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# Sensitivity of interglacial Greenland temperature and $\delta^{18}\text{O}$ to orbital and $\text{CO}_2$ forcing: climate simulations and ice core data

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## Sensitivity of interglacial Greenland temperature

V. Masson-Delmotte et al.

Title Page

Abstract

Introduction

Conclusions

References

Tables

Figures

⏪

⏩

◀

▶

Back

Close

Full Screen / Esc

Printer-friendly Version

Interactive Discussion

## Abstract

The sensitivity of interglacial Greenland temperature to orbital and CO<sub>2</sub> forcing is investigated using the NorthGRIP ice core data and coupled ocean-atmosphere IPSL-CM4 model simulations. These simulations were conducted in response to different interglacial orbital configurations, and to increased CO<sub>2</sub> concentrations. These different forcings cause very distinct simulated seasonal and latitudinal temperature and water cycle changes, limiting the analogies between the last interglacial and future climate. However, the IPSL-CM4 model shows similar magnitudes of Arctic summer warming and climate feedbacks in response to 2 × CO<sub>2</sub> and orbital forcing of the last interglacial period (126 000 yr ago).

The IPSL model produces a remarkably linear relationship between top of atmosphere incoming summer solar radiation and simulated changes in summer and annual mean central Greenland temperature. This contrasts with the stable isotope record from the Greenland ice cores, showing a multi-millennial lagged response to summer insolation. During the early part of interglacials, the observed lags may be explained by ice sheet-ocean feedbacks linked with changes in ice sheet elevation and the impact of meltwater on ocean circulation, as investigated with sensitivity studies.

A quantitative comparison between ice core data and climate simulations requires to explore the stability of the stable isotope – temperature relationship. Atmospheric simulations including water stable isotopes have been conducted with the LMDZiso model under different boundary conditions. This set of simulations allows to calculate a temporal Greenland isotope-temperature slope (0.3–0.4‰ per °C) during warmer than present Arctic climates, in response to increased CO<sub>2</sub>, increased ocean temperature and orbital forcing. This temporal slope appears twice as small as the modern spatial gradient and is consistent with other ice core estimates. A preliminary comparison with other model results implies that other mechanisms could also play a role. This suggests that further simulations and detailed inter-model comparisons are also likely to be of benefit.

## Sensitivity of interglacial Greenland temperature

V. Masson-Delmotte et al.

Title Page

Abstract

Introduction

Conclusions

References

Tables

Figures



Back

Close

Full Screen / Esc

Printer-friendly Version

Interactive Discussion



Comparisons with Greenland ice core stable isotope data reveals that IPSL/LMDZiso simulations strongly underestimate the amplitude of the ice core signal during the last interglacial, which could reach +8–10 °C at fixed-elevation. While the model-data mismatch may result from missing positive feedbacks (e.g. vegetation), it could also be explained by a reduced elevation of the central Greenland ice sheet surface by 300–400 m.

## 1 Introduction

Greenland ice cores, such as the longest NorthGRIP record (NorthGRIP-community-members, 2004), offer continuous and quantitative archives of past local climate variability at orbital time scale (e.g. Vinther et al., 2009) as well as the evidence for abrupt events (e.g. Capron et al., 2010a). Within uncertainties linked with the conversion of ice core proxies into past temperatures, with age scales, and with glaciological effects (Vinther et al., 2009), ice core data allow to explore the past magnitudes and rates of changes of central Greenland temperature, prior to the instrumental period (Masson-Delmotte et al., 2006b).

In principle, these data can provide a benchmark to test the ability of climate models to correctly represent climate feedbacks (Otto-Bliesner et al., 2006). Past changes in orbital forcing indeed provide natural externally forced experiments on the Earth's climate, leading to past interglacial periods with Arctic temperatures warmer than present-day and large changes in Greenland ice sheet volume (Kopp et al., 2009; Vinther et al., 2009). In particular, the last interglacial period, about 130–120 thousand years before present (ka), was proposed to be a good analogue for future climate change driven by anthropogenic greenhouse gas emissions (Clark and Huybers, 2009; Otto-Bliesner et al., 2006; Sime et al., 2009; Turney and Jones, 2010), especially in the Arctic.

In this manuscript, we address the following questions:

## Sensitivity of interglacial Greenland temperature

V. Masson-Delmotte et al.

Title Page

Abstract

Introduction

Conclusions

References

Tables

Figures

⏪

⏩

◀

▶

Back

Close

Full Screen / Esc

Printer-friendly Version

Interactive Discussion



**Sensitivity of  
interglacial  
Greenland  
temperature**

V. Masson-Delmotte et al.

- 5 – What is the Greenland ice core quantitative information on past surface temperature changes during the current and last interglacial, and how is it related to orbital forcing? This requires to understand the relationship between Greenland surface temperature and snowfall isotopic composition and the various processes that can modify this relationship through time.
- 10 – Which changes in Greenland climate are produced by an ocean-atmosphere model in response to different interglacial orbital configurations? For this purpose, we analyze long snapshot simulations conducted with the IPSL-CM4 model forced only by the orbital configuration of key periods of the current and last interglacial at 0, 6, 9.5, 115 and 126 ka. For 126 ka, we also consider a sensitivity test to a simple parameterization of Greenland ice sheet melt allowing to represent the impact of meltwater on the ocean circulation (Swingedouw et al., 2009).
- 15 – What are the analogies and differences between the climate response to the forcings associated to increased CO<sub>2</sub> concentrations and to changes in orbital configuration? For this purpose, we compare the IPSL-CM4 response to a higher atmospheric CO<sub>2</sub> concentrations and to the last interglacial insolation change, with a focus on Greenland climate. Indeed, climate projections (2 × and 4 × CO<sub>2</sub>) give access to climate states with 3 to 8 °C warmer central Greenland annual mean temperature (Masson-Delmotte et al., 2006b).
- 20 – Is the climate model able to capture the magnitude of changes derived from the ice core data? For direct model-data comparisons, we use the sea surface conditions from the coupled climate model to drive its atmospheric component equipped with the explicit modeling of precipitation isotopic composition (LMDZiso). This also allows to explore the stability of the isotope-temperature change through time and the mechanisms that can alter this relationship.
- 25 – What was the change in central Greenland ice sheet topography during the last interglacial? The IPSL-CM4 and LMDZiso simulations appear to underestimate the

[Title Page](#)[Abstract](#)[Introduction](#)[Conclusions](#)[References](#)[Tables](#)[Figures](#)[⏪](#)[⏩](#)[◀](#)[▶](#)[Back](#)[Close](#)[Full Screen / Esc](#)[Printer-friendly Version](#)[Interactive Discussion](#)

**Sensitivity of interglacial Greenland temperature**

V. Masson-Delmotte et al.

Title Page

Abstract

Introduction

Conclusions

References

Tables

Figures

⏪

⏩

◀

▶

Back

Close

Full Screen / Esc

Printer-friendly Version

Interactive Discussion



Discussion Paper | Discussion Paper | Discussion Paper | Discussion Paper | Discussion Paper

magnitude of last interglacial temperature and precipitation isotopic composition changes compared to the Greenland ice core data. Assuming that the model-data mismatch is mainly caused by a reduced ice sheet elevation, we can estimate the magnitude of this elevation change.

