



# 1 **Mysteriously high $\Delta^{14}\text{C}$ of the glacial atmosphere: Influence of** 2 **$^{14}\text{C}$ production and carbon cycle changes**

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10

11 **Abstract.** Despite intense focus on the ~190 permil drop in atmospheric  $\Delta^{14}\text{C}$  across the deglacial “mystery  
12 interval”, the specific mechanisms responsible for the apparent  $\Delta^{14}\text{C}$  excess in the glacial atmosphere have received  
13 considerably less attention. The computationally efficient Bern3D earth system model of intermediate complexity,  
14 designed for long-term climate simulations, allows us to address a very fundamental but still elusive question  
15 concerning the atmospheric  $\Delta^{14}\text{C}$  record: How can we explain the persistence of relatively high  $\Delta^{14}\text{C}$  values during  
16 the millennia after the Laschamp event? Large uncertainties in the pre-Holocene  $^{14}\text{C}$  production rate, as well as in  
17 the older portion of the  $\Delta^{14}\text{C}$  record, complicate our qualitative and quantitative interpretation of the glacial  $\Delta^{14}\text{C}$   
18 elevation. Here we begin with sensitivity experiments that investigate the controls on atmospheric  $\Delta^{14}\text{C}$  in more  
19 idealized settings. We show that the long-term process of sedimentation may be much more important to the  
20 simulation of  $\Delta^{14}\text{C}$  than had been previously thought. In order to provide a bounded estimate of glacial  $\Delta^{14}\text{C}$  change,  
21 the Bern3D model was integrated with five available estimates of the  $^{14}\text{C}$  production rate as well as reconstructed  
22 and hypothesized paleoclimate forcing. Model results demonstrate that none of the available reconstructions of past  
23 changes in  $^{14}\text{C}$  production can reproduce the elevated  $\Delta^{14}\text{C}$  levels during the last glacial. In order to increase  
24 atmospheric  $\Delta^{14}\text{C}$  to glacial levels, a drastic reduction of air-sea exchange efficiency in the polar regions must be  
25 assumed, though discrepancies remain for the portion of the record younger than ~33 kyr BP. We end with an  
26 illustration of how the  $^{14}\text{C}$  production rate would have had to evolve to be consistent with the  $\Delta^{14}\text{C}$  record, by  
27 combining an atmospheric radiocarbon budget with the Bern3D model. The overall conclusion is that the remaining  
28 discrepancies with respect to glacial  $\Delta^{14}\text{C}$  may be linked to an underestimation of  $^{14}\text{C}$  production and/or a biased-  
29 high reconstruction of  $\Delta^{14}\text{C}$  over the time period of interest. Alternatively, we appear to still be missing an important  
30 carbon cycle process for atmospheric  $\Delta^{14}\text{C}$ .

31

## 32 **1 Introduction**

33

34 The cosmogenic radionuclide radiocarbon ( $^{14}\text{C}$ ) is a powerful tracer for the study of several ocean processes including  
35 deep ocean circulation and ventilation. Past changes in atmospheric  $^{14}\text{C}/\text{C}$  (i.e.,  $\Delta^{14}\text{C}$ , in permil; corresponding to  $\Delta$   
36 from Stuiver and Polach, 1977), as recorded in absolutely dated tree rings, plant macrofossils, speleothems, corals,



37 and foraminifera, have been interpreted as possibly reflecting real changes in the ocean's large-scale overturning  
38 circulation (Siegenthaler et al., 1980). The extended 54,000-year record of atmospheric  $\Delta^{14}\text{C}$  from the latest IntCal  
39 compilation (i.e., IntCal13; Reimer et al., 2013) and from two Hulu Cave stalagmites (Cheng et al., 2018) suggests  
40 that large millennial-scale variations in  $\Delta^{14}\text{C}$  have occurred during the last glacial, compared to the relatively small  
41 (~30 ppm) change in atmospheric  $\text{CO}_2$  over the same time period (Fig. 1). When interpreting the implications of such  
42 changes, it is important to note that atmospheric  $\Delta^{14}\text{C}$  is controlled not only by global carbon cycle processes but also  
43 by variations in the atmospheric  $^{14}\text{C}$  production rate. Therefore, the use of atmospheric  $\Delta^{14}\text{C}$  as an indicator of past  
44 oceanic conditions, particularly those associated with air-sea exchange efficiency and deep ocean ventilation rates,  
45 requires reliable estimates of the  $^{14}\text{C}$  production rate over time.

46  
47 The vast majority of all  $^{14}\text{C}$  production changes are the result of either solar or geomagnetic modulation of  
48 the cosmic ray flux reaching the Earth (Masarik and Beer, 1999; Poluanov et al., 2016). Fig. 1 shows several different  
49 proxy records of the global production rate of  $^{14}\text{C}$  in relative units covering the full range of the  $^{14}\text{C}$  dating method,  
50 based on geomagnetic field data from marine sediments (Laj et al., 2000; Laj et al., 2004; Nowaczyk et al., 2013;  
51 Channell et al., 2018) and on  $^{10}\text{Be}$  and  $^{36}\text{Cl}$  measurements in polar ice cores (Adolphi et al., 2018). A fundamental  
52 difference between these reconstruction methods is that paleointensity-based estimates of the  $^{14}\text{C}$  production rate, by  
53 definition, do not reflect changes in the solar modulation of the cosmic radiation, whereas ice-core  $^{10}\text{Be}$ -based  
54 estimates give the combined influence of solar and geomagnetic modulation on radionuclide production. Of note is  
55 the striking coherence in all three records ( $\Delta^{14}\text{C}$ , paleointensity-based production, and ice-core  $^{10}\text{Be}$ -based production)  
56 of the Laschamp excursion (~41 kyr BP), when the Earth's geomagnetic dipole field briefly reversed and its intensity  
57 was close to zero (Nowaczyk et al., 2012; Laj et al., 2014). According to reconstructions and production rate models,  
58 this large geomagnetic event caused a doubling of the  $^{14}\text{C}$  production rate, leading to the highest  $\Delta^{14}\text{C}$  values over the  
59 last 54 kyr. Relatively high  $\Delta^{14}\text{C}$  values continued until ~25 kyr BP, then gradually diminished to preindustrial levels,  
60 interrupted by a sharp drop in  $\Delta^{14}\text{C}$  during the so-called "mystery interval" ~17.5 to 14.5 kyr BP (coincident with  
61 Heinrich Stadial 1 [HS1]; Broecker and Barker, 2007). While the Laschamp geomagnetic excursion appears to be  
62 responsible for the  $\Delta^{14}\text{C}$  peak at ~41 kyr BP, the production rate estimates during much of the pre-Holocene period  
63 are subject to considerable uncertainty.

64  
65 Paleointensity-based reconstructions are sensitive to coring disturbances of unconsolidated sediments. The  
66 last 50 kyr are represented by the uppermost, relatively slushy centimeters of marine sediment cores (Channell et al.,  
67 2018). Channell et al. (2018) used only highly accumulating sites and different coring equipment so as to avoid some  
68 of these problems, and reached very different conclusions than, e.g., Laj et al. (2000). Similarly, ice-core  $^{10}\text{Be}$  records  
69 are affected by changes in the transport and deposition of  $^{10}\text{Be}$ , which may overprint the production rate changes (e.g.,  
70 Heikkilä et al., 2013). Furthermore, in order to calculate the ice-core  $^{10}\text{Be}$  deposition fluxes, snow accumulation rates  
71 must be known for each specific ice core, which themselves have uncertainties on the order of 10 to 20 percent that  
72 propagate into the ice-core  $^{10}\text{Be}$  fluxes (Gkinis et al., 2014; Rasmussen et al., 2013). The large uncertainties associated  
73 with the reconstruction of past changes in  $^{14}\text{C}$  production lead to two distinct interpretation problems with respect to



74 the atmospheric  $\Delta^{14}\text{C}$  record: (1) the specific mechanisms responsible for high glacial  $\Delta^{14}\text{C}$  levels, and (2) the extent  
75 to which production changes contributed to the deglacial  $\Delta^{14}\text{C}$  decline.

76

77 Earlier model studies have focused heavily on the  $\sim 190$  permil drop in atmospheric  $\Delta^{14}\text{C}$  across the mystery  
78 interval and on the deglacial trends in  $\Delta^{14}\text{C}$  after HS1 (Muscheler et al., 2004; Broecker and Barker, 2007; Skinner et  
79 al., 2010; Mariotti et al., 2016; Delaygue et al., 2003; Marchal et al., 2001; Huiskamp and Meissner, 2012; Hain et al.,  
80 2014). Historically, the younger portion of the  $\Delta^{14}\text{C}$  record has received more attention than the glacial section because  
81 of the early emphasis on the general climatic trends of the North Atlantic stadials (HS1 and the Younger Dryas [YD])  
82 and the Bølling-Allerød (BA) warm period, and on the important role of an exceptionally aged ( $^{14}\text{C}$ -depleted) deep-  
83 water mass in the pulsed rise of atmospheric  $\text{CO}_2$  during the last glacial termination (e.g., Skinner et al., 2017). Less  
84 research over the last few decades has studied the mechanisms responsible for high glacial  $\Delta^{14}\text{C}$  levels. The model  
85 studies that are available point out the difficulties in simulating the correct glacial  $\Delta^{14}\text{C}$  levels (Hughen et al., 2004;  
86 Köhler et al., 2006). These studies demonstrate with box models that glacial levels of  $\Delta^{14}\text{C}$  cannot be attained without  
87 invoking significant changes in ocean circulation, air-sea gas exchange, and carbonate sedimentation. However, the  
88 box models were not able to reproduce  $\Delta^{14}\text{C}$  values higher than 700 permil, and these results still need to be scrutinized  
89 with models of higher complexity. To our knowledge, no three-dimensional ocean biogeochemical model has yet  
90 simulated the 50,000-year record of  $\Delta^{14}\text{C}$ . Many questions remain unanswered, in particular: What mechanism can  
91 account for the persistence of relatively high  $\Delta^{14}\text{C}$  values during the millennia after the Laschamp excursion?

92

93 The expected time scale for sustaining elevated levels of atmospheric  $\Delta^{14}\text{C}$  after a production peak is on the  
94 order of thousands of years, a time scale tied to the mean lifetime of  $^{14}\text{C}$  ( $\sim 8223$  years; Audi et al., 2003; Bé et al.,  
95 2013) and the time required for deep ocean ventilation (up to about 1000 years). Specifically, Muscheler et al. (2004)  
96 demonstrate that the characteristic time constant for equilibration of  $\Delta^{14}\text{C}$  after a perturbation in atmospheric  
97 production is 5000 years. By this analysis, the Laschamp event, which lasted only about 1500 to 2000 years (Laj et  
98 al., 2000), was insufficient to sustain the high  $\Delta^{14}\text{C}$  values observed over the next  $\sim 15,000$  years. The lack of significant  
99 changes (only  $\sim 10$  percent) in atmospheric  $\text{CO}_2$  during the time period of interest raises the question of what causes  
100 variations in atmospheric  $\Delta^{14}\text{C}$ , but not  $\text{CO}_2$ , on millennial time scales? The obvious answers are: cosmic ray  
101 modulation and air-sea gas exchange. Ultimately, no explanation for high glacial  $\Delta^{14}\text{C}$  levels can be complete in the  
102 absence of more robust estimates of the pre-Holocene  $^{14}\text{C}$  production rate, as well as a good understanding of the  
103 ocean carbon cycle under glacial climate conditions.

104

105 One of the major challenges associated with modelling glacial-interglacial climate cycles is that it is currently  
106 not possible to reproduce climate and atmospheric  $\text{CO}_2$  variations on the basis of orbital forcing alone. Problems  
107 include the complexity of the Earth system, making it difficult to represent all the relevant processes in models, and  
108 the long time scales involved, making simulations covering tens of thousands of years costly in computation time.  
109 Glacial-interglacial simulations with dynamic ocean and land models of intermediate complexity have begun to  
110 emerge, but these models are not yet able to reproduce the reconstructed variations in important proxy data or the



111 timing of CO<sub>2</sub> variations during the last glacial termination (Brovkin et al., 2012; Ganopolski and Brovkin, 2017;  
112 Menviel et al., 2012). A wide variety of mechanisms related to the ocean and land carbon cycles including exchange  
113 processes with marine sediments, coral reefs, and the lithosphere have been proposed to explain the glacial-interglacial  
114 variations in atmospheric CO<sub>2</sub> (Archer et al., 2000; Fischer et al., 2010; Wallmann et al., 2016), but the causes remain  
115 elusive. As long as there are still large gaps in our understanding of the glacial climate and associated ocean carbon  
116 cycle, a convenient way to examine the impact of the possible mechanisms on atmospheric CO<sub>2</sub> levels, and here on  
117  $\Delta^{14}\text{C}$ , is to perform sensitivity experiments and scenario-based simulations with models. This allows us to investigate  
118 specific phenomena in more idealized settings, permitting us to investigate in detail which parameters and processes  
119 are most important in controlling atmospheric  $\Delta^{14}\text{C}$  levels.

120

121 In this paper we extend previous modelling efforts concerning the record of atmospheric  $\Delta^{14}\text{C}$  with respect  
122 to three issues: (1) the sensitivity of the  $\Delta^{14}\text{C}$  response to carbon cycle changes and the potential importance of marine  
123 sediments, (2) the simulation of  $\Delta^{14}\text{C}$  covering the time range of the IntCal13 radiocarbon calibration curve (50,000  
124 years), the primary focus being the explanation of high glacial  $\Delta^{14}\text{C}$  levels, and (3) a new model-based reconstruction  
125 of the  $^{14}\text{C}$  production rate for the last 50 kyr as inferred from an atmospheric radiocarbon budget. In the following  
126 sections we first introduce the Bern3D earth system model of intermediate complexity and describe the carbon cycle  
127 scenarios for forcing it. We then use step changes in the  $^{14}\text{C}$  production rate and in selected parameters of the ocean  
128 carbon cycle model to gain insight into the transient and equilibrium response of atmospheric  $\Delta^{14}\text{C}$ . After these  
129 sensitivity experiments we present the results of paleoclimate simulations forced by available reconstructions of past  
130 changes in  $^{14}\text{C}$  production together with well-known and hypothesized carbon cycle changes accompanying glacial-  
131 interglacial climate cycles. Finally, we present results for a first attempt to reconstruct the history of the  $^{14}\text{C}$  production  
132 rate using the Bern3D model forced with reconstructed variations in atmospheric  $\Delta^{14}\text{C}$  and CO<sub>2</sub>. We end with a  
133 comparison of three fundamentally different (model-based, paleointensity-based, and ice-core  $^{10}\text{Be}$ -based)  
134 reconstructions of atmospheric  $^{14}\text{C}$  production.

135

## 136 **2 Materials and methods**

137

### 138 **2.1 Brief description of the Bern3D model**

139

140 Simulations are performed with the computationally efficient Bern3D earth system model of intermediate complexity  
141 (version 2.0), which is designed for long-term climate simulations over several tens of thousands of years. The Bern3D  
142 couples a frictional geostrophic 3-D ocean general circulation model (Edwards et al., 1998; Edwards and Marsh, 2005;  
143 Müller et al., 2006), a 2-D energy-moisture balance atmosphere model (Ritz et al., 2011), an ocean carbon cycle model  
144 (Müller et al., 2008; Tschumi et al., 2008; Parekh et al., 2008), a chemically active 10-layer ocean sediment model  
145 (Heinze et al., 1999; Tschumi et al., 2011; Roth et al., 2014; Jeltsch-Thömmes et al., 2019), and a four-box model  
146 representing carbon stocks in the terrestrial biosphere (Siegenthaler and Oeschger, 1987). The coarse-resolution ocean  
147 model is implemented on a 41 x 40 horizontal grid, with 32 logarithmically spaced layers in the vertical. The seasonal



148 cycle is resolved with 96 time steps per year. The tracers carried in this model study include temperature, salinity,  
149 dissolved inorganic carbon (DIC), dissolved organic carbon (DOC), carbon isotopes ( $^{13}\text{C}$  and  $^{14}\text{C}$ ) of DIC and DOC,  
150 alkalinity (ALK), phosphate, silicate, iron, dissolved oxygen ( $\text{O}_2$ ), preformed dissolved oxygen ( $\text{O}_{2,\text{pre}}$ ), and an “ideal  
151 age” tracer. The ideal age is set to zero in the surface layer, increased by  $\Delta t$  in all interior grid cells at each time step  
152 of duration  $\Delta t$ , and transported by advection, diffusion, and convection. For a more complete description of the  
153 Bern3D model, the reader is referred to Appendix A.

154

## 155 2.2 Implementation of the $^{14}\text{C}$ tracer

156

157 Natural radiocarbon ( $^{14}\text{C}$ ) is a cosmogenic radionuclide produced in the atmosphere by cosmic radiation. Once  
158 oxidized to  $^{14}\text{CO}_2$ , it participates in the global carbon cycle. Atmospheric  $^{14}\text{CO}_2$  invades the ocean by air-sea gas  
159 exchange, where it is subject to the same physical and biogeochemical processes that affect DIC. The only difference  
160 is that  $^{14}\text{C}$  is lost by radioactive decay (half-life of  $5700 \pm 30$  years; Audi et al., 2003; Bé et al., 2013). The governing  
161 natural processes, namely, atmospheric  $^{14}\text{C}$  production, air-sea gas exchange, physical transport and mixing in the  
162 water column, biological production and export of particulate and dissolved matter from the surface ocean, particle  
163 flux through the water column, particle deposition on the sea floor, remineralization and dissolution in the water  
164 column and the sediment pore waters, and vertical sediment advection and sediment accumulation, are explicitly  
165 represented in the Bern3D model (see Fig. 2).

166

167 Modelled  $^{14}\text{C}$  is often expressed relative to carbon, which is convenient for comparison to radiocarbon  
168 measurements which are generally reported as  $\Delta^{14}\text{C}$ . The purpose of  $\Delta^{14}\text{C}$ , i.e., the fractionation-corrected ratio of  
169  $^{14}\text{C}/\text{C}$  relative to that of the AD 1950 atmosphere, is to isolate the effect of radioactive decay. In this study,  $\Delta^{14}\text{C}$  is  
170 treated as a diagnostic variable using the two-tracer approach of OCMIP-2. Rather than modelling only the  $^{14}\text{C}/\text{C}$  ratio,  
171 the  $^{14}\text{C}$  tracer is carried independently from the carbon tracer. It is normalized by the standard ratio of the preindustrial  
172 atmosphere ( $^{14}r_{\text{std}} = 1.170 \times 10^{-12}$ ; Orr et al., 2017) in order to minimize the numerical error of carrying very small  
173 numbers. It is also fractionation corrected. For comparison to observations,  $\Delta^{14}\text{C}$  is calculated from the normalized  
174 and fractionation-corrected modelled  $^{14}\text{C}$  concentration:

175

$$176 \Delta^{14}\text{C} = 1000(^{14}r' - 1) \quad (1)$$

177

178 where  $^{14}r'$  is the ratio of  $^{14}\text{C}/\text{C}$  divided by  $^{14}r_{\text{std}}$ . The source-minus-sink equation for the atmospheric carbon tracer  
179 (i.e., gaseous  $\text{CO}_2$ ) is analogous to that for  $^{14}\text{C}$ , except that the equation omits the source–sink terms due to atmospheric  
180 production and radioactive decay. With respect to the ocean carbon cycle model, the oceanic carbon tracer sees a  
181 constant input of DIC from weathering on land, whereas there is no supply of  $^{14}\text{C}$  to the ocean from terrestrial  
182 weathering.

