciences, 42, 217-229, 1985.
489	Rind, D.: Components of the ice age circulation, Journal of Geophysical Research: Atmospheres,
490	92, 4241-4281, 1987.
491	Rittenour, T. M., Brigham-Grette, J., and Mann, M. E.: El Niño-Like Climate Teleconnections in
492	New England During the Late Pleistocene, Science, 288, 1039-1042, 2000.
493	Rivière, G., La n é, A., Lapeyre, G., Salas-Mélia, D., and Kageyama, M.: Links between Rossby
494	Wave Breaking and the North Atlantic Oscillation–Arctic Oscillation in Present-Day and
495	Last Glacial Maximum Climate Simulations, Journal of Climate, 23, 2987-3008, 2010.
496	Sherriff-Tadano, S. and Itoh, H.: Teleconnection Patterns Appearing in the Streamfunction Field,
497	SOLA, 9, 115-119, 2013.
498	Straus, D. M. and Shukla, J.: Does ENSO Force the PNA?, Journal of Climate, 15, 2340-2358,
499	2002.
500	Wallace, J. M. and Gutzler, D. S.: Teleconnections in the Geopotential Height Field during the
501	Northern Hemisphere Winter, Monthly Weather Review, 109, 784-812, 1981.
502	Yanase, W. and Abe-Ouchi, A.: The LGM surface climate and atmospheric circulation over East
503	Asia and the North Pacific in the PMIP2 coupled model simulations, Clim. Past, 3, 439-451,
504	2007.

505	Yanase, W. and Abe-Ouchi, A.: A Numerical Study on the Atmospheric Circulation over the
506	Midlatitude North Pacific during the Last Glacial Maximum, Journal of Climate, 23, 135-
507	151, 2010.

- Yeager, S. G., Shields, C. A., Large, W. G., and Hack, J. J.: The Low-Resolution CCSM3,
- 509 Journal of Climate, 19, 2545-2566, 2006.
- 510 Zhu, J., Liu, Z., Brady, E., Otto-Bliesner, B., Zhang, J., Noone, D., Tomas, R., Nusbaumer, J.,
- 511 Wong, T., Jahn, A., and Tabor, C.: Reduced ENSO variability at the LGM revealed by an
- isotope-enabled Earth system model, Geophysical Research Letters, 44, 6984-6992, 2017.