5 Section 2 is dedicated to the information obtained from the NorthGRIP ice core. Section 3 describes the results of the IPSL-CM4 coupled ocean-atmosphere model climate under different orbital configurations; the response of central Greenland climate to orbital forcing is also compared to its response to projections of higher greenhouse gas concentrations. An analysis of the key radiative feedbacks affecting the top of the atmosphere radiative budget is proposed. In Sect. 4, we investigate the Greenland isotope-temperature relationship for warmer than present climates using isotopic atmospheric general circulation models (LMDZiso and HadAM3iso) and discuss the implications for past central Greenland temperature and possibly elevation changes.

**2 Ice core information on past Greenland temperature**

15 **2.1 Water stable isotopes – climate relationships**

Continuous records of water stable isotopes ( $\delta^{18}\text{O}$  or  $\delta\text{D}$ ) have been measured along several deep Greenland ice cores; the longest record so far published was obtained from the NorthGRIP ice core (NorthGRIP-community-members, 2004) (Fig. 1). The initial vapour is formed by evaporation at the ocean surface. Its isotopic composition is affected by evaporation conditions through equilibrium and kinetic fractionation processes, and depends on moisture sources temperature and relative humidity. Along the air mass trajectories to Greenland, the isotopic composition of the atmospheric water vapour undergoes mixing by convection, upload of new water vapor from different sources, and distillation linked with the progressive air mass cooling and successive condensation, as well as kinetic effects on ice crystals. Altogether, these physical processes result in a linear relationship between the air temperature and the snowfall

isotopic composition in central Greenland. The slope of the modern spatial relationship is 0.7‰ of  $\delta^{18}\text{O}$  per °C for the first ice core sites (Dansgaard, 1964) and 0.8‰ of  $\delta^{18}\text{O}$  per °C for all available data including coastal stations (Dansgaard, 1964; Sjolte et al., 2011).

In addition to the impact of condensation temperature, several effects can affect the precipitation isotopic composition and modify the temporal isotope-temperature relationship

– *deposition effects*, caused by precipitation intermittency or changes in the relationship between the temperature at the condensation level and the surface temperature (Jouzel et al., 1997). Atmospheric models have shown a large deposition effect for glacial climate, due to strongly reduced winter precipitation (Krinner et al., 1997; Werner et al., 2000). In this manuscript, we assess the “precipitation weighting effect” by comparing the average temperature change to the monthly precipitation weighted temperature change;

– *source effects*, caused by changes in evaporation conditions or moisture origin (Johnsen et al., 1989; Jouzel et al., 2007; Masson-Delmotte et al., 2005a, b);

– *glaciological effects*, caused by changes in ice sheet topography which affects surface air temperature and stable isotopic composition (Vinther et al., 2009). We therefore introduce the notion of temperature estimate “at fixed elevation”, by contrast with the information on air temperature at the ice sheet surface classically derived from stable isotope data.

Alternative information on past Greenland temperature is available from the borehole temperature profiles (Dahl-Jensen et al., 1998) and from firn gas fractionation during abrupt warmings (Capron et al., 2010a; Severinghaus et al., 1998). The latter method allows the estimation of the interstadial isotope-temperature slope to range between  $0.30 \pm 0.05$  and  $0.60 \pm 0.05$ ‰ per °C (Capron et al., 2010a), therefore quite different from the spatial slope. This probably results from deposition and source effects

## Sensitivity of interglacial Greenland temperature

V. Masson-Delmotte et al.

Title Page

Abstract

Introduction

Conclusions

References

Tables

Figures

⏪

⏩

◀

▶

Back

Close

Full Screen / Esc

Printer-friendly Version

Interactive Discussion



(Masson-Delmotte et al., 2005a). In Sect. 4, we will use isotopic simulations to quantify the isotope-temperature relationship in warmer than present climate conditions.

## 2.2 Greenland Holocene climate and ice sheet elevation

Recently, (Vinther et al., 2009) conducted a synthesis of the Greenland ice core Holocene stable isotope information has recently been conducted. It combines ice core records from coastal ice caps (where changes in elevation are limited) and from the central ice sheet (where elevation changes can significantly affect the isotopic signals). The authors extract a common and homogeneous annual mean Greenland temperature signal, together with regional changes in the ice sheet topography. The new temperature history from this study (Fig. 1, central panel, blue line) reveals a pronounced Holocene climatic optimum in Greenland coinciding with a maximum thinning near the ice sheet margins. These results also imply that the NorthGRIP ice core  $\delta^{18}\text{O}$  data can be converted to temperature with a temporal slope of 0.45‰ per °C.

They calculate that the elevation of the NorthGRIP site has decreased by ~140 m since 9.5 ky and by ~60 m from 6 ka to present. The central Greenland temperature “at fixed elevation” is estimated to be ~2.3 °C higher at 9.5 ka and ~2.0 °C at 6 ka than during the last millennium, with a multi-millennial warm plateau encountered between 9.3 and 6.8 ka. This plateau occurs 1.8 to 4.3 ky (thousand years) later than the maximum in 75° N June insolation. The early Holocene warmth is partly masked in the central Greenland ice core stable isotope records because of the larger volume and elevation of the ice sheet.

## 2.3 Links between NorthGRIP $\delta^{18}\text{O}$ and 75° N summer insolation

We extract the orbital components of the NorthGRIP record using the first components of a Singular Spectrum Analysis performed on the whole series, and corresponding to periodicities longer than 3 ky (Fig. 1, bold line, central panel). With the available ice core age scales (Capron et al., 2010b; Svensson et al., 2008), the orbital component of

### Sensitivity of interglacial Greenland temperature

V. Masson-Delmotte et al.

Title Page

Abstract

Introduction

Conclusions

References

Tables

Figures

⏪

⏩

◀

▶

Back

Close

Full Screen / Esc

Printer-friendly Version

Interactive Discussion



the NorthGRIP  $\delta^{18}\text{O}$  appears to lag the reversed precession parameter (in phase with local June insolation) by several millennia (Fig. 1). A significant correlation ( $R^2 = 0.27$ ) is obtained between the smoothed NorthGRIP  $\delta^{18}\text{O}$  and 4 ky earlier  $75^\circ\text{N}$  June insolation. The four most recent optima in this smoothed NorthGRIP  $\delta^{18}\text{O}$  record lag maxima in  $75^\circ\text{N}$  June insolation by respectively 4.8, 4.8, 3.1 and 3.5 ky (Fig. 1, dashed vertical lines). These lags are significantly larger than the GICC05 age scale uncertainty ( $\sim 80$  yr at 10 ka,  $\sim 440$  yr at 20 ka,  $\sim 1000$  yr at 30 ka and  $\sim 2600$  yr at 60 ka) (Rasmussen et al., 2006; Svensson et al., 2008) and occur both under glacial and interglacial contexts.

For the Holocene, it is obvious that the Greenland optimum (at  $\sim 7$ – $10$  ka) occurs later than the 11 ka precession minimum (local June insolation maximum), likely because of the negative feedback linked with the Laurentide ice sheet albedo and weaker northward advection of heat in the Atlantic Ocean caused by the meltwater from deglaciating Northern Hemisphere ice sheets (Renssen et al., 2009). The NorthGRIP record does not allow to explore this aspect for the last interglacial because it does not span the whole length of this period (NorthGRIP-community-members, 2004). Marine sediment records of North Atlantic sea surface temperature suggest a pattern similar to the Holocene with a lag between peak insolation and peak isotopic values (Masson-Delmotte et al., 2010a). During the end of interglacials (after optima in insolation and in  $\delta^{18}\text{O}$ ), parallel decreasing trends in  $75^\circ\text{N}$  June insolation and NorthGRIP  $\delta^{18}\text{O}$  are observed. For the mid to late Holocene (the last 8 ky), the  $\delta^{18}\text{O}$ -insolation slope is  $0.02\text{‰ per }^\circ\text{W m}^{-2}$  ( $0.03$  to  $0.06^\circ\text{C per W m}^{-2}$ ), much weaker than for the end of the last interglacial (121 to 115 ka), where it reaches  $0.10\text{‰ per }^\circ\text{W m}^{-2}$  ( $\sim 0.17$  to  $0.33^\circ\text{C per W m}^{-2}$ ).