183



184 In the preindustrial spin-up simulation needed to initialize the Bern3D model, atmospheric CO<sub>2</sub> is held  
185 constant at 278.05 ppm and atmospheric  $\Delta^{14}\text{C}$  at 0 permil. During this integration time the ocean inventories of carbon  
186 and  $^{14}\text{C}$  adjust to the forcing fields. The resulting changes after >50,000 years of integration are negligibly small. Fig.  
187 3 shows the steady-state  $\Delta^{14}\text{C}$  distribution in the surface (< 100 m) and deep (> 1500 m) ocean for the preindustrial  
188 control run. The large-scale distribution of modelled  $\Delta^{14}\text{C}$  broadly resembles the observed pattern in the Global Ocean  
189 Data Analysis Project (GLODAP; Key et al., 2004). That final state (i.e., the end of the preindustrial spin-up) is used  
190 to diagnose the  $^{14}\text{C}$  production rate for the preindustrial atmosphere, such that the rate of  $^{14}\text{C}$  production is balanced  
191 by radioactive decay and the net fluxes out of the atmosphere. For transient simulations, an adjustable scale factor is  
192 applied to the preindustrial steady-state value in order to account for production changes induced by solar and/or  
193 geomagnetic modulation. The former is derived from, e.g., available reconstructions of the  $^{14}\text{C}$  production rate in  
194 relative units, as detailed in Sect. 2.5.

195

### 196 **2.3 Model configurations**

197

198 We focus in this paper on the response of atmospheric  $\Delta^{14}\text{C}$  to changes in  $^{14}\text{C}$  production and the ocean carbon cycle.  
199 For a deeper mechanistic understanding of the driving processes, step response experiments are first performed (see  
200 Sect. 3.1). These simulations include perturbations of the steady-state  $^{14}\text{C}/\text{C}$  distribution under preindustrial  
201 conditions. We investigate the impact of step changes in (1) the  $^{14}\text{C}$  production rate (“higher production” scenario),  
202 (2) wind stress and vertical diffusivity (“reduced deep ocean ventilation” scenario), and (3) the gas transfer velocity  
203 (“enhanced permanent sea ice cover” scenario). After a step change at time 0, the simulations are run to near-  
204 equilibrium over a 50,000-year integration. The following model configurations and therefore exchanging carbon  
205 reservoirs are considered: atmosphere–ocean (OCN), atmosphere–ocean–land (OCN-LND), atmosphere–ocean–  
206 sediment (OCN-SED), and atmosphere–ocean–land–sediment (ALL).

207

208 Next we examine the influence of changes that are transient in nature. We simulate atmospheric  $\Delta^{14}\text{C}$  over  
209 the full range of the  $^{14}\text{C}$  dating method (i.e., 50 to 0 kyr BP) (see Sect. 3.2 and 3.3). These transient simulations are  
210 initialized at 70 kyr BP using model configuration ALL, and forced by reconstructed changes in  $^{14}\text{C}$  production (see  
211 Sect. 2.5) over a 70,000-year integration. The first 20,000 years of the integration are considered a spin-up. Although  
212 the full record is simulated, we focus our analysis on the millennial-scale variation in atmospheric  $\Delta^{14}\text{C}$  before  
213 incipient deglaciation at ~18 kyr BP. Eight model runs are carried out for each production rate reconstruction, using  
214 different combinations of forcing fields and parameter values as described next.

215

### 216 **2.4 Carbon cycle scenarios**

217

218 In our transient simulations with the Bern3D model, eight scenarios based on different assumptions about the global  
219 carbon cycle are considered, the details of which are summarized in Table 1. The goal is to investigate the extent to



220 which changes in the ocean carbon cycle could explain high glacial  $\Delta^{14}\text{C}$  levels. We therefore consider a wide range  
221 of scenarios, including some extreme cases.

222

223 In the first scenario (MOD), the model is run with fixed preindustrial boundary conditions for the Earth's  
224 orbital parameters, radiative forcing due to well-mixed greenhouse gases, and ice sheet extent. As a consequence,  
225 atmospheric  $\text{CO}_2$  remains approximately constant at the preindustrial level of 278.05 ppm over the simulation. The  
226 second scenario (PAL) considers reasonably well-known climate forcing over the last glacial-interglacial cycle.  
227 Simulations under this scenario are initialized with output from a previous spin-up simulation forced by glacial  
228 boundary conditions with respect to orbital parameters (Berger, 1978), greenhouse gas radiative forcing based on  
229 reconstructed atmospheric greenhouse gases (Köhler et al., 2017), and ice sheet extent. In simulations under PAL, the  
230 model is integrated until 0 kyr BP following the reconstructed histories of the former. Ice sheets for the preindustrial  
231 and Last Glacial Maximum (LGM) states are taken from Peltier (1994) and linearly scaled using the global benthic  
232  $\delta^{18}\text{O}$  stack of Lisiecki and Stern (2016), which is a global ice volume proxy. Changes in the albedo, salinity and latent  
233 heat flux associated with the ice sheet buildup or melting are also taken into account (Ritz et al., 2011). Model scenario  
234 PAL appears to still be missing an important process or feedback for atmospheric  $\text{CO}_2$ , as it cannot reproduce the  
235 observed low glacial  $\text{CO}_2$  level without invoking additional changes (see, e.g., Tschumi et al., 2011; Menviel et al.,  
236 2012; Roth and Joos, 2013; Jeltsch-Thömmes et al., 2019).

237

238 In this study, we consider six scenarios that invoke additional changes to force the model toward the observed  
239 low glacial  $\text{CO}_2$  concentration. In addition to the PAL forcing, a time-varying scale factor  $F(t)$  is applied to some  
240 combination of tunable model parameters: wind stress scale factor  $\tau$ , vertical diffusivity  $K_V$ , gas transfer velocity  $k_w$ ,  
241  $\text{CaCO}_3$ -to-POC export ratio  $rr$ , and POC remineralization length scale  $\ell_{POC}$ . For the preindustrial period, the value of  
242  $F(t)$  is fixed at 1, whereas the theoretical LGM value was chosen in order to achieve an atmospheric  $\text{CO}_2$  concentration  
243 close to the LGM level of  $\sim 190$  ppm (see Table 1), as determined by sensitivity experiments. Note that the same values  
244 of  $F(t)$  apply to any of the model parameters considered in a given scenario. To obtain intermediate values,  $F(t)$  is  
245 linearly scaled using the global benthic  $\delta^{18}\text{O}$  stack (see Fig. 1). For the spin-up needed to initialize these simulations,  
246 the glacial spin-up simulation of PAL was integrated for 50,000 model years, with tunable parameters adjusted to their  
247 appropriate glacial values. Atmospheric  $\text{CO}_2$  drawdown of up to  $\sim 100$  ppm is achieved over this 50,000-year  
248 integration. From that final spun-up state, the model is run forward in time until 0 kyr BP with PAL and  $F(t)$  forcing.

249

250 The first of these scenarios (CIRC) allows us to test the sensitivity of the model results with respect to changes  
251 in ocean circulation. Tunable model parameters  $\tau$  and  $K_V$  were reduced to 40 percent of their preindustrial values  
252 throughout the global ocean during the LGM (i.e.,  $F_{\tau, K_V} = 0.4$ ). Such a drastic change in the wind stress field is not  
253 realistic. Rather, these changes should be viewed as “tuning knobs” that force the ocean model into a poorly ventilated  
254 state with an “older” ideal age and  $^{14}\text{C}$ -depleted deep waters, as suggested for the glacial ocean (e.g., Sarin et al.,  
255 2013). In the model's implementation, a change in wind stress does not affect the gas transfer velocity  $k_w$ , unlike in  
256 the real ocean where changes in wind stress and wind speed act together. The influence of a change in air-sea exchange



257 efficiency on the model results was investigated in a second scenario (VENT) where  $k_w$  is reduced in the model's  
258 north ( $> 60^\circ\text{N}$ ) and south ( $> 48^\circ\text{S}$ ) polar areas in addition to global reductions of  $\tau$  and  $K_V$  ( $F_{\tau,K_V,k_w} = 0.4$ ). A 60  
259 percent reduction of  $k_w$  is unlikely to be correct but is a straightforward way to reduce the model's gas exchange  
260 efficiency. In the third scenario (VENTx), reduction of polar  $k_w$  to 0 percent of its preindustrial value was tested  
261 ( $F_{\tau,K_V} = 0.4$ ;  $F_{k_w} = 0.0$ ). Here,  $k_w$  remains fixed at 0 percent during the last glacial and is adjusted to its preindustrial  
262 value via a linear ramp across the last glacial termination ( $\sim 18$  to 11 kyr BP). In this scenario, sea ice would  
263 permanently cover 100 percent of the Southern Ocean during the last glacial, which is not supported by the sea ice  
264 reconstructions of Gersonde et al. (2005) and Allen et al. (2011), and also the high-latitude ( $> 60^\circ\text{N}$ ) North Atlantic  
265 and Arctic Ocean, for which there is some evidence (Müller and Stein, 2014; Hoff et al., 2016).

266

267 We end by investigating the sensitivity of the model results to changes in the parameters controlling the  
268 export production of  $\text{CaCO}_3$  and the water column remineralization of POC. Model scenario BIO considers changes  
269 of the  $\text{CaCO}_3$ -to-POC export ratio (and thus also the  $\text{CaCO}_3$ -to-POC rain ratio; Archer and Maier-Reimer, 1994)  
270 ( $F_{rr} = 0.8$ ) and POC remineralization length scale (Roth et al., 2014) ( $F_{\ell_{POC}} = 1.2$ ). These changes impact the global  
271 carbon cycle by influencing the vertical gradients of DIC, ALK, and nutrients in the water column. A change in the  
272 fluxes of POC and  $\text{CaCO}_3$  to the sea floor drives a change in the magnitude of POC and  $\text{CaCO}_3$  burial in the sediments.  
273 A modest reduction in the export ratio during the last glacial is compatible with reconstructed variations in carbonate  
274 ion concentrations (Jeltsch-Thömmes et al., 2019). How the depth of POC remineralization changed over time is still  
275 unknown. The last two scenarios consider the combined effect of physical and biogeochemical changes: PHYS-BIO  
276 ( $F_{\tau,K_V,k_w,rr} = 0.7$ ) and PHYS-BIOx ( $F_{\tau,K_V,k_w,rr} = 0.8$ ;  $F_{\ell_{POC}} = 1.2$ ).

277

## 278 2.5 Field- and model-based reconstruction of $^{14}\text{C}$ production

279

280 Our ability to attribute past changes in atmospheric  $\Delta^{14}\text{C}$  to climate-related changes in the ocean carbon cycle is limited  
281 by our ability to reconstruct a precise and accurate history of the  $^{14}\text{C}$  production rate. Past changes in  $^{14}\text{C}$  production  
282 can be estimated from geomagnetic field reconstructions and from  $^{10}\text{Be}$  measurements in polar ice cores. For ice-core  
283  $^{10}\text{Be}$ -based estimates, we use the ice-core radionuclide stack of Adolphi et al. (2018), which is based on  $^{36}\text{Cl}$  data from  
284 the GRIP ice core (Baumgartner et al., 1998), and on  $^{10}\text{Be}$  data from the GRIP (Yiou et al., 1997; Baumgartner et al.,  
285 1997; Wagner et al., 2001; Muscheler et al., 2004; Adolphi et al., 2014) and GISP2 (Finkel and Nishiizumi, 1997) ice  
286 cores. It also includes  $^{10}\text{Be}$  data from the NGRIP, EDML, EDC, and Vostok ice cores around the Laschamp  
287 geomagnetic excursion (Raisbeck et al., 2017). It has been extended to the present using the  $^{10}\text{Be}$  stack of Muscheler  
288 et al. (2016). This 70,000-year  $^{10}\text{Be}$  stack provides relative changes of  $^{14}\text{C}$  production rates under the assumption that  
289  $^{14}\text{C}$  and  $^{10}\text{Be}$  production rates are directly proportional, as indicated by the most recent production rate models (e.g.,  
290 Herbst et al., 2017).

291

292 For paleointensity-based estimates, we employ (1) the North Atlantic Paleointensity Stack, or NAPIS, by Laj  
293 et al. (2000) as extended by Laj et al. (2002), (2) the Global Paleointensity Stack, or GLOPIS, by Laj et al. (2004), (3)





294 a high-resolution paleointensity stack from the Black Sea (Nowaczyk et al., 2013), and (4) a paleointensity stack from  
295 Iberian Margin sediments (Channell et al., 2018). In principle, stacks of widely distributed cores (NAPIS/GLOPIS)  
296 are expected to yield a better representation of the global geomagnetic dipole moment, whereas the paleointensity  
297 stacks from the Black Sea and the Iberian Margin avoid some of the problems associated with coring disturbances.  
298 The four different paleointensity stacks were converted to  $^{14}\text{C}$  production rates using the production rate model of  
299 Herbst et al. (2017), the local interstellar spectrum of Potgieter et al. (2014), and assuming a constant solar modulation  
300 potential of 630 MeV.

301

302 An alternative approach to estimating the  $^{14}\text{C}$  production rate is to combine an atmospheric radiocarbon  
303 budget with a prognostic carbon cycle model. Here simulations are performed with the Bern3D model and forced by  
304 reconstructed changes in atmospheric  $\Delta^{14}\text{C}$  and  $\text{CO}_2$  over the last 50 kyr. Both the IntCal13 calibration curve (Reimer  
305 et al., 2013) and the Hulu Cave  $\Delta^{14}\text{C}$  dataset (Cheng et al., 2018) are used. The  $^{14}\text{C}$  production rate  $Q$  is calculated,  
306 each model year, from the air-sea  $^{14}\text{CO}_2$  flux ( $F_{as}$ ), the atmosphere-land  $^{14}\text{CO}_2$  flux ( $F_{ab}$ ), the loss of  $^{14}\text{C}$  due to  
307 radioactive decay, and the change ( $\dot{I}_a$ ) in the atmospheric  $^{14}\text{C}$  inventory ( $I_a$ ):

308

$$309 \quad Q = F_{as} + F_{ab} + \lambda I_a + \dot{I}_a \quad (2)$$

310

311 where  $\lambda$  is the radioactive decay constant for  $^{14}\text{C}$ , i.e.,  $\lambda = \ln 2 / 5700 \text{ years} = 1.2160 \times 10^{-4} \text{ yr}^{-1}$ . The radioactive decay  
312 term  $\lambda I_a$  and the change in inventory  $\dot{I}_a$  follow the reconstructed  $\Delta^{14}\text{C}$  and  $\text{CO}_2$  records, whereas  $F_{as}$  and  $F_{ab}$  are  
313 explicitly computed by the model. The  $F_{as}$  term depends strongly on the carbon cycle scenario under consideration  
314 (see Sect. 2.4 and Table 1). For comparison with other reconstructions,  $Q$  is converted into a relative value by  
315 normalizing it by the preindustrial value.

316

### 317 **3 Results and discussion**

318

#### 319 **3.1 $\Delta^{14}\text{C}$ response to step changes**

320

321 We use step changes in the  $^{14}\text{C}$  production rate, and in selected carbon cycle parameters, to gain insight into the  
322 characteristic magnitude and time scale of the corresponding  $\Delta^{14}\text{C}$  changes (Fig. 4). Besides variations of the  
323 production rate, changes in ocean circulation and air-sea gas exchange are considered the most important factors  
324 affecting atmospheric  $\Delta^{14}\text{C}$ . Their effect on  $\Delta^{14}\text{C}$  can be understood in terms of their effect on the reservoir sizes  
325 involved in the global carbon cycle and on the exchange rates between the reservoirs. We investigate the relative  
326 importance of the major global carbon reservoirs (atmosphere, terrestrial biosphere, ocean, and sediments) by  
327 considering four different model configurations (see Sect. 2.3), with particular emphasis on the role of marine  
328 sediments.

329



330 In model studies, the process of sedimentation (here used specifically to refer to the balance between  
331 deposition and remineralization/dissolution at the sediment-water interface) is often neglected because it is a relatively  
332 minor flux. In the Bern3D model, sedimentation removes only about 0.46 Gt C and 45.31 mol  $^{14}\text{C}$  per year in the  
333 preindustrial steady state. Indeed, interaction with the sediments has little influence on the global mean value of  
334 oceanic  $\Delta^{14}\text{C}$ , and therefore atmospheric  $\Delta^{14}\text{C}$ , as long as the total oceanic amount of carbon remains approximately  
335 constant (Siegenthaler et al., 1980); however, this is not always true, particularly in the case of millennial-scale climate  
336 perturbations. The differences between the two sets of model runs shown in Fig. 4 (i.e., ALL versus OCN-LND, and  
337 OCN-SED versus OCN) are due to sedimentation-driven changes in the ocean carbon inventory. In order to facilitate  
338 our discussion, we will only make direct comparisons between model runs ALL and OCN-LND. We note that the  $^{14}\text{C}$   
339 exchange rate between the atmosphere and the terrestrial biosphere is only of minor importance for long time scales  
340 of millennia and more.

341

### 342 3.1.1 Change of $^{14}\text{C}$ production

343

344 At steady state, the relative change of atmospheric  $\Delta^{14}\text{C}$  is equal to the relative change of the  $^{14}\text{C}$  production rate,  
345 irrespective of the individual reservoirs considered. Fig. 4 shows that atmospheric  $\Delta^{14}\text{C}$  increases by about 100 permil  
346 (or 10 percent) when the production rate is increased by 10 percent. In model run ALL,  $\Delta^{14}\text{C}$  increases approximately  
347 exponentially to its new steady-state value with a characteristic time constant  $T$  of about 6170 years (i.e.,  $1 - 1/e \approx$   
348 63 percent of the total change in  $\Delta^{14}\text{C}$  occurs within 6170 years). This e-folding time scale is close to the mean lifetime  
349 of  $^{14}\text{C}$  (~8223 years), which is modulated by the time required for equilibration between the atmosphere and the ocean  
350 (i.e., the time scale for deep ocean ventilation, of the order of hundreds to 1000 years). In the next section, we will  
351 investigate the effect of ocean carbon cycle processes on atmospheric  $\Delta^{14}\text{C}$ .

352

353 Note that for simplicity, we investigated only step changes in atmospheric production, although, in reality,  
354  $^{14}\text{C}$  production varies continuously over time due to changes in the solar and/or geomagnetic modulation of the cosmic  
355 radiation. This results in a non-steady state value of atmospheric  $\Delta^{14}\text{C}$ .

356

### 357 3.1.2 Change of ocean circulation

358

359 The exchange rate between the surface and deep ocean is mainly determined by physical transport and mixing  
360 processes. The overall effect of these processes is to transport  $^{14}\text{C}$ -enriched surface waters to the thermocline and deep  
361 ocean, where waters are typically  $^{14}\text{C}$ -depleted. In addition, the nutrient supply by transport and mixing plays an  
362 important role in determining the production and export of biogenic material from the surface ocean, constituting a  
363 second pathway for transporting  $^{14}\text{C}$  to the deep ocean.

364

365 In the Bern3D model, the tunable model parameters affecting the ventilation of the deep ocean include a  
366 scale factor  $\tau$  for the wind stress field and vertical diffusivity  $K_V$ . Fig. 4 shows the atmospheric  $\Delta^{14}\text{C}$  response after a



367 sudden decrease of  $\tau$  and  $K_V$  by 50 percent. Although a halving of  $\tau$  and  $K_V$  does not represent a realistic change, the  
368 resulting state of the ocean's large-scale overturning circulation can be interpreted in terms of the "ideal age" of water,  
369 which represents the average time since a water mass last made surface boundary contact. The new steady-state ideal  
370 age after a halving of  $\tau$  and  $K_V$  is almost three times greater than the preindustrial steady-state value (i.e., ~1664 years  
371 versus ~613 years). This "ageing" of the ocean is achieved through a weakening and shoaling of the global meridional  
372 overturning circulation as evident from a moderate reduction in the meridional overturning stream function for the  
373 Indo-Pacific Ocean from about 14 to 9.5 Sv (1 Sv =  $10^6 \text{ m}^3 \text{ s}^{-1}$ ), and a very strong reduction from about 18 to 8 Sv in  
374 the Atlantic meridional overturning stream function, consistent with evidence for the glacial ocean. Here, as expected,  
375 the overall effect of deep water ageing is a stronger vertical  $\Delta^{14}\text{C}$  gradient in the water column and a subsequent  
376 increase in atmospheric  $\Delta^{14}\text{C}$ . The exact nature of the  $\Delta^{14}\text{C}$  response, however, depends on the carbon reservoirs  
377 considered.

378

379 If sedimentation is neglected, the time required for  $\Delta^{14}\text{C}$  to adjust to step changes in  $\tau$  and  $K_V$  is relatively  
380 short. Atmospheric  $\Delta^{14}\text{C}$  increases rapidly to its new steady-state value of ~159 permil, with a time constant  $T$  of about  
381 600 years. This increase of  $\Delta^{14}\text{C}$  can be explained by the fact that, owing to a weaker and shallower overturning  
382 circulation, a comparatively large amount of carbon is moved from the atmosphere to the ocean. More specifically,  
383 the atmospheric carbon inventory decreases by 14.6 percent, whereas the atmospheric  $^{14}\text{C}$  inventory decreases by only  
384 1.1 percent (Fig. 5c). The  $^{14}\text{C}$  being produced in the atmosphere is therefore diluted by a smaller carbon inventory,  
385 increasing the atmospheric  $^{14}\text{C}/\text{C}$  ratio; this asymmetry in the drawdown of  $\text{CO}_2$  and  $^{14}\text{CO}_2$  is what permits the increase  
386 of atmospheric  $\Delta^{14}\text{C}$ . Since the ocean carbon inventory changes by only +0.2 percent, the  $\Delta^{14}\text{C}$  value for the global  
387 ocean is nearly unaffected, a decrease of only ~11 permil in the new steady state (Fig. 6g).

388

389 In the model run with sedimentation, there are two distinct time constants. A rapid increase of atmospheric  
390  $\Delta^{14}\text{C}$  occurs, ~143 permil in the first few hundred years, then  $\Delta^{14}\text{C}$  gradually decreases to its final value of ~91 permil  
391 after tens of thousands of years. Reduced deep ocean ventilation is again responsible for the rapid  $\Delta^{14}\text{C}$  change and  
392 the respective time constant ( $T = \sim 480$  years). The second time constant of ~23,390 years is due to the relatively long  
393 time required for adjustment of the ocean carbon inventory to ocean circulation-driven changes in sedimentation.

394

395 The process of ocean circulation interacts with the efficiency of the ocean's biological carbon pump, via its  
396 impact on export production, ocean interior oxygen levels, and seawater carbonate chemistry/equilibria. This has  
397 important implications for the sedimentation of biogenic material at the sea floor and, on a time scale of tens of  
398 thousands of years, the total oceanic amount of carbon. Through this coupling of ocean circulation and sedimentation  
399 via the biological carbon pump, a halving of  $\tau$  and  $K_V$  leads to a 9.8 percent increase of the ocean carbon inventory in  
400 the new steady state (Fig. 5e). Qualitatively, a reduction in the ocean's overturning circulation leads to a lower surface  
401 nutrient supply, which limits the production and export of biogenic material from the surface ocean. This, in turn,  
402 decreases the fluxes of POC and  $\text{CaCO}_3$  to the sea floor, with major consequences for the magnitude of their  
403 sedimentation. At the same time, a constant amount of DIC, ALK, and nutrients is added to the ocean from terrestrial



404 weathering which is no longer balanced by the sedimentation flux (this is what permits a larger ocean carbon  
405 inventory). The overall effect is a gradual reduction of global-ocean  $\Delta^{14}\text{C}$  by  $\sim 76$  permil (Fig. 6g), which dilutes the  
406 initial atmospheric  $\Delta^{14}\text{C}$  peak by 52 permil.

407

### 408 3.1.3 Change of gas transfer velocity

409

410 It takes about a decade for the isotopic ratios of carbon to equilibrate between the atmosphere and a  $\sim 75$ -m thick  
411 surface mixed layer by air-sea gas exchange alone (Broecker and Peng, 1974). A consequence of this is that the surface  
412 ocean is undersaturated with respect to atmospheric  $\Delta^{14}\text{C}$  (see Fig. 3). The choice of gas transfer velocity  $k_w$  as a  
413 function of wind speed is critical for the efficiency of air-sea gas exchange. A reduction of  $k_w$  corresponds to a higher  
414 resistance for gas transfer across the air-sea interface, which means that the  $^{14}\text{C}$  produced in the atmosphere escapes  
415 into the surface ocean at a slower rate. The effect of a lower  $k_w$  is a larger air-sea gradient of  $\Delta^{14}\text{C}$  and higher  
416 atmospheric  $\Delta^{14}\text{C}$  values. In contrast, the  $\Delta^{14}\text{C}$  value for the surface ocean is nearly unaffected so long as the ocean  
417 carbon inventory remains approximately constant, since the vertical gradient of  $\Delta^{14}\text{C}$  in the ocean is dominated by  
418 physical transport and mixing processes. Although the exact nature of the gas transfer velocity under glacial climate  
419 conditions remains unclear,  $k_w$  represents a straightforward way to reduce the model's air-sea exchange efficiency  
420 due to theoretical changes in wind stress, sea ice, etc.

421

422 Fig. 4 shows how atmospheric  $\Delta^{14}\text{C}$  responds to a perturbation in the gas transfer velocity. In the model run  
423 without sedimentation, reduction of  $k_w$  to 0 percent of its preindustrial value, in the model's north ( $> 60^\circ\text{N}$ ) and south  
424 ( $> 48^\circ\text{S}$ ) polar areas, leads to a moderate increase of atmospheric  $\Delta^{14}\text{C}$  in the new steady state. The amplitude of  $\Delta^{14}\text{C}$   
425 change is  $\sim 42$  permil, which is achieved with an e-folding time scale  $T$  of about 180 years. This relatively short time  
426 constant can be explained by the multidecadal time scale required for  $\Delta^{14}\text{C}$  to equilibrate between the model's  
427 atmosphere, upper ocean, and terrestrial biosphere. As shown in Fig. 6, the global mean  $\Delta^{14}\text{C}$  values for the surface,  
428 deep, and global ocean in the new steady state are only slightly different from the preindustrial steady-state values, as  
429 expected from the fact that the ocean carbon inventory remains relatively stable.

430

431 Interestingly, if sedimentation is included in the model, the final value of atmospheric  $\Delta^{14}\text{C}$  is much higher  
432 ( $\sim 91$  permil). In this case, a perturbation in  $k_w$  leads to a very rapid initial increase of  $\Delta^{14}\text{C}$  ( $\sim 42$  permil), and a much  
433 slower subsequent increase of  $\Delta^{14}\text{C}$  ( $\sim 49$  permil). The latter has an e-folding time scale  $T$  of about 14,200 years. This  
434 slow doubling of the initial  $\Delta^{14}\text{C}$  increase is unexpected, but can be explained by the fact that a reduction of  $k_w$   
435 involves also a reduction of air-sea  $\text{O}_2$  gas exchange in the deep water formation regions, decreasing the oceanic  
436 oxygen that is available for transport to the deep ocean. This, in turn, implies lower oxygen concentrations in the water  
437 column and the sediment pore waters, decreasing the rate of POC remineralization in the sediments. Reducing this has  
438 the overall effect of strengthening POC sedimentation at the sea floor, causing the ocean carbon inventory to decrease.  
439 As shown in Fig. 5, the total oceanic amount of carbon decreases by 5.9 percent in the new steady state, resulting in  
440 elevated  $\Delta^{14}\text{C}$  values for the surface (+56 permil), deep (+30 permil), and global (+37 permil) ocean as well as for the



441 atmosphere (+91 permil) (see Fig. 6). Note that the increase in atmospheric  $\Delta^{14}\text{C}$  is not accompanied by a significant  
442 change in the atmospheric carbon ( $\text{CO}_2$ ) inventory (i.e., it decreases by only 2.2 to 3.3 percent).

443

444 Overall, findings from these sensitivity experiments demonstrate that (1) the response of atmospheric  $\Delta^{14}\text{C}$   
445 to changes in the internal parameters of the ocean carbon cycle, in contrast to  $^{14}\text{C}$  production changes, depends strongly  
446 on whether or not sedimentation at the sea floor is simulated, (2) the e-folding time scale for the initial adjustment of  
447 atmospheric  $\Delta^{14}\text{C}$  to ocean carbon cycle changes, i.e., changes in ocean circulation and gas exchange, is shorter than  
448 that for production changes (i.e., ~600 years and ~180 years versus ~6170 years), (3) air-sea gas exchange, in contrast  
449 to ocean circulation, has only a small effect on atmospheric  $\text{CO}_2$ , and (4) on time scales of tens of thousands of years  
450 changes in the sedimentation of biogenic material can potentially diminish (or elevate) the atmospheric  $\Delta^{14}\text{C}$  value.  
451 This is new, important information for future paleoclimate simulations, indicating that changes in  $\Delta^{14}\text{C}$  may be  
452 overestimated (or underestimated) in models that do not simulate the interaction between sediments and the water  
453 column.

454

### 455 **3.2 Role of $^{14}\text{C}$ production in past $\Delta^{14}\text{C}$ variability**

456

457 We now consider the component of past atmospheric  $\Delta^{14}\text{C}$  variability caused by production changes alone. Fig. 7  
458 shows the results of model runs using different reconstructions of the  $^{14}\text{C}$  production rate, as inferred from  
459 paleointensity data and from ice-core  $^{10}\text{Be}$  fluxes. The global carbon cycle is assumed to be constant and under  
460 preindustrial conditions for these simulations (i.e., scenario MOD is used). Our analysis is restricted to the glacial  
461 portion of the record (50 to 18 kyr BP), in part because this is the time period which experiences the largest production  
462 changes, and in part because we did not attempt to reproduce the ~80 ppm change in atmospheric  $\text{CO}_2$  that occurred  
463 during the last glacial termination. As we have already noted, much research over the last decades has attempted to  
464 explain the observed glacial-interglacial variations in atmospheric  $\Delta^{14}\text{C}$  and  $\text{CO}_2$ , and this was not the goal of this  
465 study.

466

467 At first glance, the millennial-scale structure of model-simulated  $\Delta^{14}\text{C}$  is comparable to that of the  
468 reconstructions. These similarities appear to be highest for the oldest portion of the record, roughly before 30 kyr BP.  
469 The model reproduces major features of the reconstructed  $\Delta^{14}\text{C}$  variability such as the large changes associated with  
470 the Laschamp (~41 kyr BP) and Mono Lake (~34 kyr BP) geomagnetic excursions. These two events are clearly  
471 expressed as distinct maxima in all model-simulated records. A more detailed comparison reveals a high correlation  
472 between the modelled and reconstructed  $\Delta^{14}\text{C}$  values between 50 and 33 kyr BP. Of note is the better agreement with  
473 the new Hulu Cave  $\Delta^{14}\text{C}$  dataset as compared to the IntCal13 calibration curve (i.e., Pearson correlation coefficient  $r$   
474 of 0.96 versus 0.91). This is likely due to the fact that the Laschamp excursion is smoothed/smeared out during the  
475 stacking process of the IntCal13  $\Delta^{14}\text{C}$  datasets (Adolphi et al., 2018). The correlation between modelled and  
476 reconstructed  $\Delta^{14}\text{C}$  is much weaker during the millennia after the Mono Lake excursion (33 to 18 kyr BP;  $r = 0.52$  to  
477 0.64). While it is clear that much of the millennial-scale variation in  $\Delta^{14}\text{C}$  is driven by past changes in  $^{14}\text{C}$  production,



478 the model fails to reproduce the glacial level of  $\Delta^{14}\text{C}$  and also does not capture the ~15,000-year persistent elevation  
479 of  $\Delta^{14}\text{C}$  or the subsequent decrease of  $\Delta^{14}\text{C}$  after ~25 kyr BP.

480

481 The reconstructions suggest that the highest values of  $\Delta^{14}\text{C}$  occurred during the Laschamp excursion, with a  
482 maximum value of ~595 permil at 41.1 kyr BP found in the IntCal13 record. The Hulu Cave record indicates even  
483 higher values for the Laschamp event ( $\Delta^{14}\text{C} = \sim 742$  permil, at 39.7 kyr BP). In contrast, the model is able to simulate  
484 maximum  $\Delta^{14}\text{C}$  values of only ~364 permil at 40.4 kyr BP, and ~236 permil at 40.5 kyr BP, as predicted by the  
485 paleointensity-based and ice-core  $^{10}\text{Be}$ -based production rate estimates, respectively. Although the model is unable to  
486 reproduce the reconstructed values of  $\Delta^{14}\text{C}$ , the modelled amplitude of the variation in  $\Delta^{14}\text{C}$  in response to the  
487 Laschamp event shows a reasonable agreement with the reconstructed amplitude of  $\Delta^{14}\text{C}$  change found in the IntCal13  
488 record (~240 permil). The  $\Delta^{14}\text{C}$  change predicted by paleointensity data has a maximal amplitude of about 320 permil,  
489 whereas the ice-core  $^{10}\text{Be}$  data indicate a smaller amplitude (~224 permil). Note that the IntCal13 and model-simulated  
490 amplitudes of the Laschamp-related  $\Delta^{14}\text{C}$  change are about two times smaller than that observed in the Hulu Cave  
491 record (~575 permil), which is more likely to be correct.

492

493 Moving onto the full glacial record (50 to 18 kyr BP), there are considerable discrepancies between  
494 reconstructed and modelled  $\Delta^{14}\text{C}$  ( $\Delta\Delta^{14}\text{C}$ ; see Fig. 7). The use of ice-core  $^{10}\text{Be}$  data to predict past changes in  $\Delta^{14}\text{C}$   
495 results in the largest  $\Delta\Delta^{14}\text{C}$ , with offsets between the records as high as ~544 to 558 permil (root-mean-square error  
496  $RMSE = 404$  to 408 permil). Model-simulated  $\Delta^{14}\text{C}$  given by paleointensity data varies widely between the four  
497 available reconstructions, yielding  $\Delta\Delta^{14}\text{C}$  values of ~325 to 639 permil ( $RMSE = 206$  to 455 permil). Note that the  
498 upper limit of the paleointensity-based  $\Delta\Delta^{14}\text{C}$  overlaps with the ice-core  $^{10}\text{Be}$ -based  $\Delta\Delta^{14}\text{C}$ . Given the uncertainties  
499 associated with the reconstruction of past changes in  $^{14}\text{C}$  production, accurate predictions of its contribution to past  
500 changes in  $\Delta^{14}\text{C}$  are challenging. Nonetheless, the substantial systematic offsets between the model-simulated and  
501 reconstructed  $\Delta^{14}\text{C}$  records after ~33 kyr BP point toward insufficiently high  $^{14}\text{C}$  production rates over this period of  
502 time. The question arises as to whether another factor besides geomagnetic modulation of the cosmic ray intensity  
503 was responsible for elevated glacial  $\Delta^{14}\text{C}$  levels. The effect of ocean carbon cycle changes on the evolution of  
504 atmospheric  $\Delta^{14}\text{C}$  is considered next.

505

### 506 **3.3 Carbon cycle contribution to high glacial $\Delta^{14}\text{C}$ levels**

507

508 Here we investigate the magnitude and timing of the maximum possible  $\Delta^{14}\text{C}$  change during the last glacial period,  
509 obtained by running the Bern3D model with eight different carbon cycle scenarios (see Table 1). For the sake of  
510 clarity, we will discuss only the results of model runs using the mean paleointensity-based  $^{14}\text{C}$  production rate, though  
511 all available reconstructions were used. We emphasize that this is not a best-guess estimate of paleointensity-based  
512  $^{14}\text{C}$  production. One should focus on the relative changes of  $\Delta^{14}\text{C}$  between model scenarios, and how specific carbon  
513 cycle processes affect the glacial level of  $\Delta^{14}\text{C}$ .

514



515 Modelled 50,000-year records of atmospheric  $\Delta^{14}\text{C}$  and  $\text{CO}_2$  plus corresponding reconstructed histories are  
516 shown in Fig. 8. In order to provide a basis for comparison of modelling efforts, the results of model run MOD (which  
517 assumes a constant preindustrial carbon cycle) are presented. The influence of ocean carbon cycle changes on  
518 atmospheric  $\Delta^{14}\text{C}$  was tested in the other model runs. Interestingly, the forcing fields for model run PAL (orbital  
519 parameters, greenhouse gas radiative forcing, and ice sheet extent) have only a minimal impact on  $\Delta^{14}\text{C}$ . The PAL  
520 forcing fields also do not achieve sufficiently low glacial  $\text{CO}_2$  concentrations. Only a slight reduction of atmospheric  
521  $\text{CO}_2$  by  $\sim 20$  ppm could be achieved, which unrealistically occurs during the last glacial termination ( $\text{CO}_2 = 258.07$   
522 ppm, at 14.6 kyr BP). With hypothesized carbon cycle changes, the agreement between observed and modelled  $\text{CO}_2$   
523 during the last glacial period is good (as by design), but the deglacial  $\text{CO}_2$  rise is lagged and  $\sim 60$  ppm too small at 11  
524 kyr BP. Since this study focuses on glacial  $\Delta^{14}\text{C}$  levels before incipient deglaciation at  $\sim 18$  kyr BP, we will not discuss  
525 the lag any further.

526

527 Running the model with additional carbon cycle changes leads to an improvement of modelled  $\Delta^{14}\text{C}$  as  
528 compared to model run PAL. The amplitude of atmospheric  $\Delta^{14}\text{C}$  change is highest for runs CIRC, VENT, and  
529 VENTx. This behavior is due to the fact that, owing to a reduction of  $\tau$ ,  $K_V$ , and  $k_w$ , strong vertical  $\Delta^{14}\text{C}$  gradients in  
530 the ocean, as well as a large air-sea  $\Delta^{14}\text{C}$  gradient, are established. As shown in Fig. 8, the reduction in ocean  
531 ventilation is clearly expressed in the observed increase of the model ocean's ideal age as well as surface- and deep-  
532 water reservoir ages, the latter two being equivalent to radiocarbon reservoir age offsets following Soulet et al. (2016).  
533 The deep-water reservoir age (i.e., B-Atm age offset) provides a measure of the radiocarbon disequilibrium between  
534 the deep ocean and the atmosphere, which arises due to the combined effect of air-sea gas exchange efficiency and  
535 the strength of the ocean's overturning circulation (i.e., the deep ocean ventilation rate). As a consequence of the  
536 reduction in the meridional overturning circulation, model run CIRC predicts a substantial increase in B-Atm for the  
537 glacial ocean, with values reaching as high as  $\sim 3236$   $^{14}\text{C}$  years, nearly double the preindustrial value of  $\sim 1626$   $^{14}\text{C}$   
538 years. Model run VENT predicts a slightly larger increase ( $+166$   $^{14}\text{C}$  years) in the glacial ocean's B-Atm, due to the  
539 inhibition of air-sea gas exchange. The "oldest" waters are found in run VENTx where air-sea gas exchange is severely  
540 restricted, giving a glacial B-Atm value of  $\sim 3576$   $^{14}\text{C}$  years. This would imply that the model's glacial ocean was more  
541 than two times older than its preindustrial counterpart. Further support comes from the model-simulated ideal age,  
542 which indicates that the ventilation time scale for the glacial ocean was about two times longer than for the  
543 preindustrial ocean (i.e.,  $\sim 1297$  years versus  $\sim 613$  years).

544

545 Indirect evidence for deep water ageing can be provided by the occurrence of depleted ocean interior oxygen  
546 levels, due to the progressive consumption of dissolved oxygen during POC remineralization. This situation would  
547 signal higher values of apparent oxygen utilization ( $\text{AOU} = \text{O}_{2,\text{pre}} - \text{O}_2$ ). Model runs CIRC, VENT, and VENTx do  
548 indeed indicate a large increase in AOI of about  $95 \text{ mmol m}^{-3}$  from its preindustrial value of  $\sim 150 \text{ mmol m}^{-3}$ . The  
549 reason for this AOI increase is that a reduction of deep ocean ventilation permits enhanced accumulation of  
550 remineralized carbon in the ocean interior and therefore a more efficient biological carbon pump. Model runs BIO,  
551 PHYS-BIO, and PHYS-BIOx allow us to investigate the impact of other biological carbon pump changes on



552 atmospheric  $\Delta^{14}\text{C}$  and  $\text{CO}_2$  (i.e., changes in the  $\text{CaCO}_3$ -to-POC export ratio and POC remineralization length scale).  
553 While these changes lead to an effective atmospheric  $\text{CO}_2$  drawdown mechanism, model results confirm that their  
554 effect on atmospheric  $\Delta^{14}\text{C}$  is much less important (see Fig. 8).