## 2.4 Greenland last interglacial climate

The ice core information on central Greenland climate during the last interglacial is not as precise as for the Holocene due to the age scale uncertainty, the end of the

### Sensitivity of interglacial Greenland temperature

V. Masson-Delmotte et al.

Title Page

Abstract

Introduction

Conclusions

References

Tables

Figures

⏪

⏩

◀

▶

Back

Close

Full Screen / Esc

Printer-friendly Version

Interactive Discussion





NorthGRIP record at  $\sim 122$  ka, and the lack of information from borehole thermometry to constrain the isotope-temperature-elevation histories. Based on the shape of north Atlantic SST records synchronized on the EDC3 age scale (Masson-Delmotte et al., 2010a), one may assume that the isotopic values of the deepest part of the NorthGRIP ice core may be representative of a multi-millennial temperature plateau. Considering the uncertainty on the isotope-temperature relationship (between 0.3 and 0.8 ‰ per °C), the NorthGRIP Last Interglacial  $\sim 3$ ‰  $\delta^{18}\text{O}$  anomaly would translate into a 3.8–10.0 °C surface temperature anomaly. The signal for the last interglacial is without doubt larger than for the early to mid Holocene (see Sect. 2.2), as also expected from the larger orbital forcing (Fig. 1).

The data do not allow by themselves to quantify the deposition or glaciological effects affecting this temperature estimate, motivating the use of climate models to explore the mechanisms controlling precipitation isotopic composition.

### 3 Climate modelling

#### 3.1 IPSL Coupled climate model simulations

The IPSL-CM4 coupled climate model has been intensively used for CMIP3 and PMIP2 simulations (Alkama et al., 2008; Born et al., 2010; Braconnot et al., 2007, 2008a; Kageyama et al., 2009; Marti et al., 2010; Swingedouw et al., 2006). The model couples the atmospheric component LMDZ (Hourdin et al., 2006) with the OPA ocean component (Madec and Imbard, 1996). A sea ice model (Fichefet and Maqueda, 1997) which computes the ice thermodynamics and physics is coupled with the ocean-atmosphere model. The ocean and atmosphere exchange momentum, heat and fresh-water fluxes, as well as surface temperature and sea ice once a day, using the OASIS coupler (Valcke, 2006). None of the fluxes are corrected or adjusted. The model is run with a horizontal resolution of 96 points in longitude and 71 points in latitude ( $3.78^\circ \times 2.58^\circ$ ) for the atmosphere and 182 points in longitude and 149 points in latitude for the ocean. There are 19 vertical levels in the atmosphere and 31 levels in the

## Sensitivity of interglacial Greenland temperature

V. Masson-Delmotte et al.

Title Page

Abstract

Introduction

Conclusions

References

Tables

Figures

⏪

⏩

◀

▶

Back

Close

Full Screen / Esc

Printer-friendly Version

Interactive Discussion





thermohaline circulation (Swingedouw et al., 2009). The model response has also been analysed in monsoon areas (Braconnot et al., 2008b).

The IPSL-CM4 model results have previously been compared with the ice core information and other model results in terms of polar amplification under glacial conditions or climate projection scenarios (Masson-Delmotte et al., 2006a, b) as well as briefly for the last interglacial (Masson-Delmotte et al., 2010b). These previous studies showed that the IPSL-CM4 model response is comparable to other climate models and generally seems to underestimate the magnitude of temperature changes compared to those derived from the ice core data.

### 3.2 Impact of orbital forcing on IPSL simulated central Greenland climate

Figure 2 displays the model results for central Greenland using the same definition as in Masson-Delmotte et al. (2006b), that is the temperature averaged at places where ice sheet elevation is above 1.3 km. For each simulation, monthly mean values of central Greenland temperatures are displayed as a function of monthly mean values of 75° N top of atmosphere incoming solar radiation. The elliptic shape of the plots reflects the one month seasonal lag between surface air temperature and insolation, mostly because of the thermal inertia of the surrounding oceans affecting heat advection to central Greenland. Orbital forcing alone has limited impacts on the simulated winter temperature (because of a weak incoming insolation at that season and latitude) and a strong impact on summer-fall temperatures.

Because the change in summer temperature (with a range of July temperature changes from  $-2.5^{\circ}\text{C}$  for 115 ka to  $+5.8^{\circ}\text{C}$  for 126 ka) is dominating the annual mean temperature change (Table 2), the IPSL model simulates the same sign for annual mean and summer temperature changes but a weaker annual mean temperature change ( $-0.5$  for 115 ka to  $+0.9^{\circ}\text{C}$  for 126 ka). The model results for summer and annual mean temperature are depicted in Fig. 1 with respectively red and green open circles. This comparison suggests that the IPSL model simulation has the right sign of temperature changes, but underestimates the magnitude of annual mean changes

## Sensitivity of interglacial Greenland temperature

V. Masson-Delmotte et al.

Title Page

Abstract

Introduction

Conclusions

References

Tables

Figures

⏪

⏩

◀

▶

Back

Close

Full Screen / Esc

Printer-friendly Version

Interactive Discussion



compared to the ice core derived information. We now explore the simulated deposition effects which can impact the model-data comparison, focusing on the precipitation weighting effect.

For all orbital contexts, the IPSL model shows a positive precipitation weighting effect (difference between monthly precipitation weighted temperature and annual mean temperature) (Table 2, last column). This effect is minimum at 115 ka (1.8 °C), maximum at 126 ka (5.2 °C) and is strongly enhanced with increasing local summer insolation. This is due to a strong (non linear) enhancement of summer precipitation for warmer summer temperature (Table 2). The IPSL model therefore points to a large deposition effect, suggesting that the Greenland ice core warm interglacial proxy records (stable isotopes, but also <sup>10</sup>Be... ) may be biased towards summer. The simulated changes in precipitation weighted temperature is intermediate between the summer and annual mean temperature, and vary between -1.1 °C (at 115 ka) and +3.6 °C (at 126 ka) (Table 2).

In the IPSL simulations, the maximum summer temperature change (occurring in July) appears to be strongly linearly related ( $R^2 = 0.99$ ) with maximum 75° N incoming summer insolation (occurring in June), with a slope of 0.08 °C per  $W m^{-2}$  (Fig. 2b). We first observe that, even considering this largest signal (July temperature), the model response to summer insolation therefore appears at least twice as small as derived from the ice core data for the transition from 122 to 115 ka (0.17 to 0.33 °C per  $W m^{-2}$ , see Sect. 2.3). In Sect. 3.2, we will investigate the changes affecting the top of the atmosphere radiative budget and key radiative feedbacks in order to better describe the processes responsible for such a linear model response to the orbital forcing.