555

556 Model run VENTx gives the best results with respect to glacial levels of  $\Delta^{14}\text{C}$ , with a maximum  
557 underestimation of  $\sim 202$  to  $229$  permil ( $RMSE = 103$  to  $110$  permil) and a relatively good correlation ( $r = 0.79$  to  
558  $0.91$ ). Only one model parameter was changed for run VENTx as compared to runs CIRC and VENT, namely, the  
559 polar gas transfer velocity  $k_w$  was reduced to 0 percent of its preindustrial value during the last glacial. In this extreme  
560 scenario, we assume that sea ice cover extended in the northern hemisphere as far south as  $60^\circ\text{N}$  and in the southern  
561 hemisphere as far north as  $48^\circ\text{S}$ , which is not supported by the reconstructions (Gersonde et al., 2005; Allen et al.,  
562 2011). Nonetheless, considering extreme assumptions about polar air-sea exchange efficiency under glacial climate  
563 conditions is interesting for two reasons: (1) a change in gas exchange hardly affects the atmospheric  $\text{CO}_2$   
564 concentration, and (2) an additional change of atmospheric  $\Delta^{14}\text{C}$  could be achieved on a time scale of tens of thousands  
565 of years through a change in sedimentation at the sea floor. This behavior has important implications for the glacial  
566 atmosphere, which is characterized by high  $\Delta^{14}\text{C}$  levels in conjunction with low but relatively stable  $\text{CO}_2$   
567 concentrations. In contrast to a change in ocean circulation, air-sea gas exchange is a dedicated  $\Delta^{14}\text{C}$  control knob that  
568 can be invoked for a further increase of atmospheric  $\Delta^{14}\text{C}$  without changing atmospheric  $\text{CO}_2$ . Here, an additional  
569 increase of  $\Delta^{14}\text{C}$  by  $\sim 130$  permil, as compared to runs CIRC and VENT, is achieved if gas exchange is reduced  
570 permanently to 0 percent in the polar regions.

571

572 While the modelled  $\Delta^{14}\text{C}$  values for run VENTx show rather good agreement with the reconstructions  
573 between 50 and 33 kyr BP ( $r = 0.92$  to  $0.96$ ;  $RMSE = 74$  to  $102$  permil), considerable discrepancies remain for the  
574 younger portion of the record. The analysis shown in Fig. 9 illustrates that even with extreme changes in the ocean  
575 carbon cycle it is very difficult to reproduce the reconstructed  $\Delta^{14}\text{C}$  values after  $\sim 33$  kyr BP. During this period of  
576 time, model run VENTx underestimates  $\Delta^{14}\text{C}$  by up to  $\sim 203$  permil ( $RMSE = 118$  to  $128$  permil), and is very poorly  
577 ( $r = 0.1$ ) correlated with the reconstructions, confirming that there are still considerable gaps in our understanding.  
578 Although it may be possible that permanent North Atlantic-Arctic and Antarctic sea ice cover extended to lower and  
579 higher latitudes than previously reconstructed, we conclude from our model study that even extreme assumptions  
580 about sea ice cover are insufficient to explain the elevated  $\Delta^{14}\text{C}$  levels after  $\sim 33$  kyr BP. It appears instead that the  
581 glacial  $^{14}\text{C}$  production rate was higher than previously estimated and/or the reconstruction of glacial atmospheric  $\Delta^{14}\text{C}$   
582 levels is biased high. The older portion of the  $\Delta^{14}\text{C}$  record is based on data from archives other than tree rings (i.e.,  
583 plant macrofossils, speleothems, corals, and foraminifera) (Reimer et al., 2013), providing, except for the Lake  
584 Suigetsu plant macrofossil data (Bronk Ramsey et al., 2012), only indirect measurements of atmospheric  $\Delta^{14}\text{C}$ . Note  
585 that these data show uncertainty in calendar age that propagate into the estimation of past atmospheric  $\Delta^{14}\text{C}$  levels.

586

587 Large uncertainties in the pre-Holocene  $^{14}\text{C}$  production rate also hamper our qualitative and quantitative  
588 interpretation of the atmospheric  $\Delta^{14}\text{C}$  record. There is considerable disagreement between the available





589 reconstructions of past changes in  $^{14}\text{C}$  production (Fig. 1). Paleointensity-based estimates typically predict higher  $^{14}\text{C}$   
590 production rates than ice-core  $^{10}\text{Be}$ -based ones. An exception is the paleointensity stack from Channell et al. (2018),  
591 which predicts lower production rates. But, irrespective of the scatter, it is clear that all of the  $^{14}\text{C}$  production rate  
592 estimates are insufficiently high to explain the elevated  $\Delta^{14}\text{C}$  levels during the last glacial. Given the uncertainties in  
593 these estimates, it is very difficult to quantitatively describe the role of the ocean carbon cycle in determining the  $\Delta^{14}\text{C}$   
594 and  $\text{CO}_2$  levels in the glacial atmosphere.

595

### 596 **3.4 Reconstructing the $^{14}\text{C}$ production rate by deconvolving the $\Delta^{14}\text{C}$ record**

597

598 The unresolved discrepancy between reconstructed and model-simulated  $\Delta^{14}\text{C}$  raises the question how the  $^{14}\text{C}$   
599 production rate would have had to evolve to be consistent with the IntCal13 calibration curve or the new Hulu Cave  
600  $\Delta^{14}\text{C}$  dataset. This question is addressed by deconvolving the atmospheric  $\Delta^{14}\text{C}$  reconstruction over the last 50 kyr,  
601 using the Bern3D carbon cycle model forced with reconstructed atmospheric histories of  $\Delta^{14}\text{C}$  and  $\text{CO}_2$  (see Eq. [2]).  
602 The carbon cycle scenarios described in Table 1, with the exception of MOD, are used in order to provide an estimate  
603 of the uncertainty associated with the model's glacial ocean carbon cycle. We note that the carbon cycle scenarios are  
604 not designed to capture the specific features of the last glacial termination, and therefore the results of the  
605 deconvolution over this time period must be considered very preliminary (and viewed tentatively). A detailed analysis  
606 of the Holocene  $^{14}\text{C}$  production rate is available in the literature (Roth and Joos, 2013). Finally, we consider the  
607 uncertainties associated with the older portion of the  $\Delta^{14}\text{C}$  record by deconvolving both the IntCal13 and Hulu Cave  
608  $\Delta^{14}\text{C}$  records. Hulu Cave data overlap with IntCal13 between  $\sim 10.6$  and  $33.3$  kyr BP (Cheng et al., 2018), as expected  
609 from the fact that IntCal13 between  $10.6$  and  $26.8$  kyr BP is based in part on Hulu Cave stalagmite H82 (Southon et  
610 al., 2012), whereas there are substantial offsets before  $\sim 30$  kyr BP.

611

612 Fig. 10 shows the new, model-based reconstruction of past changes in  $^{14}\text{C}$  production. Before the onset of  
613 the Laschamp excursion at  $\sim 42$  kyr BP, production rates as inferred from the Hulu Cave record are near modern levels,  
614 whereas those obtained from the IntCal13 record are somewhat higher than modern. As expected, peak production  
615 occurs during the Laschamp event ( $\sim 42$  to  $40$  kyr BP), with the Hulu Cave dataset yielding the largest amplitude  
616 (factor of  $\sim 2$  greater than modern). The IntCal13 record predicts a smaller amplitude of  $\sim 1.6$  times the modern value.  
617 Both  $\Delta^{14}\text{C}$  records predict production minima at  $\sim 37$  kyr BP ( $\sim 7$  percent higher than modern) and  $\sim 32$  kyr BP ( $\sim 5$   
618 percent higher than modern), interrupted by a prominent peak (factors of  $\sim 1.5$  and  $\sim 1.4$ , respectively) during the Mono  
619 Lake geomagnetic excursion ( $\sim 34$  kyr BP), though the details of the timing and structure differ between the two  
620 records. Between  $32$  and  $22$  kyr BP, model-based estimates of the  $^{14}\text{C}$  production rate are  $\sim 1.3$  times the modern value,  
621 which then decrease to around modern levels by HS1 ( $\sim 18$  kyr BP).

622

623 Model-based estimates of  $^{14}\text{C}$  production during the last glacial are typically higher than paleointensity-based  
624 and ice-core  $^{10}\text{Be}$ -based ones, as expected from the analysis in Sect. 3.2. Between  $32$  and  $22$  kyr BP, the deconvolutions  
625 of the IntCal13 and Hulu Cave  $\Delta^{14}\text{C}$  records give estimates that are about  $17.5$  percent higher than the reconstructions.



626 It is important to note that the differences between the production rate estimates inferred from the proxy data (i.e.,  
627 paleointensity data and ice-core  $^{10}\text{Be}$  fluxes) are as large as the differences between our deconvolution results and the  
628 reconstructions (see Table 2). As shown in Fig. 11, it is extremely difficult to reconcile the discrepancies between  
629 reconstructed and model-based  $^{14}\text{C}$  production on the basis of carbon cycle changes alone. Nonetheless, the fact  
630 remains that two independent estimates of the  $^{14}\text{C}$  production rate (estimates inferred from paleointensity data and  
631 from ice-core  $^{10}\text{Be}$  fluxes) show systematically lower rates than those obtained by our model-based deconvolution of  
632 atmospheric  $\Delta^{14}\text{C}$ . The differences between the results shown in Fig. 10 and Fig. 11 and Table 2 stem from various  
633 uncertainties that are discussed next.

634

635         Uncertainties associated with the glacial ocean carbon cycle (Fig. 10, colored shading; Fig. 11, colored lines)  
636 are systematic in our approach. The deconvolutions, e.g., of the Hulu Cave  $\Delta^{14}\text{C}$  record, under different model  
637 scenarios are offset against one another, whereas the millennial-scale variability is maintained (see Fig. 11). We do  
638 not attempt to resolve uncertainties associated with Dansgaard-Oeschger warming events and related Antarctic and  
639 tropical climatic excursions in the model runs. Such climatic events may have influenced the atmospheric radiocarbon  
640 budget, but their influence on long-term variations in atmospheric  $\Delta^{14}\text{C}$ , and therefore inferred production rates, is  
641 presumably limited. As may be expected, the lowest production rates (the lowest  $F_{as}$  values) are found in extreme  
642 scenario VENTx and the highest in scenarios PAL and BIO, mirroring the high and low glacial  $\Delta^{14}\text{C}$  levels achieved  
643 by these model scenarios as discussed in Sect. 3.3. Note that there is a large uncertainty in the model-based  $^{14}\text{C}$   
644 production rate stemming from uncertainties associated with the reconstruction of past changes in atmospheric  $\Delta^{14}\text{C}$ ,  
645 in particular the older portion of the  $\Delta^{14}\text{C}$  record.

646

647         A shortcoming of paleointensity-based reconstructions of the  $^{14}\text{C}$  production rate is that they neglect changes  
648 in the solar modulation of the cosmic radiation. The solar modulation potential, which describes the impact of the  
649 solar magnetic field on isotope production, varied between 100 and 1200 MeV during the Holocene on decadal to  
650 centennial time scales, with a median value of approximately 565 MeV (Roth and Joos, 2013). A halving of the solar  
651 modulation potential (e.g., from 600 to 300 MeV) increases the  $^{14}\text{C}$  production rate by about 25 percent for the modern  
652 geomagnetic field strength (Roth and Joos, 2013; see their Fig. 13). This sensitivity remains similar when changes in  
653 the strength of the geomagnetic field are limited as during the last ~35 kyr (Muscheler and Heikkilä, 2011). A shift to  
654 lower solar modulation potential could have materialized if the sun spent on average more time in the postulated  
655 “Grand Minimum” mode (Usoskin et al., 2014) during the last glacial than during the Holocene. The sensitivity of  
656 isotope production to variations in solar modulation potential becomes large during the Laschamp event when the  
657 intensity of the geomagnetic field was close to zero and changes in the solar modulation of the cosmic ray flux may  
658 have a discernible impact on the high  $\Delta^{14}\text{C}$  levels found over this period. A reduction of the solar modulation potential  
659 from 600 to 0 MeV would double  $^{14}\text{C}$  production during times of zero geomagnetic field strength (Masarik and Beer,  
660 2009). However, it is likely that changes in the solar modulation potential were insufficient to explain the discrepancy  
661 between paleointensity-based production rate estimates and the results of our deconvolution, in particular for the post-  
662 Laschamp period and for the reconstruction by Channell et al. (2018). Uncertainties associated with the paleointensity-



663 based reconstructions stem also from uncertainties in estimating the age-scales of the marine sediments and the  
664 geomagnetic field data.

665

666 The ice-core  $^{10}\text{Be}$ -based reconstruction of past changes in  $^{14}\text{C}$  production reflects, by definition, the combined  
667 influence of changes in the solar and geomagnetic modulation of the cosmic ray flux reaching the Earth. This method,  
668 therefore, avoids a fundamental shortcoming of reconstructions based on geomagnetic field data. The assumption is  
669 that the  $^{10}\text{Be}$  and  $^{36}\text{Cl}$  deposited on polar ice and measured in ice cores scales with the amount of cosmogenic isotopes  
670 in the atmosphere. A difficulty is to extrapolate measurements from a single or a few locations to the global  
671 atmosphere. Changes in climate influence atmospheric transport and deposition of  $^{10}\text{Be}$  as well as the snow  
672 accumulation rate, which affect the ice-core  $^{10}\text{Be}$  concentration (Elsässer et al., 2015). Furthermore, the sensitivity of  
673  $^{10}\text{Be}$  in polar ice versus the sensitivity of total production to magnetic field variations, or “polar bias”, is a point of  
674 debate, but atmospheric transport models (Heikkilä et al., 2009; Field et al., 2006) and data analyses (Bard et al.,  
675 1997; Adolphi and Muscheler, 2016; Adolphi et al., 2018) reach different conclusions about its existence and  
676 magnitude. If a polar bias was present, it would lead to an underestimation of the geomagnetic modulation of the ice-  
677 core  $^{10}\text{Be}$  flux and therefore the  $^{14}\text{C}$  production rate.

678

679 Given the uncertainties associated with the proxy records, it may not be surprising that estimates of the  $^{14}\text{C}$   
680 production rate for the last 50 kyr, as obtained by three fundamentally different methods (geomagnetic field data from  
681 marine sediments,  $^{10}\text{Be}$  and  $^{36}\text{Cl}$  measurements in polar ice cores, and model-based deconvolution of atmospheric  
682  $\Delta^{14}\text{C}$ ), disagree with one another, typically by order 10 percent and sometimes by up to 100 percent. At the same time,  
683 it is intriguing that two independent estimates of the  $^{14}\text{C}$  production rate (i.e., estimates inferred from paleointensity  
684 and ice-core  $^{10}\text{Be}$  data) give values that are systematically lower than what is required to match the  $\Delta^{14}\text{C}$  reconstruction.

685

#### 686 **4 Summary and conclusions**

687

688 It is generally assumed that atmospheric  $\Delta^{14}\text{C}$  is controlled by abiotic processes such as atmospheric  $^{14}\text{C}$  production,  
689 air-sea gas exchange, and ocean circulation and mixing. Here, results from sensitivity experiments with the Bern3D  
690 earth system model of intermediate complexity suggest that atmospheric  $\Delta^{14}\text{C}$  is potentially quite sensitive to the  
691 interaction between sediments and the water column on multimillennial timescales. This rather surprising result is due  
692 to the coupling of ocean circulation and biogenic particle sedimentation via the biological carbon pump, which has  
693 important implications for the ocean carbon inventory. If the model’s ocean carbon cycle is sufficiently perturbed,  
694 e.g., by changing the inputs or parameters controlling ocean circulation and/or gas exchange, the resulting shift in  
695 sedimentation has a significant impact on the total oceanic amount of carbon and therefore the average  $\Delta^{14}\text{C}$  value of  
696 the ocean. On time scales of tens of thousands of years the impact of these changes on atmospheric  $\Delta^{14}\text{C}$  is significant  
697 because of the long time scale associated with changes in the sedimentation of biogenic material. This is important  
698 information for long-term climate studies and paleoclimate modelling efforts concerning  $\Delta^{14}\text{C}$ . Note that the  
699 representation of ocean-sediment interactions in the Bern3D is necessarily simplified compared to reality.



700 Nonetheless, a change in the ocean carbon inventory linked with changing sedimentation should be discussed as one  
701 of the potentially important factors affecting atmospheric  $\Delta^{14}\text{C}$  during the last glacial period.

702

703 The reason for the high  $\Delta^{14}\text{C}$  values exhibited by the glacial atmosphere is still not clear. In order to  
704 investigate potential mechanisms governing glacial  $\Delta^{14}\text{C}$  levels, the Bern3D model is again used as a tool. Results of  
705 model simulations forced only by production changes point out that none of the available reconstructions of the  $^{14}\text{C}$   
706 production rate can explain the full amplitude of  $\Delta^{14}\text{C}$  change during the last glacial. In order to test the sensitivity of  
707 the model results with respect to the ocean carbon cycle state, various model parameters, i.e., different sets of physical  
708 and biogeochemical parameters, were “tuned” to match the glacial  $\text{CO}_2$  level. From this, we find that atmospheric  
709  $\Delta^{14}\text{C}$  is most sensitive to changes in physical model parameters, in particular those controlling ocean circulation and  
710 gas exchange. In order to achieve an atmospheric  $\Delta^{14}\text{C}$  value close to the glacial level, the gas transfer velocity in the  
711 polar regions had to be reduced by 100 percent. If interpreted as being due to a greater extent of permanent sea ice  
712 cover, a reduction in polar air-sea exchange efficiency is a possible explanation for high glacial  $\Delta^{14}\text{C}$  levels. Although  
713 this hypothesis is compelling, such a scenario is not supported by the proxy records of Antarctic sea ice cover  
714 (Gersonde et al., 2005; Allen et al., 2011) and the  $^{13}\text{C}/^{12}\text{C}$  ratio of atmospheric  $\text{CO}_2$  (Eggleston et al., 2016).

715

716 Before model-simulated  $\Delta^{14}\text{C}$  can be taken seriously, it must be demonstrated that the reconstruction of past  
717 changes in  $^{14}\text{C}$  production is reliable. There is, however, a substantial amount of scatter in the paleointensity-based  
718 and ice-core  $^{10}\text{Be}$ -based estimates of  $^{14}\text{C}$  production. Here we adopt an alternative approach to estimating the  $^{14}\text{C}$   
719 production rate, which would indeed benefit from further constraints and lines of supporting evidence. Our  
720 deconvolution-based approach assumes that the  $^{14}\text{C}$  production rate can be derived from an atmospheric radiocarbon  
721 budget, constructed using a prognostic carbon cycle model combined with the atmospheric  $\Delta^{14}\text{C}$  record. Our model  
722 results suggest that the glacial  $^{14}\text{C}$  production rate as inferred from paleointensity data and ice-core  $^{10}\text{Be}$  fluxes may  
723 be underestimated by about 15 percent between 32 and 22 kyr BP. Note that our model-based estimates are associated  
724 with uncertainties arising from the reconstruction of the older portion of the  $\Delta^{14}\text{C}$  record and the model simulation of  
725 the glacial ocean carbon cycle (e.g., a too high glacial air-sea  $\text{CO}_2$  flux). Future improvements in the reconstruction  
726 of past changes in  $^{14}\text{C}$  production and atmospheric  $\Delta^{14}\text{C}$  would open up the possibility of attributing model deficiencies  
727 to real changes in the ocean carbon cycle. Ultimately, an improved knowledge of  $^{14}\text{C}$  production during the last glacial,  
728 as well as more robust constraints on the prevailing climate conditions (e.g., ocean circulation, sea ice cover, and wind  
729 speed), are necessary to elucidate the processes permitting mysteriously high  $\Delta^{14}\text{C}$  levels in the glacial atmosphere.

730

#### 731 **Appendix A: Complete description of the Bern3D model**

732

733 The physical core of the Bern3D model is based on the 3-D rigid-lid ocean model of Edwards et al. (1998) as updated  
734 by Edwards and Marsh (2005). The forcing fields for the model integration are monthly mean wind stress data taken  
735 from NCEP/NCAR (Kalnay et al., 1996). Diapycnal mixing is parameterized with a uniform vertical diffusivity  $K_V$  of  
736  $2 \times 10^{-5} \text{ m s}^{-1}$ . The parameterization of eddy-induced transport is separated from that of isopycnal mixing, using the



737 Gent-McWilliams skew flux (Griffies, 1998). Running at the same temporal and horizontal resolution, the one-layer  
738 energy-moisture balance atmosphere model performs an analysis of the energy budget of the Earth by involving solar  
739 radiation, infrared fluxes, evaporation and precipitation, and sensible and latent heat. The zonally averaged surface  
740 albedo climatology is taken from Kukla and Robinson (1980). Transport of moisture is performed by diffusion and  
741 advection and heat by eddy diffusion.

742

743 The Bern3D ocean carbon cycle model is based on the Ocean Carbon-Cycle Model Intercomparison Project  
744 (OCMIP-2) protocols. Air-sea gas exchange is parameterized using the standard gas transfer formulation adopted for  
745 OCMIP-2, except that the gas transfer velocity  $k_w$  parameterization is a linear function of wind speed (Krakauer et  
746 al., 2006) to which we have added a scale factor of 0.81 to match the observed global ocean inventory of bomb  $^{14}\text{C}$   
747 (Müller et al., 2008). It is assumed that  $\text{CO}_2$  and  $\text{O}_2$  are well-mixed in the atmosphere. Surface boundary conditions  
748 also include a virtual-flux term for biogeochemical tracers (e.g., DIC and ALK) to account for their dilution or  
749 concentration due to implicit freshwater fluxes. Modifications from the original OCMIP-2 biotic protocol include the  
750 prognostic formulation of new/export production as a function of light, temperature, and limiting nutrient  
751 concentrations, where the nutrient uptake follows Michaelis-Menten kinetics. The production of biogenic  $\text{CaCO}_3$  and  
752 opal is computed on the basis of the modelled particulate organic carbon (POC) production and availability of silicate,  
753 with a maximum possible fraction of  $\text{CaCO}_3$  material that can be produced. This threshold value is represented by the  
754  $\text{CaCO}_3$ -to-POC export ratio. In the preindustrial control run, the global mean export ratio  $rr$  is 0.082.

755

756 Biogenic material that has been produced in the 75-m production zone is redistributed over the water column  
757 in order to parameterize the downward particle flux through the water column. A power-law model referred to as the  
758 Martin curve is used to describe the vertical POC flux profile, whereas both  $\text{CaCO}_3$  and opal export are redistributed  
759 over the water column with an exponential curve. POC is remineralized instantaneously back to dissolved form  
760 according to Redfield stoichiometry and with a 250-m length scale  $l_{\text{POC}}$  (i.e., in 250 m, the POC flux declines by  $1 -$   
761  $1/e \approx 63$  percent). Likewise,  $\text{CaCO}_3$  and opal are dissolved within one time step, with  $e$ -folding depths of 5066 and  
762 10,000 m, respectively. Biogenic material reaching the model's sea floor forms the upper boundary condition of the  
763 10-layer sediment model after Heinze et al. (1999) and Gehlen et al. (2006). The sediment model includes four solid  
764 sediment components (POC,  $\text{CaCO}_3$ , opal, and clay) and is based on the sediment advection and accumulation scheme  
765 as in the work of Archer et al. (1993). The rate of POC remineralization in the sediments is primarily determined by  
766 the pore water concentration of oxygen, whereas the sediment mineral dissolution rate is governed by the saturation  
767 state of pore waters with respect to  $\text{CaCO}_3$  or opal. The weathering input of DIC, ALK, and nutrients into the ocean  
768 is added as a constant increment to each wet grid cell along the coastlines. Any material input from terrestrial  
769 weathering is considered "radiocarbon dead". The values for these fluxes were chosen so that at the end of the  
770 preindustrial spin-up, input (weathering) and output (sedimentation) are balanced. Additional details concerning the  
771 sediment model are provided in Tschumi et al. (2011) and the appendix of Jeltsch-Thömmes et al. (2019).

772



773           The exchange of any isotopic perturbation between the atmosphere and the terrestrial biosphere is simulated  
774 by use of the four-box model of Siegenthaler and Oeschger (1987). The terrestrial biosphere is represented by four  
775 well-mixed compartments (ground vegetation plus leaves, wood, detritus, and soils), with a fixed total carbon  
776 inventory of 2220 Gt C. Net primary production is balanced by respiration of detritus and soils, and is set to 60 Gt C  
777 per year.

778

779 **Data availability.** All data generated or analyzed during this study can be made available upon request to the  
780 corresponding author (A.D.).

781

782 **Author contribution.** This study was designed by F.J. and A.D. with input from F.A. A.D. developed and  
783 performed the model simulations. F.A. provided production data. A.D. wrote the manuscript with contributions from  
784 the co-authors.

785

786 **Competing interests.** The authors declare that they have no conflict of interest.

787

788 **Acknowledgements.** This work was made possible by the Swiss National Science Foundation (#200020\_172476)  
789 and by the UniBE international 2021 fellowship program of the U. Bern. F.A. was supported by the Swedish  
790 Research Council (Vetenskapsrådet DNR: 2016-00218).

791

## 792 **References**

793

794 Adolphi, F., and Muscheler, R.: Synchronizing the Greenland ice core and radiocarbon timescales over the Holocene  
795 – Bayesian wiggle-matching of cosmogenic radionuclide records, *Climate of the Past*, 12, 15–30, 2016.

796 Adolphi, F., Muscheler, R., Svensson, A., Aldahan, A., Possnert, G., Beer, J., . . . Thiéblemont, R.: Persistent link  
797 between solar activity and Greenland climate during the Last Glacial Maximum, *Nature Geoscience*, 7,  
798 662–666, 2014.

799 Adolphi, F., Ramsey, C. B., Erhardt, T., Edwards, R. L., Cheng, H., Turney, C. S., . . . Muscheler, R.: Connecting  
800 the Greenland ice-core and U/Th timescales via cosmogenic radionuclides: testing the synchronicity of  
801 Dansgaard–Oeschger events, *Climate of the Past*, 14, 1755–1781, 2018.

802 Allen, C. S., Pike, J., and Pudsey, C. J.: Last glacial–interglacial sea-ice cover in the SW Atlantic and its potential  
803 role in global deglaciation, *Quaternary Science Reviews*, 30, 2446–2458, 2011.

804 Archer, D., and Maier-Reimer, E.: Effect of deep-sea sedimentary calcite preservation on atmospheric CO<sub>2</sub>  
805 concentration, *Nature*, 367, 260–263, 1994.

806 Archer, D., Lyle, M., Rodgers, K., and Froelich, P.: What controls opal preservation in tropical deep-sea sediments?,  
807 *Paleoceanography*, 8, 7–21, 1993.

808 Archer, D., Winguth, A., Lea, D., and Mahowald, N.: What caused the glacial/interglacial atmospheric pCO<sub>2</sub>  
809 cycles?, *Reviews of Geophysics*, 38, 159–189, 2000.



- 810 Audi, G., Bersillon, O., Blachot, J., and Wapstra, A. H.: The Nubase evaluation of nuclear and decay properties,  
811 Nuclear Physics A, 729, 3–128, 2003.
- 812 Bard, E., Raisbeck, G. M., Yiou, F., and Jouzel, J.: Solar modulation of cosmogenic nuclide production over the last  
813 millennium: comparison between  $^{14}\text{C}$  and  $^{10}\text{Be}$  records, Earth and Planetary Science Letters, 150, 453–462,  
814 1997.
- 815 Baumgartner, S., Beer, J., Wagner, G., Kubik, P., Suter, M., Raisbeck, G. M., and Yiou, F.:  $^{10}\text{Be}$  and dust, Nuclear  
816 Instruments and Methods in Physics Research Section B: Beam Interactions with Materials and Atoms,  
817 123, 296–301, 1997.
- 818 Baumgartner, S., Beer, J., Masarik, J., Wagner, G., Meynadier, L., and Synal, H.-A.: Geomagnetic Modulation of  
819 the  $^{36}\text{Cl}$  Flux in the GRIP Ice Core, Greenland, Science, 279, 1330–1332, 1998.
- 820 Bé, M.-M., Chisté, V., Dulieu, C., Mougeot, X., Chechev, V., Kondev, F., . . . Wang, B.: Table of Radionuclides  
821 (Comments on evaluations), Monographie BIPM-5, 7, 2013.
- 822 Berger, A. L.: Long-term variations of daily insolation and Quaternary climatic changes, Journal of the Atmospheric  
823 Sciences, 35, 2362–2367, 1978.
- 824 Broecker, W., and Barker, S.: A 190‰ drop in atmosphere's  $\Delta^{14}\text{C}$  during the "Mystery Interval" (17.5 to 14.5 kyr),  
825 Earth and Planetary Science Letters, 256, 90–99, 2007.
- 826 Broecker, W. S., and Peng, T.-H.: Gas exchange rates between air and sea, Tellus, 26, 21–35, 1974.
- 827 Bronk Ramsey, C., Staff, R. A., Bryant, C. L., Brock, F., Kitagawa, H., van der Plicht, J., Schlolaut, G., Marshall,  
828 M. H., Brauer, A., Lamb, H. F., Payne, R. L., Tarasov, P. E., Haraguchi, T., Gotanda, K., Yonenobu, H.,  
829 Yokoyama, Y., Tada, R., and Nakagawa, T.: A complete terrestrial radiocarbon record for 11.2 to 52.8 kyr  
830 B.P., Science, 338, 370–374, 2012.
- 831 Brovkin, V., Ganopolski, A., Archer, D., and Munhoven, G.: Glacial  $\text{CO}_2$  cycle as a succession of key physical and  
832 biogeochemical processes, Climate of the Past, 8, 251–264, 2012.
- 833 Channell, J. E., Hodell, D. A., Crowhurst, S. J., Skinner, L. C., and Muscheler, R.: Relative paleointensity (RPI) in  
834 the latest Pleistocene (10–45 ka) and implications for deglacial atmospheric radiocarbon, Quaternary  
835 Science Reviews, 191, 57–72, 2018.
- 836 Cheng, H., Edwards, R. L., Southon, J., Matsumoto, K., Feinberg, J. M., Sinha, A., . . . Ning, Y.: Atmospheric  
837  $^{14}\text{C}/^{12}\text{C}$  changes during the last glacial period from Hulu Cave, Science, 362, 1293–1297, 2018.
- 838 Delaygue, G., Stocker, T. F., Joos, F., and Plattner, G.-K.: Simulation of atmospheric radiocarbon during abrupt  
839 oceanic circulation changes: trying to reconcile models and reconstructions, Quaternary Science Reviews,  
840 22, 1647–1658, 2003.
- 841 Edwards, N. R., and Marsh, R.: Uncertainties due to transport-parameter sensitivity in an efficient 3-D ocean-  
842 climate model, Climate Dynamics, 24, 415–433, 2005.
- 843 Edwards, N. R., Willmott, A. J., and Killworth, P. D.: On the Role of Topography and Wind Stress on the Stability  
844 of the Thermohaline Circulation, Journal of Physical Oceanography, 28, 756–778, 1998.
- 845 Eggleston, S., Schmitt, J., Bereiter, B., Schneider, R., and Fischer, H.: Evolution of the stable carbon isotope  
846 composition of atmospheric  $\text{CO}_2$  over the last glacial cycle, Paleoceanography, 31, 434–452, 2016.



- 847 Elsässer, C., Wagenbach, D., Levin, I., Stanzick, A., Christl, M., Wallner, A., . . . Dibb, J.: Simulating ice core  $^{10}\text{Be}$   
848 on the glacial–interglacial timescale, *Climate of the Past*, 11, 115–133, 2015.
- 849 Field, C. V., Schmidt, G. A., Koch, D., and Salyk, C.: Modeling production and climate-related impacts on  $^{10}\text{Be}$   
850 concentration in ice cores, *Journal of Geophysical Research: Atmospheres*, 111,  
851 <https://doi.org/10.1029/2005JD006410>, 2006.
- 852 Finkel, R. C., and Nishiizumi, K.: Beryllium 10 concentrations in the Greenland Ice Sheet Project 2 ice core from 3–  
853 40 ka, *Journal of Geophysical Research: Oceans*, 102, 26699–26706, 1997.
- 854 Fischer, H., Schmitt, J., Lüthi, D., Stocker, T. F., Tschumi, T., Parekh, P., . . . Wolff, E.: The role of Southern Ocean  
855 processes in orbital and millennial  $\text{CO}_2$  variations – A synthesis, *Quaternary Science Reviews*, 29, 193–  
856 205, 2010.
- 857 Ganopolski, A., and Brovkin, V.: Simulation of climate, ice sheets and  $\text{CO}_2$  evolution during the last four glacial  
858 cycles with an Earth system model of intermediate complexity, *Climate of the Past*, 13, 1695–1716, 2017.
- 859 Gehlen, M., Bopp, L. E., Aumont, O., Heinze, C., and Ragueneau, O.: Reconciling surface ocean productivity,  
860 export fluxes and sediment composition in a global biogeochemical ocean model, *Biogeosciences*, 3, 521–  
861 537, 2006.
- 862 Gersonde, R., Crosta, X., Abelmann, A., and Armand, L.: Sea-surface temperature and sea ice distribution of the  
863 Southern Ocean at the EPILOG Last Glacial Maximum—a circum-Antarctic view based on siliceous  
864 microfossil records, *Quaternary Science Reviews*, 24, 869–896, 2005.
- 865 Gkinis, V., Simonsen, S. B., Buchardt, S. L., White, J. W., and Vinther, B. M.: Water isotope diffusion rates from  
866 the NorthGRIP ice core for the last 16,000 years – Glaciological and paleoclimatic implications, *Earth and*  
867 *Planetary Science Letters*, 405, 132–141, 2014.
- 868 Griffies, S. M.: The Gent-McWilliams Skew Flux, *Journal of Physical Oceanography*, 28, 831–841, 1998.
- 869 Hain, M. P., Sigman, D. M., and Haug, G. H.: Distinct roles of the Southern Ocean and North Atlantic in the  
870 deglacial atmospheric radiocarbon decline, *Earth and Planetary Science Letters*, 394, 198–208, 2014.
- 871 Heikkilä, U., Beer, J., and Feichter, J.: Meridional transport and deposition of atmospheric  $^{10}\text{Be}$ , *Atmospheric*  
872 *Chemistry and Physics*, 9, 515–527, 2009.
- 873 Heikkilä, U., Phipps, S. J., and Smith, A. M.:  $^{10}\text{Be}$  in late deglacial climate simulated by ECHAM5-HAM – Part 1:  
874 Climatological influences on  $^{10}\text{Be}$  deposition, *Climate of the Past*, 9, 2641–2649, 2013.
- 875 Heinze, C., Maier-Reimer, E., Winguth, A. M., and Archer, D.: A global oceanic sediment model for long-term  
876 climate studies, *Global Biogeochemical Cycles*, 13, 221–250, 1999.
- 877 Herbst, K., Muscheler, R., and Heber, B.: The new local interstellar spectra and their influence on the production  
878 rates of the cosmogenic radionuclides  $^{10}\text{Be}$  and  $^{14}\text{C}$ , *Journal of Geophysical Research: Space Physics*, 122,  
879 23–34, 2017.
- 880 Hoff, U., Rasmussen, T. L., Stein, R., Ezat, M. M., and Fahl, K.: Sea ice and millennial-scale climate variability in  
881 the Nordic seas 90 kyr ago to present, *Nature Communications*, 7, doi:10.1038/ncomms12247, 2016.
- 882 Hughen, K., Lehman, S., Southon, J., Overpeck, J., Marchal, O., Herring, C., and Turnbull, J.:  $^{14}\text{C}$  Activity and  
883 Global Carbon Cycle Changes over the Past 50,000 Years, *Science*, 303, 202–207, 2004.





- 884 Huiskamp, W. N., and Meissner, K. J.: Oceanic carbon and water masses during the Mystery Interval: A model-data  
885 comparison study, *Paleoceanography and Paleoclimatology*, 27, <https://doi.org/10.1029/2012PA002368>,  
886 2012.
- 887 Jeltsch-Thömmes, A., Battaglia, G., Cartapanis, O., Jaccard, S. L., and Joos, F.: Low terrestrial carbon storage at the  
888 Last Glacial Maximum: constraints from multi-proxy data, *Climate of the Past*, 15, 849–879, 2019.
- 889 Köhler, P., Muscheler, R., and Fischer, H.: A model-based interpretation of low-frequency changes in the carbon  
890 cycle during the last 120,000 years and its implications for the reconstruction of atmospheric  $\Delta^{14}\text{C}$ ,  
891 *Geochemistry Geophysics Geosystems*, 7, 1–22, 2006.
- 892 Köhler, P., Nehrbass-Ahles, C., Schmitt, J., Stocker, T. F., and Fischer, H.: A 156 kyr smoothed history of the  
893 atmospheric greenhouse gases  $\text{CO}_2$ ,  $\text{CH}_4$ , and  $\text{N}_2\text{O}$  and their radiative forcing, *Earth System Science Data*,  
894 9, 363–387, 2017.
- 895 Kalnay, E., Kanamitsu, M., Kistler, R., Collins, W., Deaven, D., Gandin, L., . . . Joseph, D.: The NCEP/NCAR 40-  
896 Year Reanalysis Project, *Bulletin of the American Meteorological Society*, 77, 437–471, 1996.
- 897 Key, R. M., Kozyr, A., Sabine, C. L., Lee, K., Wanninkhof, R., Bullister, J. L., . . . Peng, T.-H.: A global ocean  
898 carbon climatology: Results from Global Data Analysis Project (GLODAP), *Global Biogeochemical*  
899 *Cycles*, 18, <https://doi.org/10.1029/2004GB002247>, 2004.
- 900 Kovaltsov, G. A., and Usoskin, I. G.: A new 3D numerical model of cosmogenic nuclide  $^{10}\text{Be}$  production in the  
901 atmosphere, *Earth and Planetary Science Letters*, 291, 182–188, 2010.
- 902 Krakauer, N. Y., Randerson, J. T., Primeau, F. W., Gruber, N., and Menemenlis, D.: Carbon isotope evidence for the  
903 latitudinal distribution and wind speed dependence of the air-sea gas transfer velocity, *Tellus B: Chemical*  
904 *and Physical Meteorology*, 58, 390–417, 2006.
- 905 Kukla, G., and Robinson, D.: Annual Cycle of Surface Albedo, *Monthly Weather Review*, 108, 56–68, 1980.
- 906 Laj, C., Kissel, C., Mazaud, A., Channell, J. E., and Beer, J.: North Atlantic palaeointensity stack since 75ka  
907 (NAPIS-75) and the duration of the Laschamp event, *Philosophical Transactions of the Royal Society of*  
908 *London. Series A: Mathematical, Physical and Engineering Sciences*, 358, 1009–1025, 2000.
- 909 Laj, C., Kissel, C., Mazaud, A., Michel, E., Muscheler, R., and Beer, J.: Geomagnetic field intensity, North Atlantic  
910 Deep Water circulation and atmospheric  $\Delta^{14}\text{C}$  during the last 50 kyr, *Earth and Planetary Science Letters*,  
911 200, 177–190, 2002.
- 912 Laj, C., Kissel, C., and Beer, J.: High resolution global paleointensity stack since 75 kyr (GLOPIS-75) calibrated to  
913 absolute values, *Timescales of the Paleomagnetic Field*, 145, 255–265, 2004.
- 914 Laj, C., Guillou, H., and Kissel, C.: Dynamics of the earth magnetic field in the 10–75 kyr period comprising the  
915 Laschamp and Mono Lake excursions: New results from the French Chaîne des Puy in a global  
916 perspective, *Earth and Planetary Science Letters*, 387, 184–197, 2014.
- 917 Lisiecki, L. E., and Stern, J. V.: Regional and global benthic  $\delta^{18}\text{O}$  stacks for the last glacial cycle,  
918 *Paleoceanography*, 31, 1368–1394, 2016.