When taking into account the ocean circulation changes linked with a parameterization of Greenland melt at 126 ka, the IPSL model simulates a 0.6 °C weaker July (resp. 0.4 °C annual) warming than in the standard 126 ka simulation (not shown in Table 2). In this simulation, the AMOC is reduced because deep water formation in the North Atlantic / Nordic Seas is weakened by the Greenland ice sheet meltwater. The meridional transport by the atmospheric circulation is enhanced to compensate for the reduction

## Sensitivity of interglacial Greenland temperature

V. Masson-Delmotte et al.

[Title Page](#)[Abstract](#)[Introduction](#)[Conclusions](#)[References](#)[Tables](#)[Figures](#)[⏪](#)[⏩](#)[◀](#)[▶](#)[Back](#)[Close](#)[Full Screen / Esc](#)[Printer-friendly Version](#)[Interactive Discussion](#)

in ocean heat transport but the Arctic cools because of a larger sea ice extent. Taking into account the impact of ice sheet melt on the ocean circulation therefore increases the model-data mismatch.

### 3.3 Differences between increased CO<sub>2</sub> and orbitally forced IPSL climate responses

The orbital forcing has a negligible impact as such on the global and annual radiative forcing ( $<0.3 \text{ W m}^{-2}$  over the last 130 ka), which contrasts with the  $3.7 \text{ W m}^{-2}$  radiative forcing for  $2 \times \text{CO}_2$  (resp.  $7.4 \text{ W m}^{-2}$  for  $4 \times \text{CO}_2$ ). Moreover, the diurnal and seasonal distributions of these two forcings are drastically different. Anomalies in summer insolation exceed  $50 \text{ W m}^{-2}$  at mid and high northern latitudes (Fig. 3ab, showing top of atmosphere radiative budgets) at 126 ka, with large seasonal and latitudinal contrasts. This differs from the more homogeneous forcing caused by increased CO<sub>2</sub> concentrations. Note that obliquity affects the latitudinal distribution of annual insolation, with opposite effects at low and high latitudes, and a range of variations of resp.  $4.5$  to  $10.5 \text{ W m}^{-2}$  at  $75^\circ \text{ N}$  along the current and last interglacial (0–12 ka and 115–130 ka).

We now focus on the 126 ka simulation, because of the large magnitude of the seasonal insolation change caused by the combination of precession and eccentricity for this period, and compare it with the  $2 \times \text{CO}_2$  simulation. Figure 3 (panels c and d) shows the differences between last interglacial (126 ka) and present day for JJA, DJF and annual mean temperature, as well as their zonal mean, and compares them to the differences between  $2 \times \text{CO}_2$  and present day for JJA, DJF and annual mean temperature. Increased CO<sub>2</sub> leads to simulated warming at low latitudes and a larger magnitude of warming at both poles (with reference to present day reference simulations), especially in the winter season. By contrast the 126 ka orbital forcing leads to a small annual mean cooling at low to mid latitudes, a small annual mean warming anomaly around  $60^\circ \text{ S}$  and a large ( $\sim 4^\circ \text{ C}$ ) warming in the Arctic. The model response to 126 ka orbital forcing follows the latitudinal and seasonal anomalies of insolation (Fig. 3a and b), with the exception of the Arctic, where a year round persistent warming is simu-

## Sensitivity of interglacial Greenland temperature

V. Masson-Delmotte et al.

Title Page

Abstract

Introduction

Conclusions

References

Tables

Figures

⏪

⏩

◀

▶

Back

Close

Full Screen / Esc

Printer-friendly Version

Interactive Discussion



lated. Such a feature is model-dependent, as shown by the comparison between the IPSL results and other coupled model simulations for the seasonal cycle of simulated last interglacial temperature anomalies for Greenland (Masson-Delmotte et al., 2010b).

Figure 4 shows that the control simulation correctly captures the amplitude and extrema of the observed Northern Hemisphere sea ice cover (Rayner et al., 2003), but has a slight shift (one month earlier than in the data) in the seasonal cycle. In response to 126 ka orbital forcing, the model produces a summer sea ice retreat of  $\sim 3$  million  $\text{km}^2$ . This represents half of the retreat ( $\sim 6$  million  $\text{km}^2$ ) simulated for  $2 \times \text{CO}_2$ . In the 126 ka simulation, a small winter sea ice retreat is also simulated. This can be attributed to the large uptake of heat during summer in the high latitude ocean as well as to an enhanced AMOC, which brings warm surface waters to the high latitudes (Born et al., 2010). This winter sea ice retreat is probably the cause for the warmer winter temperatures at 126 ka compared to the control simulation (Fig. 3).

Winter Arctic warming is particularly large under  $2 \times \text{CO}_2$  forcing, reaching  $\sim 8^\circ\text{C}$ , to be compared to the  $\sim 2^\circ\text{C}$  Arctic warming for 126 ka conditions. While this comparison highlights the differences between the two types of simulations, and therefore the limitations of analogies between the last interglacial and future climate change, we would like to stress that the simulated summer Arctic warming at 126 ka reaches a magnitude ( $\sim 4^\circ\text{C}$ ) comparable to summer Arctic warming forced by  $2 \times \text{CO}_2$  (see also Fig. 5d).

The different climate responses to orbital (126 ka) and  $2 \times \text{CO}_2$  forcing also have a signature on patterns of evaporation changes. Figure 3 shows a strong increase in north Atlantic evaporation at 126 ka, in contrast with a strong increase in Nordic Seas evaporation in response to  $2 \times \text{CO}_2$  forcing (probably linked with reduced sea ice cover). We expect that changes in moisture sources affect moisture distillation and Greenland precipitation isotopic depletion, and therefore the isotope-temperature relationships. Before presenting the isotopic calculation results (in Sect. 4), we perform a simple analysis of radiative feedbacks in order to understand the causes for the linear behavior of the IPSL-CM4 model in response to orbital forcing, and to further compare the model response to  $\text{CO}_2$  and orbital forcing.

## Sensitivity of interglacial Greenland temperature

V. Masson-Delmotte et al.

[Title Page](#)[Abstract](#)[Introduction](#)[Conclusions](#)[References](#)[Tables](#)[Figures](#)[⏪](#)[⏩](#)[◀](#)[▶](#)[Back](#)[Close](#)[Full Screen / Esc](#)[Printer-friendly Version](#)[Interactive Discussion](#)

### 3.4 Analysis of radiative feedbacks

Following (Braconnot et al., 2007), a simple feedback analysis was performed in order to quantify the main drivers of changes in the top of the atmosphere radiative budget (TOA) at high latitudes (60–80° N). The methodology for this analysis is described in Appendix A.

Figure 5 displays analyses of the top of atmosphere radiative budget terms and key feedbacks (represented by symbols) around Greenland. The different simulations are represented by the same colors as in Fig. 2. The specific radiative budgets for June–July are shown for all the orbitally forced simulations (Fig. 5a) and for each month for 126 ka (Fig. 5b) and  $2 \times \text{CO}_2$  simulations (Fig. 5c). We do not display the changes in heat and water transport and only focus on the local radiation fluxes within the atmospheric column.

Figure 5a allows to better characterise the radiative feedbacks involved in the linear response of the IPSL simulated summer Greenland surface temperature with respect to summer insolation. At high northern latitudes, the different components of the radiative budget depict a linear relationship with respect to the change in incoming solar radiation at the top of the atmosphere  $\Delta \text{SWi}_{\text{simul}}$ . The net top of atmosphere shortwave flux ( $\Delta \text{SWi}_{\text{simul}}$ , represented by “x” symbols) appears relatively close to the prescribed insolation change and only partially compensated for by increased longwave emission ( $\Delta \text{LW}_{\text{simul}}$ , represented by filled rectangles) so that the net radiative budget is positive (not shown).