- 919 Müller, J., and Stein, R.: High-resolution record of late glacial and deglacial sea ice changes in Fram Strait  
920 corroborates ice–ocean interactions during abrupt climate shifts, *Earth and Planetary Science Letters*, 403,  
921 446–455, 2014.
- 922 Müller, S. A., Joos, F., Edwards, N. R., and Stocker, T. F.: Water Mass Distribution and Ventilation Time Scales in  
923 a Cost-Efficient, Three-Dimensional Ocean Model, *Journal of Climate*, 19, 5479–5499, 2006.
- 924 Müller, S. A., Joos, F., Plattner, G.-K., Edwards, N. R., and Stocker, T. F.: Modeled natural and excess radiocarbon:  
925 Sensitivities to the gas exchange formulation and ocean transport strength, *Global Biogeochemical Cycles*,  
926 22, <https://doi.org/10.1029/2007GB003065>, 2008.
- 927 Marchal, O., Stocker, T. F., and Muscheler, R.: Atmospheric radiocarbon during the Younger Dryas: production,  
928 ventilation, or both?, *Earth and Planetary Science Letters*, 185, 383–395, 2001.
- 929 Mariotti, V., Paillard, D., Bopp, L., Roche, D. M., and Bouttes, N.: A coupled model for carbon and radiocarbon  
930 evolution during the last deglaciation, *Geophysical Research Letters*, 43, 1306–1313, 2016.
- 931 Masarik, J., and Beer, J.: Simulation of particle fluxes and cosmogenic nuclide production in the Earth's atmosphere,  
932 *Journal of Geophysical Research: Atmospheres*, 104, 12099–12111, 1999.
- 933 Masarik, J., and Beer, J.: An updated simulation of particle fluxes and cosmogenic nuclide production in the Earth's  
934 atmosphere, *Journal of Geophysical Research: Atmospheres*, 114, <https://doi.org/10.1029/2008JD010557>,  
935 2009.
- 936 Menviel, L., Joos, F., and Ritz, S. P.: Simulating atmospheric CO<sub>2</sub>, <sup>13</sup>C and the marine carbon cycle during the Last  
937 Glacial–Interglacial cycle: possible role for a deepening of the mean remineralization depth and an increase  
938 in the oceanic nutrient inventory, *Quaternary Science Reviews*, 56, 46–68, 2012.
- 939 Muscheler, R., and Heikkilä, U.: Constraints on long-term changes in solar activity from the range of variability of  
940 cosmogenic radionuclide records, *Astrophysics and Space Sciences Transactions*, 7, 355–364, 2011.
- 941 Muscheler, R., Beer, J., Wagner, G., Laj, C., Kissel, C., Raisbeck, G. M., . . . Kubike, P. W.: Changes in the carbon  
942 cycle during the last deglaciation as indicated by the comparison of <sup>10</sup>Be and <sup>14</sup>C records, *Earth and  
943 Planetary Science Letters*, 219, 325–340, 2004.
- 944 Muscheler, R., Adolphi, F., Herbst, K., and Nilsson, A.: The Revised Sunspot Record in Comparison to Cosmogenic  
945 Radionuclide-Based Solar Activity Reconstructions, *Solar Physics*, 291, 3025–3043, 2016.
- 946 Nowaczyk, N. R., Arz, H. W., Frank, U., Kind, J., and Plessen, B.: Dynamics of the Laschamp geomagnetic  
947 excursion from Black Sea sediments, *Earth and Planetary Science Letters*, 351–352, 54–69, 2012.
- 948 Nowaczyk, N. R., Frank, U., Kind, J., and Arz, H. W.: A high-resolution paleointensity stack of the past 14 to 68 ka  
949 from Black Sea sediments, *Earth and Planetary Science Letters*, 384, 1–16, 2013.
- 950 Orr, J. C., Najjar, R. G., Aumont, O., Bopp, L., Bullister, J. L., Danabasoglu, G., . . . Yool, A.: Biogeochemical  
951 protocols and diagnostics for the CMIP6 Ocean Model Intercomparison Project (OMIP), *Geoscientific  
952 Model Development*, 10, 2169–2199, 2017.
- 953 Parekh, P., Joos, F., and Müller, S. A.: A modeling assessment of the interplay between aeolian iron fluxes and iron-  
954 binding ligands in controlling carbon dioxide fluctuations during Antarctic warm events, *Paleoceanography  
955 and Paleoclimatology*, 23, <https://doi.org/10.1029/2007PA001531>, 2008.



- 956 Peltier, W. R.: Ice Age Paleotopography, *Science*, 265, 195–201, 1994.
- 957 Poluianov, S. V., Kovaltsov, G. A., Mishev, A. L., and Usoskin, I. G.: Production of cosmogenic isotopes  $^7\text{Be}$ ,  $^{10}\text{Be}$ ,  
958  $^{14}\text{C}$ ,  $^{22}\text{Na}$ , and  $^{36}\text{Cl}$  in the atmosphere: Altitudinal profiles of yield functions, *Journal of Geophysical*  
959 *Research: Atmospheres*, 121, 8125–8136, 2016.
- 960 Potgieter, M. S., Vos, E. E., Boezio, M., De Simone, N., Di Felice, V., and Formato, V.: Modulation of Galactic  
961 Protons in the Heliosphere During the Unusual Solar Minimum of 2006 to 2009, *Solar Physics*, 289, 391–  
962 406, 2014.
- 963 Raisbeck, G. M., Cauquoin, A., Jouzel, J., Landais, A., Petit, J.-R., Lipenkov, V. Y., . . . Yiou, F.: An improved  
964 north-south synchronization of ice core records around the 41 kyr  $^{10}\text{Be}$  peak, *Climate of the Past*, 13, 217–  
965 229, 2017.
- 966 Rasmussen, S. O., Abbott, P. M., Blunier, T., Bourne, A. J., Brook, E. J., Buchardt, S. L., . . . Winstrup, M.: A first  
967 chronology for the North Greenland Eemian Ice Drilling (NEEM) ice core, *Climate of the Past*, 9, 2713–  
968 2730, 2013.
- 969 Reimer, P., Bard, E., Bayliss, A., Beck, J., Blackwell, P., Ramsey, C., . . . Van der Plicht, J.: IntCal13 and Marine13  
970 Radiocarbon Age Calibration Curves 0–50,000 Years cal BP, *Radiocarbon*, 55, 1869–1887, 2013.
- 971 Ritz, S. P., Stocker, T. F., and Joos, F.: A Coupled Dynamical Ocean–Energy Balance Atmosphere Model for  
972 Paleoclimate Studies, *Journal of Climate*, 24, 349–375, 2011.
- 973 Roth, R., and Joos, F.: A reconstruction of radiocarbon production and total solar irradiance from the Holocene  $^{14}\text{C}$   
974 and  $\text{CO}_2$  records: implications of data and model uncertainties, *Climate of the Past*, 9, 1879–1909, 2013.
- 975 Roth, R., Ritz, S. P., and Joos, F.: Burial-nutrient feedbacks amplify the sensitivity of atmospheric carbon dioxide to  
976 changes in organic matter remineralisation, *Earth System Dynamics*, 5, 321–343, 2014.
- 977 Sarnthein, M., Schneider, B., and Grootes, P. M.: Peak glacial  $^{14}\text{C}$  ventilation ages suggest major draw-down of  
978 carbon into the abyssal ocean, *Climate of the Past*, 9, 2595–2614, 2013.
- 979 Siegenthaler, U., and Oeschger, H.: Biospheric  $\text{CO}_2$  emissions during the past 200 years reconstructed by  
980 deconvolution of ice core data, *Tellus*, 39B, 140–154, 1987.
- 981 Siegenthaler, U., Heimann, M., and Oeschger, H.:  $^{14}\text{C}$  Variations Caused by Changes in the Global Carbon Cycle,  
982 *Radiocarbon*, 22, 177–191, 1980.
- 983 Skinner, L. C., Fallon, S., Waelbroeck, C., Michel, E., and Barker, S.: Ventilation of the Deep Southern Ocean and  
984 Deglacial  $\text{CO}_2$  Rise, *Science*, 328, 1147–1151, 2010.
- 985 Skinner, L. C., Primeau, F., Freeman, E., de la Fuente, M., Goodwin, P. A., Gottschalk, J., . . . Scrivner, A. E.:  
986 Radiocarbon constraints on the glacial ocean circulation and its impact on atmospheric  $\text{CO}_2$ , *Nature*  
987 *Communications*, 8, 16010, 2017.
- 988 Soulet, G., Skinner, L. C., Beaupré, S. R., and Galy, V.: A Note on Reporting of Reservoir  $^{14}\text{C}$  Disequilibria and  
989 Age Offsets, *Radiocarbon*, 58, 205–211, 2016.
- 990 Southon, J., Noronha, A. L., Cheng, H., Edwards, R. L., and Wang, Y.: A high-resolution record of atmospheric  $^{14}\text{C}$   
991 based on Hulu Cave speleothem H82, *Quaternary Science Reviews*, 33, 32–41, 2012.
- 992 Stuiver, M., and Polach, H. A.: Discussion: Reporting of  $^{14}\text{C}$  Data, *Radiocarbon*, 19, 355–363, 1977.



993 Tschumi, T., Joos, F., and Parekh, P.: How important are Southern Hemisphere wind changes for low glacial carbon  
994 dioxide? A model study, *Paleoceanography and Paleoclimatology*, 23,  
995 <https://doi.org/10.1029/2008PA001592>, 2008.

996 Tschumi, T., Joos, F., Gehlen, M., and Heinze, C.: Deep ocean ventilation, carbon isotopes, marine sedimentation  
997 and the deglacial CO<sub>2</sub> rise, *Climate of the Past*, 7, 771–800, 2011.

998 Usoskin, I. G., Hulot, G., Gallet, Y., Roth, R., Licht, A., Joos, F., . . . Khokhlov, A.: Evidence for distinct modes of  
999 solar activity, *Astronomy & Astrophysics*, 562, 1–4, 2014.

1000 Wagner, G., Beer, J., Masarik, J., Muscheler, R., Kubik, P. W., Mende, W., . . . Yiou, F.: Presence of the solar de  
1001 Vries cycle (~205 years) during the last ice age, *Geophysical Research Letters*, 28, 303–306, 2001.

1002 Wallmann, K., Schneider, B., and Sarnthein, M.: Effects of eustatic sea-level change, ocean dynamics, and nutrient  
1003 utilization on atmospheric pCO<sub>2</sub> and seawater composition over the last 130 000 years: a model study,  
1004 *Climate of the Past*, 12, 339–375, 2016.

1005 Yiou, F., Raisbeck, G. M., Baumgartner, S., Beer, J., Hammer, C., Johnsen, S., . . . Yiou, P.: Beryllium 10 in the  
1006 Greenland Ice Core Project ice core at Summit, Greenland, *Journal of Geophysical Research: Oceans*, 102,  
1007 26783–26794, 1997.

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1024 Table 1. Summary of model scenarios considered in this study. Initial conditions refer to the boundary conditions used  
 1025 for the precursor spin-up simulation needed to initialize the transient simulation. These correspond either to  
 1026 preindustrial (PI) or last glacial conditions. The paleoclimate forcing fields, i.e., Orb-GHG-Ice, are reconstructed  
 1027 changes in orbital parameters (Berger, 1978), greenhouse gas radiative forcing based on reconstructed atmospheric  
 1028 greenhouse gas histories (Köhler et al., 2017), and varying ice sheet extent scaled using the global benthic  $\delta^{18}\text{O}$  stack  
 1029 of Lisiecki and Stern (2016). Numbers refer to the scale factor values applied to the tunable model parameters  $\tau$  (wind  
 1030 stress scale factor),  $K_V$  (vertical diffusivity),  $k_w$  (gas transfer velocity),  $rr$  ( $\text{CaCO}_3$ -to-POC export ratio), and  $\ell_{POC}$   
 1031 (POC remineralization length scale) at the last glacial maximum (LGM). These values were chosen in order to achieve  
 1032 an atmospheric  $\text{CO}_2$  concentration close to the LGM level, and are varied over time using the global benthic  $\delta^{18}\text{O}$   
 1033 stack.

Scenario	Initial conditions	Paleoclimate forcing	Tunable parameters: scale factor at LGM				
			$\tau$	$K_V$	$k_w$	$rr$	$\ell_{POC}$
MOD	PI	-	-	-	-	-	-
PAL	Glacial	Orb-GHG-Ice	-	-	-	-	-
CIRC	Glacial	Orb-GHG-Ice	0.4	0.4	-	-	-
VENT	Glacial	Orb-GHG-Ice	0.4	0.4	0.4	-	-
VENTx	Glacial	Orb-GHG-Ice	0.4	0.4	0.0	-	-
BIO	Glacial	Orb-GHG-Ice	-	-	-	0.8	1.2
PHYS-BIO	Glacial	Orb-GHG-Ice	0.7	0.7	0.7	0.7	-
PHYS-BIOx	Glacial	Orb-GHG-Ice	0.8	0.8	0.8	0.8	1.2

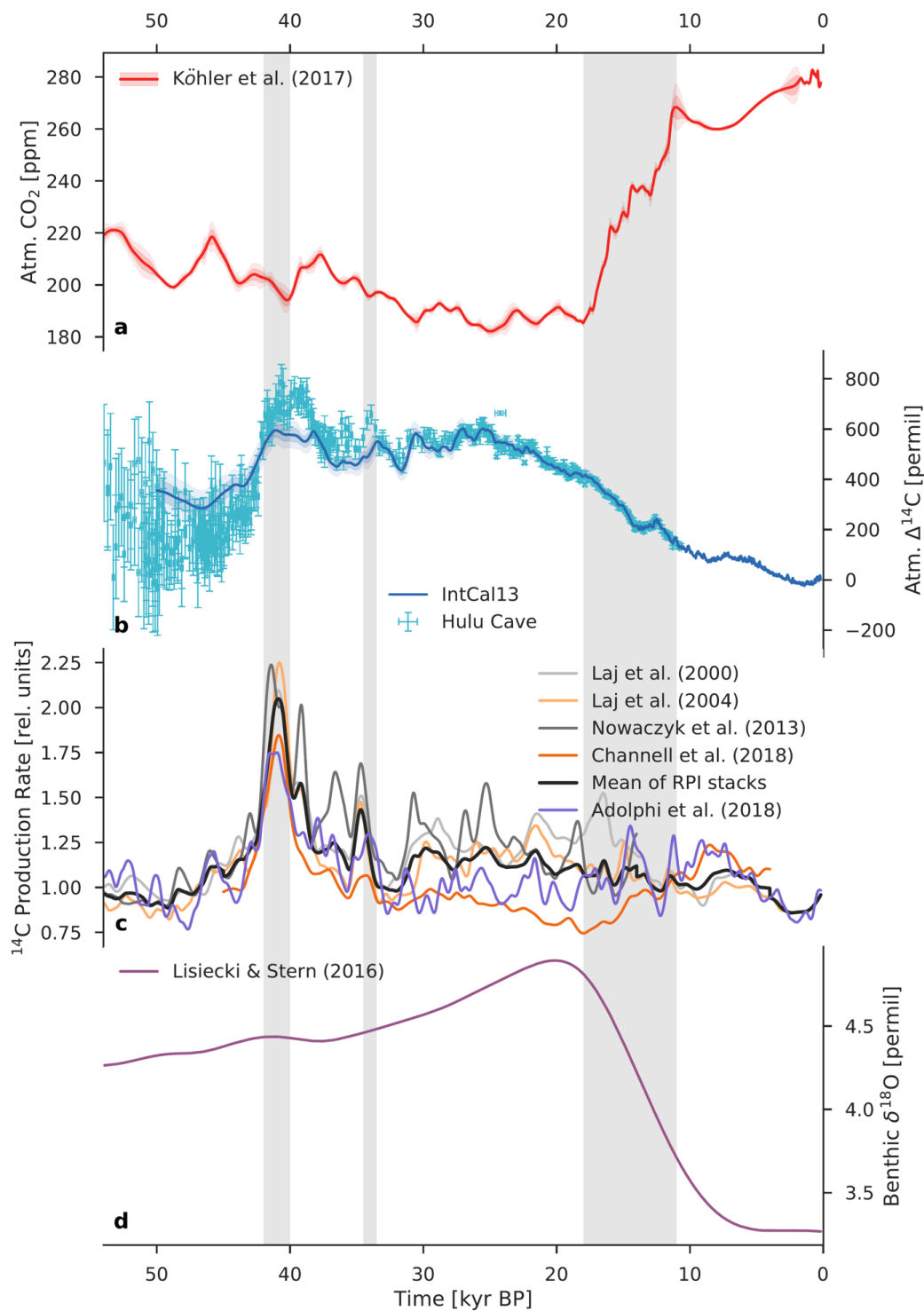
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1046 Table 2. Production rate estimates in relative units inferred from three fundamentally different reconstruction methods:  
 1047 geomagnetic field data from marine sediments,  $^{10}\text{Be}$  and  $^{36}\text{Cl}$  measurements in polar ice cores, and model-based  
 1048 deconvolution of atmospheric  $\Delta^{14}\text{C}$ . Laj00, Laj04, Now13, and Chn18 refer to the paleointensity-based reconstructions  
 1049 of Laj et al. (2000), Laj et al. (2004), Nowaczyk et al. (2013), and Channell et al. (2018), respectively. Adp18 refers  
 1050 to the ice-core  $^{10}\text{Be}$ -based reconstruction of Adolphi et al. (2018). Int13 and Hul18 refer to the model-based  
 1051 reconstructions from this study, using the IntCal13 calibration curve (Reimer et al., 2013) and the new Hulu Cave  
 1052  $\Delta^{14}\text{C}$  dataset (Cheng et al., 2018). The bold numbers show the mean production rates during the last glacial (50 to 18  
 1053 kyr BP).