At 6, 9.5, 122 and 126 ka, a strong positive shortwave feedback is linked with the total (surface and cloud) albedo effect ( $\Delta \text{albedo}_{\text{simul}}$ , represented by “+” symbols). This effect is dominated by the clear sky (surface) albedo effect ( $\Delta \text{albedo}_{\text{cs}_{\text{simul}}}$  represented by the triangle symbols), only partly compensated by an enhanced negative cloud shortwave feedback (difference between  $\Delta \text{albedo}_{\text{cs}_{\text{simul}}}$  and  $\Delta \text{albedo}_{\text{simul}}$ ). The albedo feedback is consistent with changes in sea ice (Fig. 4). It increases almost linearly with the insolation forcing, stressing that the changes in clear sky shortwave

## Sensitivity of interglacial Greenland temperature

V. Masson-Delmotte et al.

Title Page

Abstract

Introduction

Conclusions

References

Tables

Figures

⏪

⏩

◀

▶

Back

Close

Full Screen / Esc

Printer-friendly Version

Interactive Discussion



## Sensitivity of interglacial Greenland temperature

V. Masson-Delmotte et al.

Title Page

Abstract

Introduction

Conclusions

References

Tables

Figures

⏪

⏩

◀

▶

Back

Close

Full Screen / Esc

Printer-friendly Version

Interactive Discussion

surface radiation drive the surface radiative budget, surface temperature and thereby the snow and ice extent. Note that by construction, the total albedo feedback between the different simulations lies on a line proportional to the planetary albedo of the control simulation. At 115 ka, clear sky and cloud albedo feedbacks have opposite signs and have a much smaller magnitude (with respect to the magnitude of the orbital forcing) compared to other orbital simulations. The different effects are thus not symmetrical for increased or reduced insolation, certainly due to the temperature thresholds needed to build and melt snow and ice.

In addition, the longwave radiative budget changes ( $\Delta LWn_{\text{simul}}$ , filled diamonds) appear to be driven by the changes in Planck emission directly caused by changes in surface temperature ( $\Delta PI_{\text{simul}}$ , open diamonds). There is only a small increase in the atmospheric greenhouse effect caused by changes in the vertical temperature profile, water vapour content, and infra-red cloud radiative feedbacks (difference between the filled and open diamond symbols). This greenhouse feedback is too small to drive a non linear response of the radiative budget around Greenland.

While this approach ignores the dynamical heat advection effects, it suggests that the top of the atmosphere radiative budget at high northern latitude is relatively linear with respect to orbital forcing and highlights the importance of the positive feedbacks linked with the surface albedo. The magnitude of the atmospheric greenhouse effect and the shortwave cloud negative feedback increase with the magnitude of the insolation forcing. In this model, the cloud feedback is enhanced in a warmer Arctic. Compensations of non linearities of the Planck, albedo and cloud radiative effects at 115 ka likely explain the overall linearity of the IPSL model high northern latitude temperature response to summer insolation forcing.

Figure 5b and c allow to compare the seasonal cycle and magnitude feedbacks at play in 126 ka and  $2 \times CO_2$  simulations, which reach similar magnitudes of summer temperature change in Greenland. As expected, the changes in greenhouse effect are larger for the  $2 \times CO_2$  forcing than for insolation forcing. The net shortwave radiation is the same here as the total albedo effects and is also positive in summer, as is the



net radiative budget. Again, the clear sky albedo feedback is the dominant contribution during summer, and the cloud feedback only accounts for a small fraction of changes in shortwave radiation, even though its magnitude is larger than for the insolation forcing. This comparison shows that the albedo, cloud and atmospheric greenhouse feedbacks have comparable magnitude and sign in summer. These simulated feedbacks seem consistent with ongoing changes related with Arctic sea ice retreat and warming (Screen and Simmonds, 2010). It also highlights the different seasonality effects, with larger winter greenhouse feedbacks for  $2 \times \text{CO}_2$  in winter, as well as an earlier albedo feedbacks for  $2 \times \text{CO}_2$  likely caused by the strongly reduced winter sea ice cover in this simulation than for 126 ka (Fig. 4).

In order to better characterize the links between changes in surface temperature and atmospheric water content, Fig. 5d compares the seasonal cycle of atmospheric precipitable water anomaly as a function of surface temperature anomaly for 115, 126 ka and  $2 \times \text{CO}_2$  simulations. The asymmetry between atmospheric moisture changes at 115 and 126 ka is obvious. Despite a completely different seasonality of the changes (with the  $2 \times \text{CO}_2$  simulations showing its largest temperature changes in winter), the 126 ka and  $2 \times \text{CO}_2$  simulations again depict similar magnitudes of temperature and precipitable water changes, in summer.

## 4 Atmospheric modeling of water stable isotopes

### 4.1 Set up of the LMDZiso simulations

While water stable isotopes are not yet available in the coupled IPSL model, they have been implemented in its atmospheric component, LMDZ4 (Risi et al., 2010b), with a standard resolution of  $2.5^\circ \times 3.75^\circ$ . The ability of the model to capture the modern and LGM Greenland precipitation isotopic composition has already been analysed (Risi et al., 2010b; Steen-Larsen et al., 2011). These comparisons have shown that the model correctly captures the 0.8‰ per  $^\circ\text{C}$  modern spatial isotope-temperature relationship. In

## Sensitivity of interglacial Greenland temperature

V. Masson-Delmotte et al.

Title Page

Abstract

Introduction

Conclusions

References

Tables

Figures

◀

▶

◀

▶

Back

Close

Full Screen / Esc

Printer-friendly Version

Interactive Discussion



central and North Greenland, the model has a warm bias (up to 8 °C) and produces too depleted precipitation (by 5 ‰). This contrasts with a cold and enriched bias at coastal stations. Comparable biases are found by other atmospheric models (e.g. ECHAM and REMO-iso) (Sjolte et al., 2011).

A suite of simulations has been conducted with the LMDZ4iso model, forced by the sea surface conditions and associated external forcings (6 ka or 126 ka orbital parameters, or increased greenhouse gas concentration) simulated by the IPSL model; a sensitivity test with 4 °C homogeneous artificial increase in sea surface temperature compared to present-day (AMIP) has also been performed (Table 1).

## 4.2 LMDZiso isotope-temperature relationships

Consistent with the coupled IPSL model simulations, annual mean temperature changes simulated in central Greenland remain very small for the different simulations corresponding to changes in orbital configurations (<1 °C); they reach 4 °C for 2×CO<sub>2</sub>, 6 °C for SST+4 °C and ~9 °C for 4xCO<sub>2</sub> simulations (Fig. 7). Deposition effects can be considered both for temperature and δ<sup>18</sup>O by calculating either annual mean or precipitation weighted values (Fig. 7). As discussed previously, this effect is particularly large for the orbitally forced simulations (up to 2 °C and 1 ‰ reaching magnitudes comparable to the climate change signal). Because the CO<sub>2</sub> forcing increases both winter and summer temperature and precipitation (Fig. 6), the resulting precipitation weighting effect is smaller (typically 1 °C and 0.5 ‰ for 4 × CO<sub>2</sub>). This effect enhances the magnitude of precipitation weighted δ<sup>18</sup>O anomalies (Fig. 7) and therefore slightly increases the “warm climate” isotope-temperature slope (from 0.30 to 0.36 ‰ per °C). Within all the studied simulations, the strength of the correlation is comparable between annual mean precipitation isotopic composition and temperature, and precipitation weighted isotopic composition and temperature ( $R^2 > 0.95$ ,  $n = 6$ ) and significantly larger than the correlation between precipitation weighted isotopic signal and annual mean temperature ( $R^2 = 0.86$ ,  $n = 6$ ). This suggests that the ice core data

## Sensitivity of interglacial Greenland temperature

V. Masson-Delmotte et al.

Title Page

Abstract

Introduction

Conclusions

References

Tables

Figures

⏪

⏩

◀

▶

Back

Close

Full Screen / Esc

Printer-friendly Version

Interactive Discussion

(capturing precipitation weighted information) should best be interpreted in terms of changes in precipitation weighted temperature.

When considering all the available simulations, a linear regression leads to a mean “warm climate” isotope-temperature slope of 0.31 ‰ per °C, with values ranging from 0.26 to 0.39 ‰ per °C. This uncertainty is estimated by using either annual mean or precipitation weighting for temperature and  $\delta^{18}\text{O}$ , and by selections of 5 of the 6 simulations to assess the uncertainty on each slope, which is about 0.03 ‰ per °C. This simulated slope is consistent with the lowest values derived from interstadial warming events (Capron et al., 2010a), with the slopes obtained using the borehole information at the glacial-interglacial scale (Cuffey and Clow, 1997; Dahl-Jensen et al., 1998), and lower than the slopes estimated during the current interglacial period after accounting for elevation changes (Vinther et al., 2009). This finding is also consistent with a small isotope-temperature slope simulated by the GISS model for the Holocene for Greenland (Legrande and Schmidt, 2009).

At 126 ka, the simulated change in Greenland precipitation isotopic composition is very small (0.75 ‰) compared to the ice core data. Indeed, a ~3 ‰ anomaly above the last millennium level is consistently recorded in the deepest part of the NorthGRIP ice core (at 123 ka), in Eemian ice found in the disturbed bottom layers at Summit (Landais et al., 2004; Masson-Delmotte et al., 2010a; Suwa et al., 2006) and in preliminary measurements from the NEEM ice core, recently drilled in north west Greenland (unpublished data).

The deposition effect alone cannot explain why the isotope-temperature slope is particularly weak for these warmer than present climates. Larger spring-summer temperature anomalies in the  $4 \times \text{CO}_2$  simulation are only associated with a small Greenland precipitation  $\delta^{18}\text{O}$  anomaly. This is also the case, but in a weaker proportion, for the 126 ka simulation. Source effects linked with geographical shifts of the origin of the moisture source (as hinted by changes in evaporation, Fig. 3) are likely the cause for a reduced isotopic depletion despite strong summer Arctic warming.

## Sensitivity of interglacial Greenland temperature

V. Masson-Delmotte et al.

Title Page

Abstract

Introduction

Conclusions

References

Tables

Figures

⏪

⏩

◀

▶

Back

Close

Full Screen / Esc

Printer-friendly Version

Interactive Discussion



### 4.3 LMDZiso changes in moisture origin

We conducted a water tagging experiment (Risi et al., 2010a) in which the high latitude (North of 50° N) oceanic evaporation was tagged for the control and 4 × CO<sub>2</sub> experiments (in order to explore the largest anomaly). For central Greenland, 14 % of present-day moisture originates from high latitude (>50° N) evaporation. High latitude moisture is strongly isotopically enriched compared to the global mean atmospheric water vapour. The modern spatial slope in Greenland is 0.8‰ per °C including all moisture sources. The water tagging simulation allows to show that, without the Arctic moisture source, this spatial slope would be reduced to 0.7‰ per °C. This arises from a spatial gradient in the contribution of (enriched) high latitude moisture to Greenland precipitation. This contribution decreases poleward, because air mass trajectories reaching northern Greenland are transported at high elevation and are less exposed to high latitude evaporation.

In the 4 × CO<sub>2</sub> experiment, the proportion of high latitude moisture decreases by about 40 % in winter and 60 % in summer, due to enhanced poleward moisture transport from the subtropics and decreased high latitude evaporation (Fig. 3). This source effect quantitatively explains the difference between the Rayleigh isotope-temperature slope (0.7‰ per °C) and the actual temporal isotope-temperature slope (0.3‰ per °C). This analysis shows that changes in high latitude recycling explain why the isotope-temperature slopes for warmer climates is much smaller in LMDZiso than the modern spatial slope. We now compare the LMDZiso model results with other available isotopic model results.

### 4.4 Comparison with other isotope model results

Small slopes are simulated by the LMDZiso model for Greenland for projections and interglacial configurations, and by the GISS model for the Holocene for Greenland (Legrande and Schmidt, 2009).

Here we also briefly examine results from Greenland using HadAM3iso simulations

CPD

7, 1585–1630, 2011

## Sensitivity of interglacial Greenland temperature

V. Masson-Delmotte et al.

Title Page

Abstract

Introduction

Conclusions

References

Tables

Figures

⏪

⏩

◀

▶

Back

Close

Full Screen / Esc

Printer-friendly Version

Interactive Discussion



previously published for Antarctica (Sime et al., 2008). We focus on a snapshot simulation for year 2100 in response to SST and sea ice outputs from the coupled Hadley model simulation using the A1B greenhouse concentration scenario. This is relatively comparable to the LMDZiso 2 × CO<sub>2</sub> simulation.

5 Whilst seasonal cycles of LMDZiso 2 × CO<sub>2</sub> and HadAM3iso 2100 outputs show relatively comparable magnitudes of Arctic sea ice, central Greenland temperature and precipitation changes, albeit with slightly different seasonal aspects (Appendix Fig. A1),  $\delta^{18}\text{O}$  anomalies (with respect to the reference period) are higher for HadAM3iso (not shown). The HadAM3iso  $\delta^{18}\text{O}$  anomalies are positive all year round, while LMDZiso  
10 2 × CO<sub>2</sub> shows very small (or slightly negative)  $\delta^{18}\text{O}$  anomalies for that season. As a result, the HadAM3iso model produces larger shifts in  $\delta^{18}\text{O}$  for a comparable warming, compared with LMDZiso. The average central Greenland shift is about 3‰ in HadAM3iso, which is slightly closer to the observed interglacial shift, compared with LMDZiso. However, note that since this shift occurs due to CO<sub>2</sub> forcing, rather than  
15 a more realistic orbital forced warming, so it is difficult to know the pertinence of this result for the last interglacial climate.

The difference between the models likely arises from differences in moisture advection to central Greenland in the two models. HadAM3iso 2100 evaporation changes have comparable patterns but larger magnitudes at high northern latitudes, compared  
20 to 2 × CO<sub>2</sub> LMD4iso results (Figs. 3 and S1). This suggests that, whilst LMDZ4iso enhances the transport of depleted subtropical moisture towards Greenland (see previous section), the specific 2100 simulation examined here may be allowing HadAM3iso to transport more moisture from nearby sea ice free high latitude oceans during the CO<sub>2</sub> warming. Present day observations also depict shifts between local and advected  
25 moisture during the autumn ice growth season with distinct isotopic fingerprints which also tends to support the idea that this local-distal moisture transport balance mechanism could be important (Kurita, 2011).

We conclude from these sections that changes in deposition (bias towards summer precipitation for orbitally driven warm climates) and source effects (varying contribution

## Sensitivity of interglacial Greenland temperature

V. Masson-Delmotte et al.

[Title Page](#)[Abstract](#)[Introduction](#)[Conclusions](#)[References](#)[Tables](#)[Figures](#)[⏪](#)[⏩](#)[◀](#)[▶](#)[Back](#)[Close](#)[Full Screen / Esc](#)[Printer-friendly Version](#)[Interactive Discussion](#)

of Arctic moisture for all simulations) are responsible for the LMDZiso Greenland isotope-temperature slope smaller than the modern spatial slope for warmer than present climates. The magnitude of changes in moisture origins and transport pathways could affect the isotope-temperature slope between different models and different simulations. Additional investigations are needed to assess better understand the reasons for inter-model differences.