Time (kyr BP)	Mean production rate (relative units)						
	Laj00	Laj04	Now13	Chn18	Adp18	Int13	Hul18
50 to 42	1.08	1.04	1.12	1.08	1.01	1.23	1.14
42 to 37	1.57	1.56	1.71	1.36	1.44	1.45	1.67
37 to 32	1.19	1.09	1.35	0.98	1.10	1.25	1.28
32 to 22	1.22	1.15	1.29	0.92	0.99	1.31	1.31
22 to 18	1.31	1.20	1.17	0.81	0.98	1.11	1.11
<b>50 to 18</b>	<b>1.25</b>	<b>1.18</b>	<b>1.31</b>	<b>1.01</b>	<b>1.08</b>	<b>1.28</b>	<b>1.29</b>

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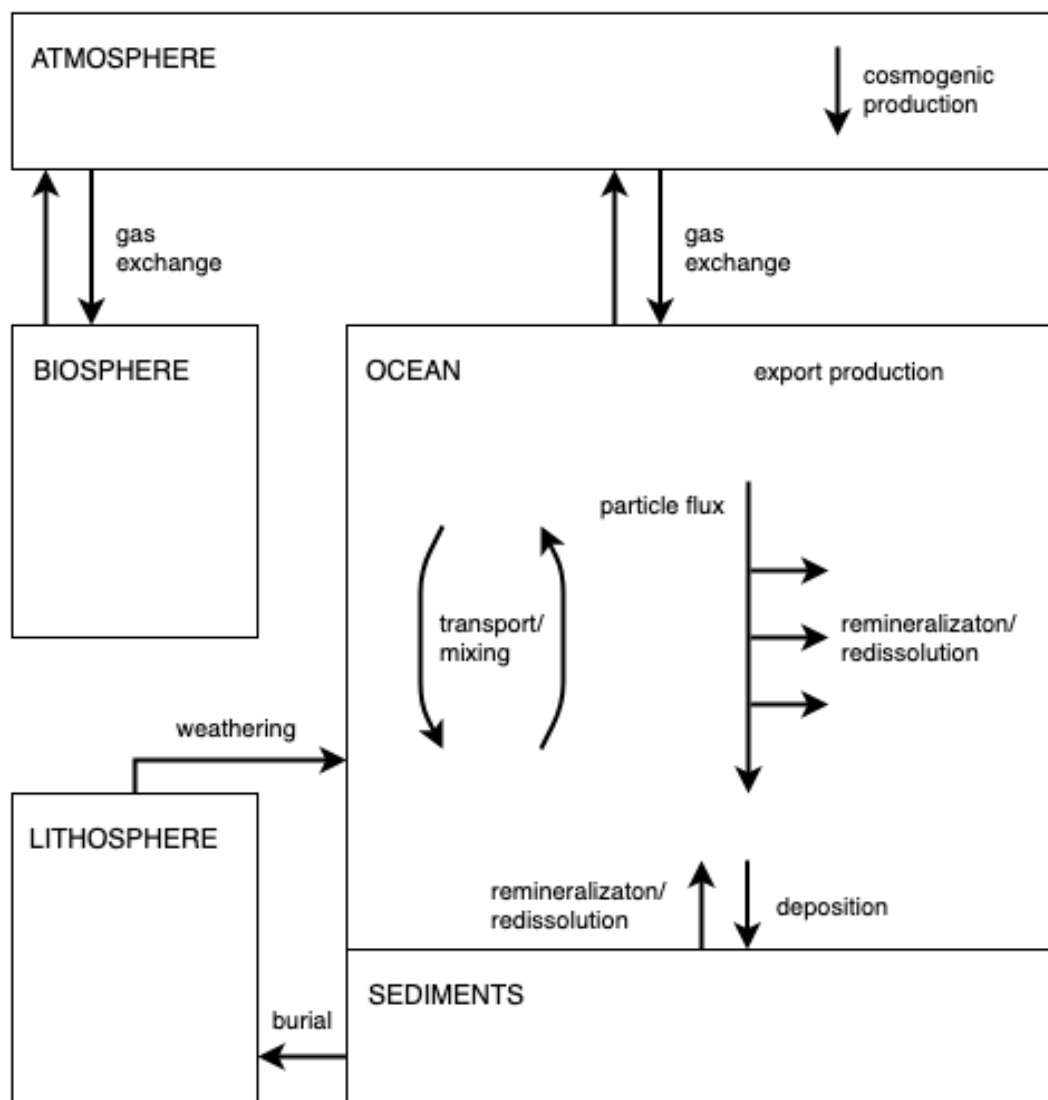




1072 Fig. 1. Comparison of various paleoclimate records for the last 54 kyr. (a) Atmospheric CO<sub>2</sub> from the data compilation  
1073 of Köhler et al. (2017). The light red envelope shows the uncertainty (2σ). (b) Atmospheric Δ<sup>14</sup>C reconstructed from  
1074 <sup>14</sup>C measurements on tree rings, plant macrofossils, speleothems, corals, and foraminifera. The light blue envelope  
1075 shows the uncertainty (2σ) in the IntCal13 calibration curve (Reimer et al., 2013), whereas the Hulu Cave data (Cheng  
1076 et al., 2018) are shown with error bars (1σ). Hulu Cave data are consistent with IntCal13 between ~10.6 and 33.3 kyr  
1077 BP. For both records Δ<sup>14</sup>C values were adjusted to the presently accepted value of the radiocarbon half-life (5700  
1078 years). (c) <sup>14</sup>C production rate in relative units reconstructed from paleointensity data (Laj et al., 2000; Laj et al., 2004;  
1079 Nowaczyk et al., 2013; Channell et al., 2018) and from polar ice-core <sup>10</sup>Be fluxes (Adolphi et al., 2018). The heavy  
1080 dark gray line is the mean paleointensity-based <sup>14</sup>C production rate. (d) Global benthic δ<sup>18</sup>O stack, a proxy for ice  
1081 volume, from Lisiecki and Stern (2016). Three vertical light gray bars indicate the Laschamp excursion (~41 kyr BP),  
1082 when the Earth's geomagnetic dipole field intensity was close to zero, the Mono Lake geomagnetic excursion (~34  
1083 kyr BP), and the last glacial termination (~18 to 11 kyr BP), respectively.

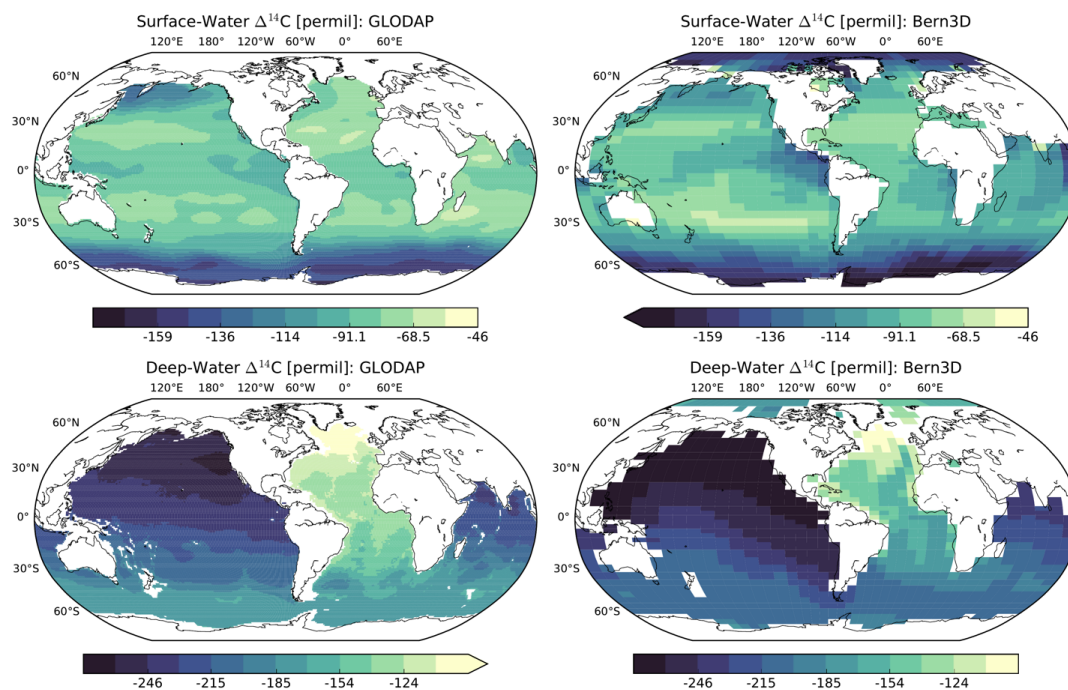
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1102 Fig. 2. Schematic diagram of the Bern3D carbon cycle model. The fully coupled model includes the major global  
1103 carbon reservoirs (atmosphere, terrestrial biosphere, ocean, and sediments) and the exchange fluxes between them.  
1104 Biogeochemical processes, namely, air-sea gas exchange, biological export production, and particle flux through the  
1105 water column, are parameterized by refined OCMIP-2 formulations. Details concerning the model are provided in  
1106 Sect. 2 and Appendix A.

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Fig. 3. Steady-state distribution of  $\Delta^{14}\text{C}$  in the surface ( $> 100$  m) and deep ( $< 1500$  m) ocean for the preindustrial control run (right), compared to the distribution of  $\Delta^{14}\text{C}$  based on the Global Ocean Data Analysis Project (GLODAP).

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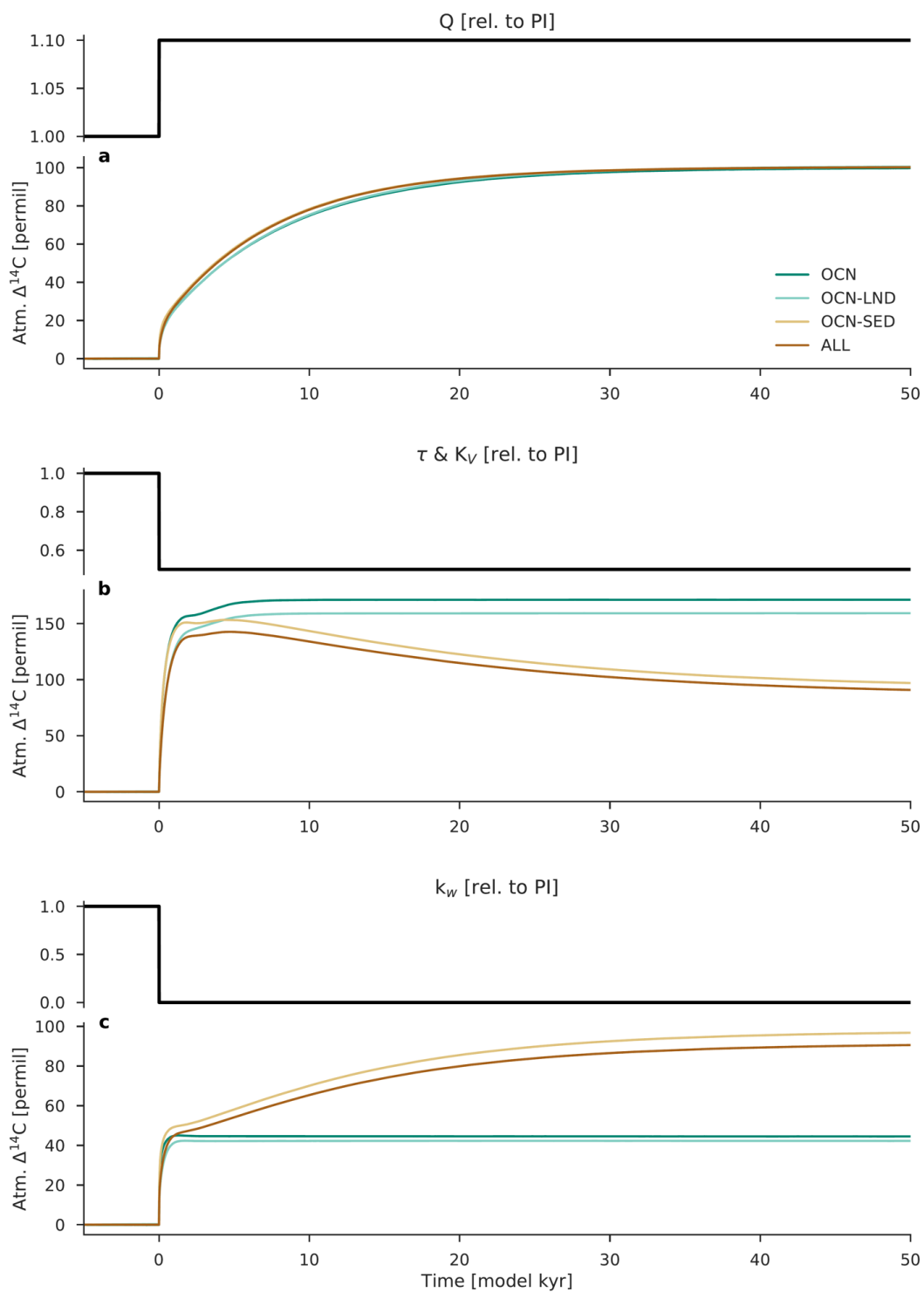
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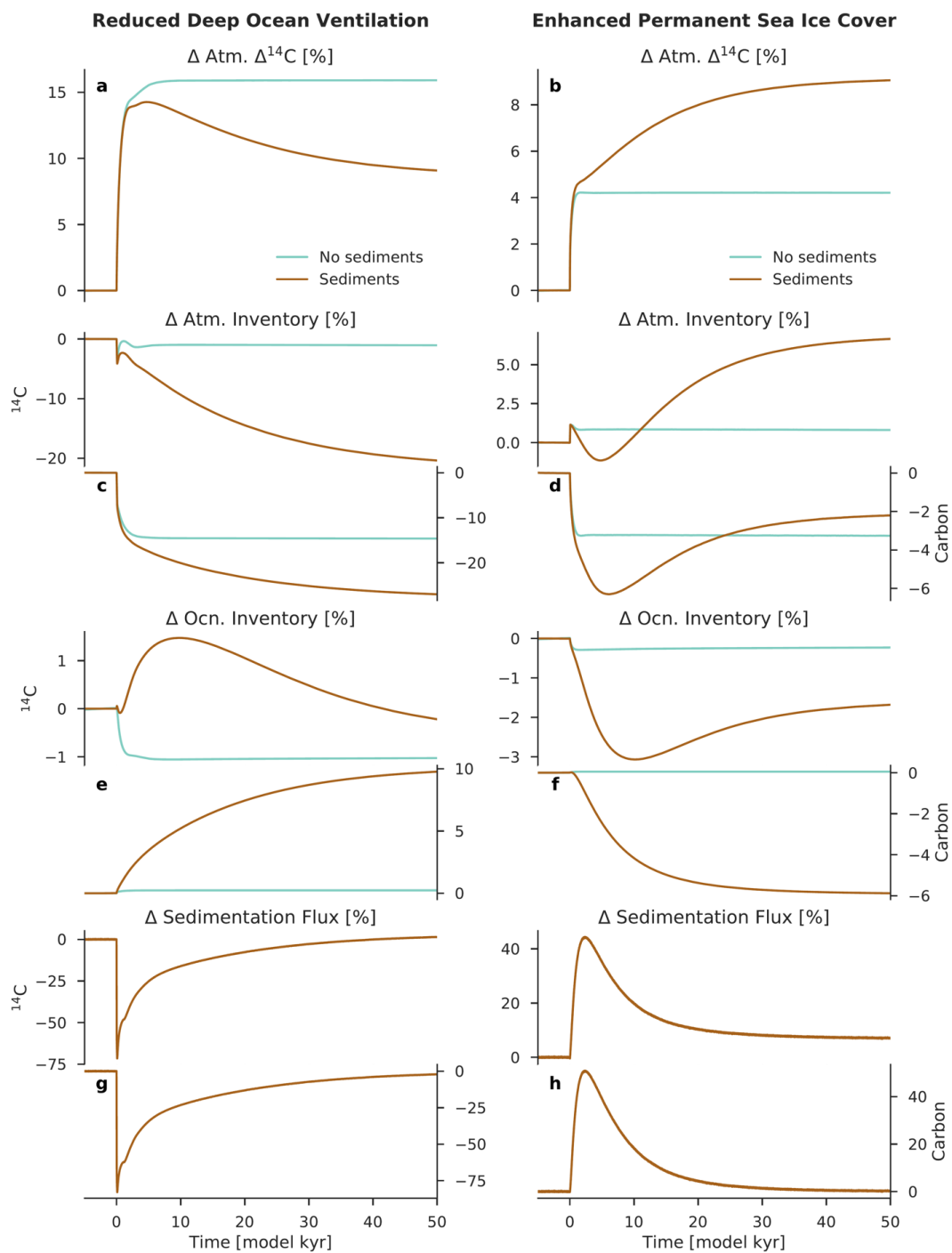
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1130 Fig. 4. Response of atmospheric  $\Delta^{14}\text{C}$  to step changes in  $^{14}\text{C}$  production, followed by step changes in the tunable model  
1131 parameters of the ocean carbon cycle. (a)  $^{14}\text{C}$  production  $Q$  is increased at time 0 from 100 to 110 percent of its  
1132 preindustrial value (“higher production” scenario). (b) Wind stress scale factor  $\tau$  and vertical diffusivity  $K_V$  are  
1133 decreased at time 0 from 100 to 50 percent of their preindustrial values (“reduced deep ocean ventilation” scenario).  
1134 (c) Gas transfer velocity  $k_w$  is decreased at time 0 from 100 to 0 percent of its preindustrial value at the north ( $> 60^\circ\text{N}$ )  
1135 and south ( $> 48^\circ\text{S}$ ) poles (“enhanced permanent sea ice cover” scenario). Four model configurations are considered.  
1136 The dark turquoise line shows the model results using the atmosphere–ocean (OCN) configuration, the light turquoise  
1137 line is the atmosphere–ocean–land (OCN-LND) configuration, the light brown line is the atmosphere–ocean–sediment  
1138 (OCN-SED) configuration, and the dark brown line is the atmosphere–ocean–land–sediment (ALL) configuration.

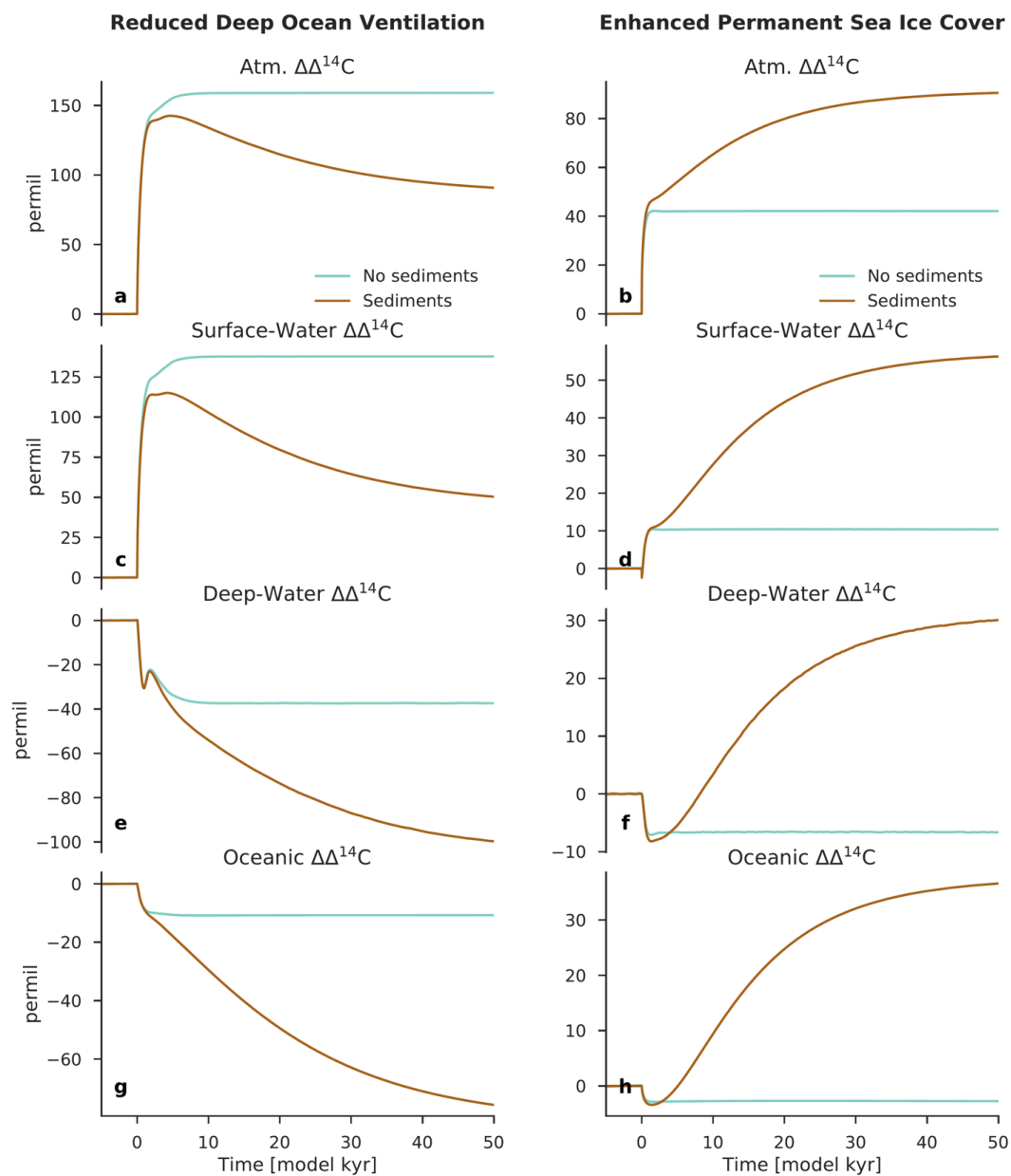
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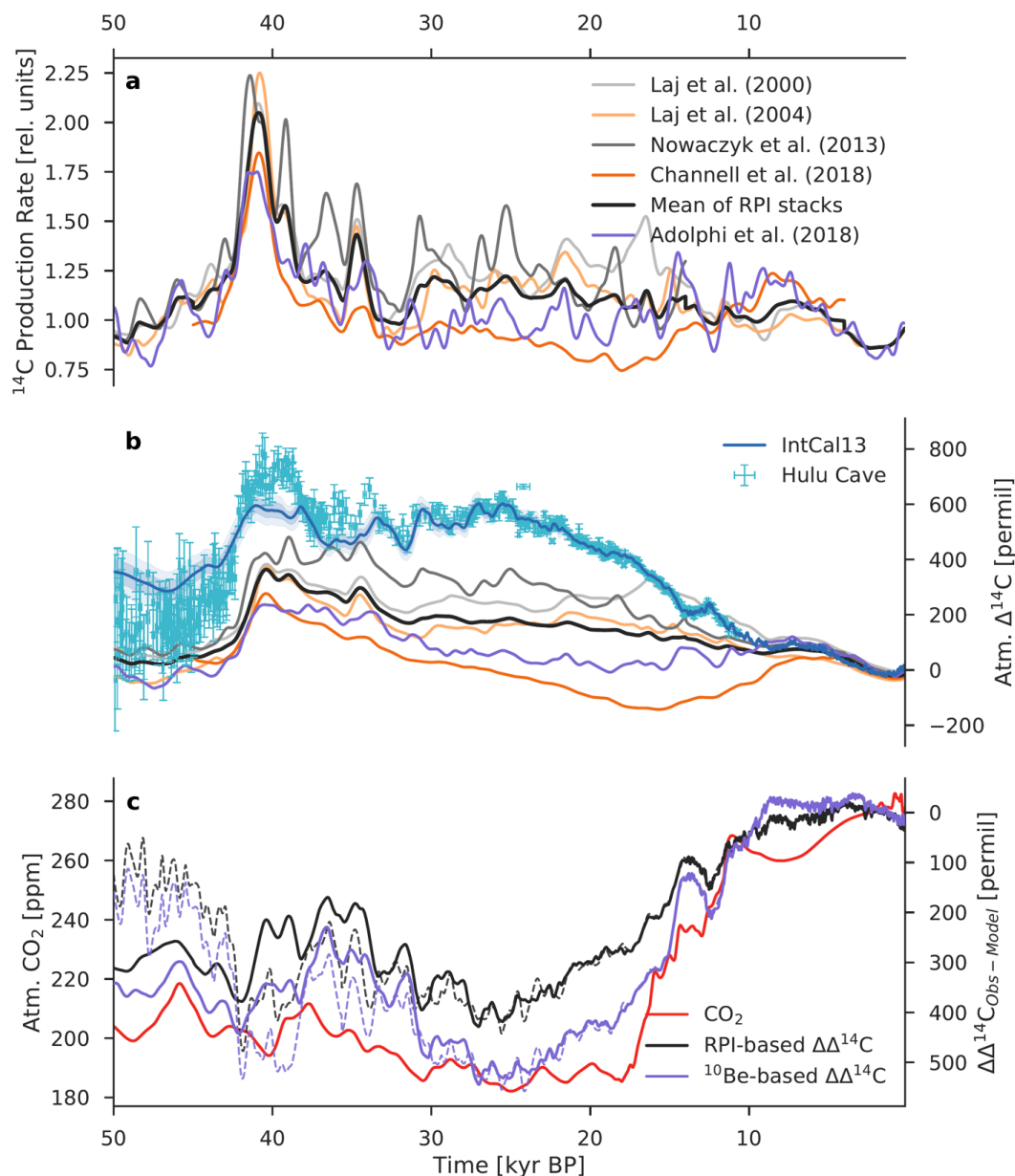
1160 Fig. 5. Changes in carbon reservoir sizes and the sedimentation flux for the scenarios “reduced deep ocean ventilation”  
1161 (left) and “enhanced permanent sea ice cover” (right). The change in atmospheric  $\Delta^{14}\text{C}$  is also shown (a, b). Anomalies  
1162 are expressed here as differences relative to the preindustrial steady state (in percent). Turquoise lines show the model  
1163 results using configuration OCN-LND (no sediments) and brown lines are configuration ALL (model includes  
1164 sediments). The y-axis on the left-hand side of each panel refers to changes in the  $^{14}\text{C}$  inventory, whereas the y-axis  
1165 on the right-hand side of each panel refers to changes in the carbon inventory or flux.