#### 4.5 Implications of IPSL/LMDZiso results for central Greenland ice sheet elevation during the last interglacial

The LMDZiso low temporal slope appears consistent with previous results obtained for glacial climate (Capron et al., 2010a) and Holocene climate (Vinther et al., 2009). The IPSL-CM4 and LMDZiso models do underestimate the magnitude of temperature and precipitation isotopic composition changes compared to the ice core data. This mismatch may result from either missing feedbacks (e.g. vegetation changes), model sensitivity to forcings (e.g. magnitude of sea ice, water vapour and moisture origin, cloud etc. feedbacks), or, alternatively, from changes in Greenland elevation which are not considered in the climate simulations.

Assuming that the LMDZ/IPSL model correctly captures the first order of the response to 126 ka insolation, the model-data comparison leaves a  $\delta^{18}\text{O}$  anomaly of  $\sim 2.25\%$  to explain. Given the modern  $\sim -0.6\%$  per 100 m  $\delta^{18}\text{O}$ -elevation gradient in Greenland (Vinther et al., 2009), this suggests that the central Greenland ice sheet elevation may have been reduced by at most 325–450 m at the end of the last interglacial. Such a reduced elevation in central Greenland is expected to result from stronger melt in the coastal ablation zone and dynamical ice sheet response during the last interglacial compared to today. So far, no information can be extracted from the deepest parts of the NorthGRIP ice core regarding elevation changes. Air content measurements from the deepest parts of the GRIP ice core (Raynaud et al., 1997) suggest little change in Summit elevation. It is expected that the undisturbed parts of the NEEM ice core could bring further constraints.

## Sensitivity of interglacial Greenland temperature

V. Masson-Delmotte et al.

Title Page

Abstract

Introduction

Conclusions

References

Tables

Figures



Back

Close

Full Screen / Esc

Printer-friendly Version

Interactive Discussion





melt on ocean circulation (as forcings for coupled ocean-atmosphere models), and also by the uncertainties on the isotope-temperature slopes. Here, we first make use of coupled ocean-atmosphere climate simulations run by one model, IPSL-CM4, under different orbital and CO<sub>2</sub> forcings.

At 126 ka, this model has a strong summer temperature response compared to earlier published runs e.g. (Otto-Bliesner et al., 2006), (Gröger et al., 2007), and propagates the orbitally forced Arctic warming into winter. There is evidence for a strong sea ice retreat in some Arctic areas during the last interglacial (Polyak et al., 2010). New sea ice proxy records would be extremely useful to assess the realism of model sea ice outputs. In these simulations, the IPSL-CM4 model does not include the feedbacks associated with vegetation changes. Increased boreal forest cover (CAPE, 2006) could be expected to induce continental spring warming due to albedo effect and summer cooling due to increased evapotranspiration (Otto et al., 2011).

The IPSL model depicts a very strong linear relationship between simulated summer Greenland temperature and summer insolation forcing from 6 orbital configurations (0, 6, 9.5, 115, 122 and 126 ka). The slope of this relationship appears smaller than the one which can be estimated from the NorthGRIP data for the late interglacial trends. This may be due to the lack of feedbacks such as ice sheet elevation changes. Sensitivity tests with parameterisations of Greenland melt however highlight the fact that a large Greenland meltwater flux (about 10 mm year<sup>-1</sup>) (Swingedouw et al., 2009) acts as a local negative feedback through the impact of a reduced AMOC, reducing the magnitude of 126 ka Greenland warming by about 0.5 °C. These tests, however, do not account for any changes in Greenland ice sheet topography.

The quantitative interpretation of the ice core data relies on estimates of the temporal isotope-temperature relationship. Because the simulated 126 ka annual mean temperature change is modest (<1 °C), and lower than expected from the ice core data, we also explore simulations conducted using boundary conditions from 2 × CO<sub>2</sub> and 4 × CO<sub>2</sub> as well as 4 °C warmer SST climates. We stress the fact that there is no physical analogy between the greenhouse and orbital forcings: the IPSL model response

## Sensitivity of interglacial Greenland temperature

V. Masson-Delmotte et al.

Title Page

Abstract

Introduction

Conclusions

References

Tables

Figures



Back

Close

Full Screen / Esc

Printer-friendly Version

Interactive Discussion





strongly differs in terms of seasonal and latitudinal temperature or water cycle changes.

During the last interglacial, the mid to high latitude summer warming occurs without a clear tropical or global anomaly and persists in winter at high latitudes; obliquity changes indeed induce reduced annual mean tropical insolation and ocean temperatures. This strongly differs from the impact of increased greenhouse gas concentrations, marked by year round tropical warming and strong winter warming at high latitudes. However, the magnitude of summer Arctic warming is very similar in the IPSL 126 ka and  $2 \times \text{CO}_2$  simulations. Moreover, our simple analysis of feedbacks affecting the top of atmosphere radiative budget has also demonstrated comparable magnitudes of changes in the albedo, cloud and atmospheric greenhouse feedbacks in summer. Given the importance of summer temperature on ice sheet ablation, these comparable magnitudes have relevance regarding the assessment of climate model feedbacks, changes in Greenland ice sheet mass balance, and implications for sea level.

The LMDZiso model outputs show strong shifts in the precipitation seasonality due to increased summer precipitation in response to the 6 ka and 126 ka orbital forcings (proportionally stronger than for increased  $\text{CO}_2$  simulations). If true, this suggests that the Greenland ice core interglacial data must be cautiously interpreted in terms of precipitation weighted signals with a summer bias. In the warm climate simulations, LMDZiso produces an isotope-temperature slope of  $\sim 0.3\text{‰}$  (within a 30% uncertainty). Shifts in moisture origin under warm summer conditions clearly reduce the imprint of Greenland temperature changes in the simulated  $\delta^{18}\text{O}$ . Such changes may be caused by changes in storm tracks or in the Hadley cell (Fischer and Jungclaus, 2010), in response to changing latitudinal temperature gradients, sea ice and land sea contrasts. The differences between isotopic model  $\delta^{18}\text{O}$  shifts may be due to different changes in moisture origin (especially the proportion of Arctic versus low latitude moisture). This aspect would deserve to be further investigated, perhaps using water tagging methods, and/or second order stable isotope information (e.g. deuterium excess, oxygen 17-excess) which could allow to test the realism of changes in moisture source

## Sensitivity of interglacial Greenland temperature

V. Masson-Delmotte et al.

Title Page

Abstract

Introduction

Conclusions

References

Tables

Figures

⏪

⏩

◀

▶

Back

Close

Full Screen / Esc

Printer-friendly Version

Interactive Discussion



characteristics (Kurita, 2011).

For LMDZiso, the simulated 6 ka and 126 ka  $\delta^{18}\text{O}$  is much weaker than the ice core signals. Given the range of isotope-temperature responses obtained under strongly warmer climates (+4°C SST, 4xCO<sub>2</sub>), the last interglacial ice core signal (~3‰) is only compatible with very large (precipitation weighted) temperature shifts (8 to 10°C) (at fixed elevation). The 126 ka LMDZiso simulation can also be reconciled with the ice core data, assuming a 300–400 m reduced elevation in central Greenland (and even larger surface elevation changes when considering the impact of meltwater on climate). In the future, this should be compared with information obtained from air content data (Raynaud et al., 1997) from the recent NEEM deep ice core. The robustness of this finding should be assessed by comparing last interglacial precipitation isotopic composition simulations conducted with different climate models.