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1192 Fig. 6. Change in  $\Delta^{14}\text{C}$  for the atmosphere, surface ocean, deep ocean, and global ocean for the scenarios “reduced  
1193 deep ocean ventilation” (left) and “enhanced permanent sea ice cover” (right). Anomalies are expressed here as  
1194 differences relative to the preindustrial steady state (in permil). Turquoise lines show the model results using  
1195 configuration OCN-LND (no sediments) and brown lines are configuration ALL (model includes sediments).

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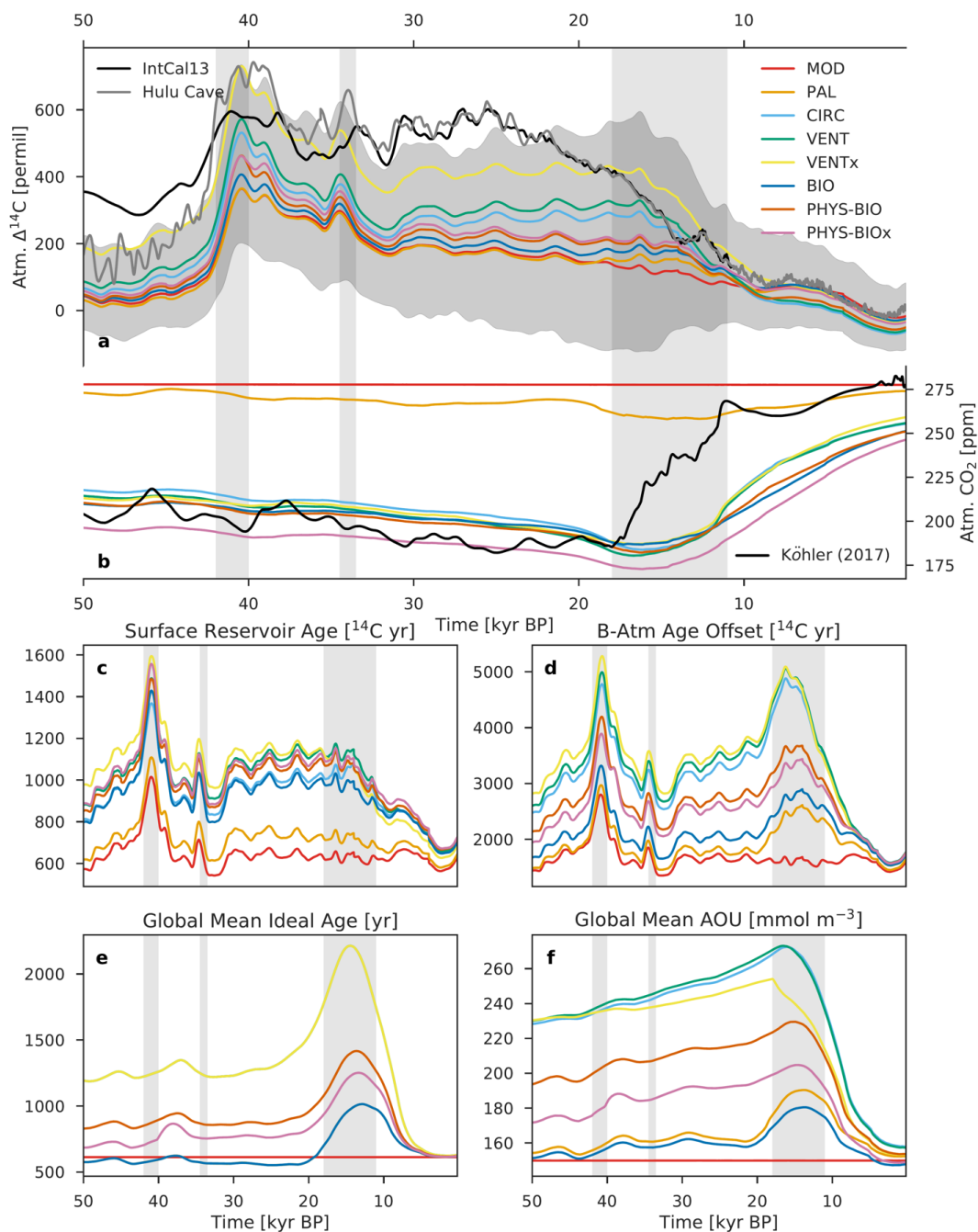
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1198 Fig. 7. Component of atmospheric  $\Delta^{14}\text{C}$  variability caused by production changes alone. (a) Relative  $^{14}\text{C}$  production  
1199 rate as inferred from paleointensity data (gray) and from polar ice-core  $^{10}\text{Be}$  fluxes (purple). The heavy dark gray line  
1200 is the mean paleointensity-based  $^{14}\text{C}$  production rate. (b) Modelled  $\Delta^{14}\text{C}$  records based only on  $^{14}\text{C}$  production changes,  
1201 compared with the reconstructed IntCal13 and Hulu Cave  $\Delta^{14}\text{C}$  records. The modelled records are given by scenario  
1202 MOD that assumes a constant preindustrial carbon cycle. (c) Difference between reconstructed  $\Delta^{14}\text{C}$  and model-  
1203 simulated  $\Delta^{14}\text{C}$  using paleointensity data (RPI-based  $\Delta\Delta^{14}\text{C}$ ; gray) and ice-core  $^{10}\text{Be}$  data ( $^{10}\text{Be}$ -based  $\Delta\Delta^{14}\text{C}$ ; purple),





1204 compared with the atmospheric CO<sub>2</sub> record (red). Solid lines show the IntCal13–model difference, whereas dashed  
1205 lines show the Hulu–model difference. The  $\Delta\Delta^{14}\text{C}$  curve indicates changes in  $\Delta^{14}\text{C}$  that can be attributed to some  
1206 combination of carbon cycle changes, uncertainties in the reconstruction of the <sup>14</sup>C production rate, and uncertainties  
1207 in the IntCal13 and Hulu Cave  $\Delta^{14}\text{C}$  records.

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Fig. 8. Modelled records of atmospheric (a)  $\Delta^{14}\text{C}$  and (b)  $\text{CO}_2$ , compared with their reconstructed atmospheric histories (black and gray lines). Also shown are modelled records of (c) surface reservoir age, (d) B-Atm  $^{14}\text{C}$  age offset, (e) ideal age, and (f) apparent oxygen utilization (AOU). Colored lines show the results of model runs using the mean



1236 paleointensity-based  $^{14}\text{C}$  production rate and the eight different carbon cycle scenarios described in Sect. 2.4 and Table  
1237 1. The gray envelope in (a) shows the uncertainty ( $2\sigma$ ) from all production rate reconstructions and carbon cycle  
1238 scenarios, providing a bounded estimate of  $\Delta^{14}\text{C}$  change. Radiocarbon ventilation ages are expressed here as  
1239 radiocarbon reservoir age offsets following Soulet et al. (2016) which are used extensively by the radiocarbon dating  
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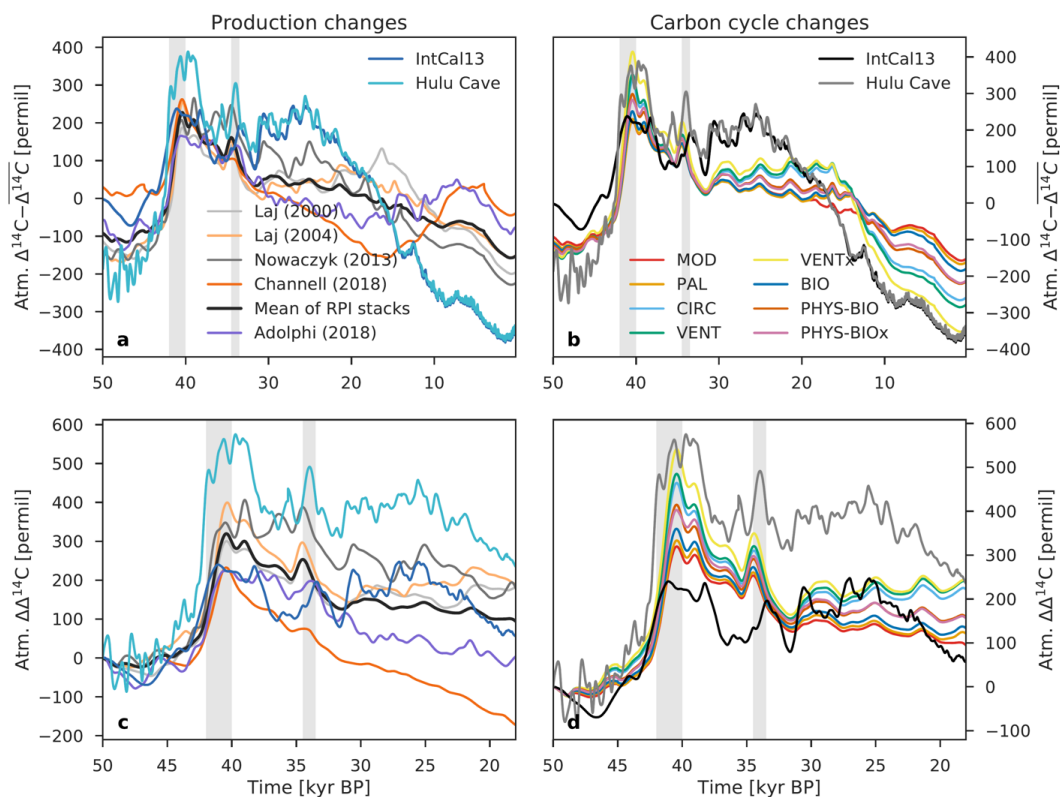
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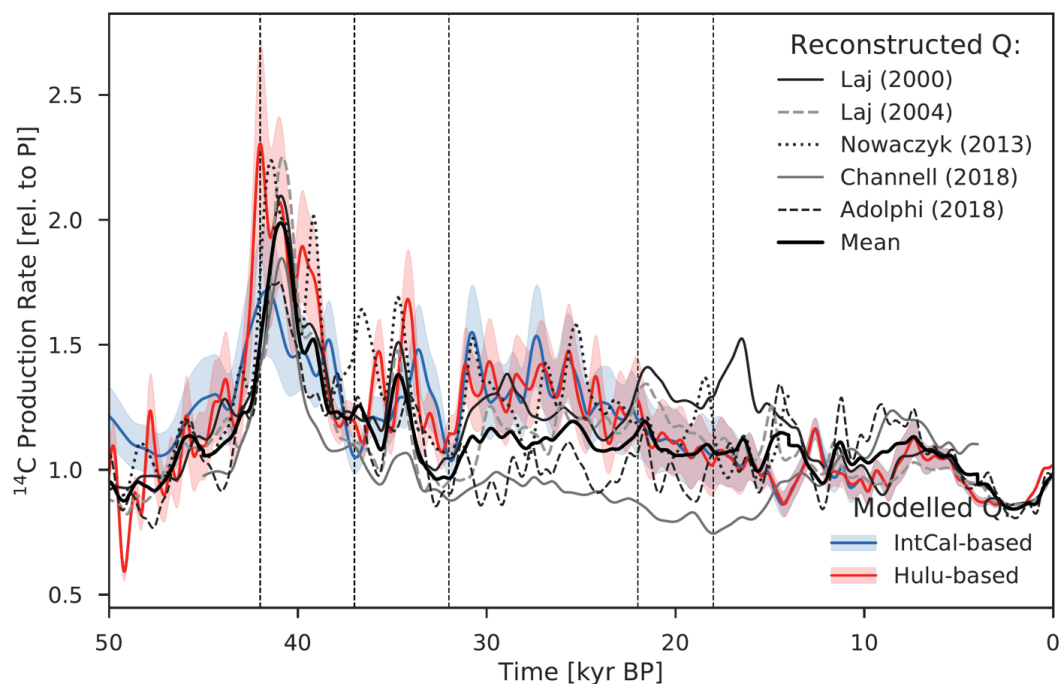
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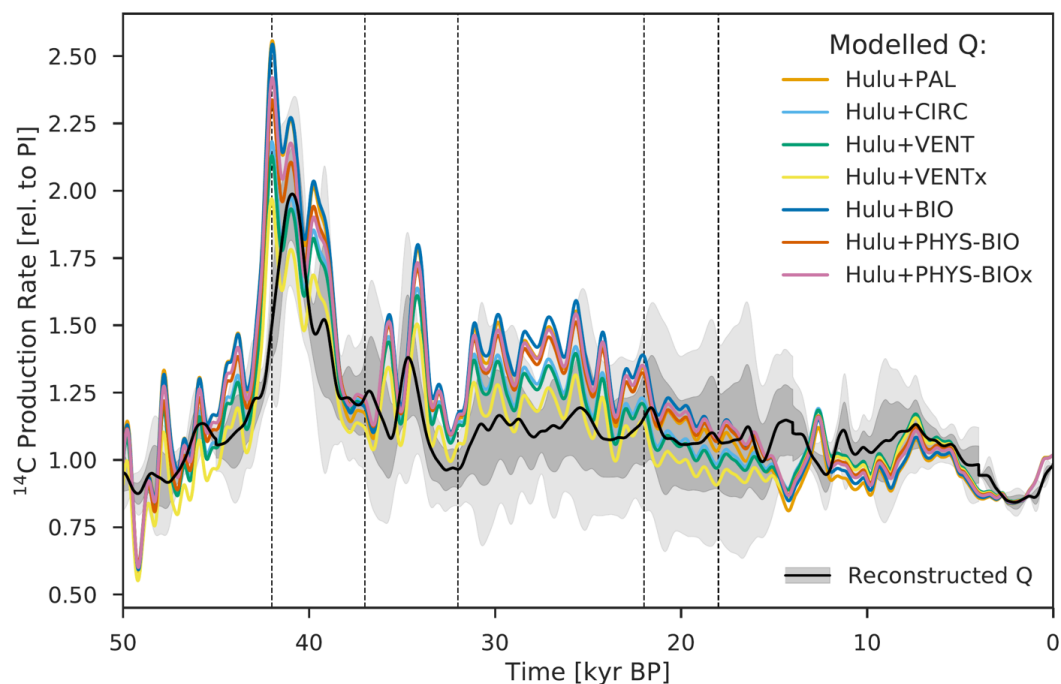
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1268 Fig. 9. Comparison of atmospheric  $\Delta^{14}\text{C}$  variability caused by changes in the ocean carbon cycle (b, d) with  
1269 production-driven changes in atmospheric  $\Delta^{14}\text{C}$  using scenario MOD (a, c). For the analysis of carbon cycle changes,  
1270 only the results of model runs using the mean paleointensity-based  $^{14}\text{C}$  production rate are shown. The  $\Delta^{14}\text{C}$  records  
1271 in the upper panel (a, b) have been detrended to remove the mean, whereas the lower panel (c, d) shows  $\Delta^{14}\text{C}$  anomalies  
1272 expressed as differences relative to the  $\Delta^{14}\text{C}$  value at 50 kyr BP, only for the portion of the record older than 18 kyr  
1273 BP. Two vertical light gray bars indicate the Laschamp (~41 kyr BP) and Mono Lake (~34 kyr BP) geomagnetic  
1274 excursions.

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1282 Fig. 10. Model-based  $^{14}\text{C}$  production rate in relative units compared with paleointensity-based and ice-core  $^{10}\text{Be}$ -based  
1283 estimates. Here, the  $^{14}\text{C}$  production rate is inferred from an atmospheric radiocarbon budget, using the Bern3D carbon  
1284 cycle model forced with reconstructed variations in atmospheric  $\Delta^{14}\text{C}$  and  $\text{CO}_2$ . Results of model runs using the  
1285 IntCal13 calibration curve and seven carbon cycle scenarios are shown in the light blue envelope ( $2\sigma$ ). The light red  
1286 envelope ( $2\sigma$ ) shows the results obtained using the composite Hulu Cave (10.6 to 50 kyr BP) and IntCal13 (0 to 10.6  
1287 kyr BP)  $\Delta^{14}\text{C}$  record. The heavy black line is the mean of the five available production rate reconstructions: Laj et al.  
1288 (2000), Laj et al. (2004), Nowaczyk et al. (2013), Channell et al. (2018), and Adolphi et al. (2018).

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1298 Fig. 11. Relative  $^{14}\text{C}$  production rate as inferred from the Bern3D model under seven carbon cycle scenarios (see Sect.  
1299 2.4). Estimates shown here are obtained by the composite Hulu Cave and IntCal13  $\Delta^{14}\text{C}$  record. The black line is the  
1300 mean of the five available production rate reconstructions shown in Fig. 10; the gray envelope shows its uncertainty  
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