In the coming years, the PMIP3 project is expected to allow climate model inter-comparison with standardized boundary conditions for the last interglacial. We also aim to perform simulations at 126 ka with a prescribed reduced Greenland ice sheet, in order to better assess the impact of elevation changes on temperature and precipitation isotopic composition. Intercomparisons of isotopic simulations both under last interglacial and increased CO<sub>2</sub> boundary conditions are needed, in order to better understand the robustness of the results. Finally, the consistency between changes in elevation changes, accumulation changes and meltwater flux really needs to be assessed, and the proper framework for this lies in interactive ice sheet- climate coupling, including water stable isotope tracers.

## Appendix A

### Method for radiative feedbacks analysis

Following (Braconnot et al., 2007), a simple feedback analysis was performed in order to quantify the main drivers of changes in the top of the atmosphere radiative budget

## Sensitivity of interglacial Greenland temperature

V. Masson-Delmotte et al.

Title Page

Abstract

Introduction

Conclusions

References

Tables

Figures

⏪

⏩

◀

▶

Back

Close

Full Screen / Esc

Printer-friendly Version

Interactive Discussion



(TOA) over and around Greenland (60–80° N, 60–10° W):

$$\Delta\text{TOA}_{\text{simul}} = \Delta\text{SWn}_{\text{simul}} + \Delta\text{LWn}_{\text{simul}} \quad (\text{A1})$$

where  $\Delta_{\text{simul}}$  is the change between a forced simulation (6, 9.5, 115, 122, 126 ka and  $2 \times \text{CO}_2$ ) and the control simulation (ctrl); SWn is the net shortwave radiation at the top of the atmosphere (positive downwards) and LWn the net longwave radiation (positive downwards).

$\Delta\text{SWn}_{\text{simul}}$  is driven by interplay between the insolation forcing and the albedo feedbacks. The actual insolation forcing  $\Delta\text{SWf}_{\text{simul}}$  corresponds to the net change in shortwave radiative forcing under the assumption of a constant planetary albedo (Hewitt and Mitchell, 1996). The shortwave radiative forcing (SWf) (at fixed planetary albedo) is estimated using the control simulation planetary albedo ( $\alpha_{\text{ctrl}}^{\text{tot}}$ ) and the prescribed change in insolation  $\Delta\text{SWi}_{\text{simul}}$  as :

$$\Delta\text{SWf}_{\text{simul}} = (1 - \alpha_{\text{ctrl}}^{\text{tot}}) \Delta\text{SWi}_{\text{simul}} \quad (\text{A2})$$

The albedo feedback then results from the changes in surface albedo, atmospheric diffusion and clouds:

$$\Delta\text{ALB}_{\text{simul}} = \Delta\text{SWn}_{\text{simul}} - \Delta\text{SWf}_{\text{simul}} \quad (\text{A3})$$

At first approximation, for clear sky conditions (cs), the change in shortwave radiation at the top of the atmosphere is primary due to changes in surface albedo (even though one cannot distinguish the effects of changes in atmospheric properties from changes in surface albedo). The snow and sea ice albedo effect can be thus be approximated from the difference in simulated clear sky (cs) net shortwave radiative fluxes, as:

$$\Delta\text{ALBcs}_{\text{simul}} = \Delta\text{SWn}_{\text{cs}_{\text{simul}}} - \Delta\text{SWf}_{\text{simul}} \quad (\text{A4})$$

The role of clouds on  $\Delta\text{SWn}_{\text{simul}}$  can then be estimated as the difference between the total and clear sky albedo feedbacks, or equivalently, by the change in cloud shortwave

## Sensitivity of interglacial Greenland temperature

V. Masson-Delmotte et al.

Title Page

Abstract

Introduction

Conclusions

References

Tables

Figures

⏪

⏩

◀

▶

Back

Close

Full Screen / Esc

Printer-friendly Version

Interactive Discussion



radiative forcing (with small uncertainties resulting from the differences in the area covered by clouds in the different simulations).

It is not easy to estimate the contribution of surface temperature, water vapour content, trace gases and lapse rate on the long wave emission at the top of the atmosphere ( $\Delta LWn_{\text{simul}}$ ). In the case of orbital forcing, all the terms that affect the longwave radiation are considered as feedbacks, which contrasts with the  $2 \times \text{CO}_2$  forcing which exerts a direct longwave forcing. Here, we only consider a bulk estimate of the total greenhouse effect ( $g$ ), considering the difference between the long wave emission at the surface and at the top of the atmosphere

$$\Delta g_{\text{simul}} = \Delta LWn_{\text{simul}} - \Delta PI_{\text{simul}} \quad (\text{A5})$$

with  $\Delta PI_{\text{simul}}$  the change in direct (Planck) emission at the surface temperature  $T_{S_{\text{simul}}}$  with respect to the control simulation, which can be approximated by:

$$\Delta PI_{\text{simul}} = 4\sigma T_{\text{ctrl}}^3 (T_{S_{\text{simul}}} - T_{S_{\text{ctrl}}}) \quad (\text{A6})$$

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## Sensitivity of interglacial Greenland temperature

V. Masson-Delmotte et al.

Title Page

Abstract

Introduction

Conclusions

References

Tables

Figures

⏪

⏩

◀

▶

Back

Close

Full Screen / Esc

Printer-friendly Version

Interactive Discussion





## Sensitivity of interglacial Greenland temperature

V. Masson-Delmotte et al.

Title Page

Abstract

Introduction

Conclusions

References

Tables

Figures

⏪

⏩

◀

▶

Back

Close

Full Screen / Esc

Printer-friendly Version

Interactive Discussion



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**Sensitivity of  
interglacial  
Greenland  
temperature**

V. Masson-Delmotte et al.

Title Page

Abstract

Introduction

Conclusions

References

Tables

Figures

◀

▶

◀

▶

Back

Close

Full Screen / Esc

Printer-friendly Version

Interactive Discussion



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## Sensitivity of interglacial Greenland temperature

V. Masson-Delmotte et al.

Title Page

Abstract

Introduction

Conclusions

References

Tables

Figures

⏪

⏩

◀

▶

Back

Close

Full Screen / Esc

Printer-friendly Version

Interactive Discussion



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**Sensitivity of  
interglacial  
Greenland  
temperature**

V. Masson-Delmotte et al.

[Title Page](#)[Abstract](#)[Introduction](#)[Conclusions](#)[References](#)[Tables](#)[Figures](#)[⏪](#)[⏩](#)[◀](#)[▶](#)[Back](#)[Close](#)[Full Screen / Esc](#)[Printer-friendly Version](#)[Interactive Discussion](#)

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## Sensitivity of interglacial Greenland temperature

V. Masson-Delmotte et al.

[Title Page](#)
[Abstract](#)
[Introduction](#)
[Conclusions](#)
[References](#)
[Tables](#)
[Figures](#)




[Back](#)
[Close](#)
[Full Screen / Esc](#)
[Printer-friendly Version](#)
[Interactive Discussion](#)


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## Sensitivity of interglacial Greenland temperature

V. Masson-Delmotte et al.

**Table 1.** Description of the simulations. The LMDZiso simulations were run for 5 yr with climatological forcing averaged from the IPSL-CM4 outputs, and results analysed for the last 3 yr of this simulation.

Name	Orbital forcing	Atmospheric composition	Greenland melt	Ocean surface
IPSL-0ka	0 ka	pre-industrial	No	calculated
IPSL-6 ka	6 ka	6 ka	No	calculated
IPSL-9.5 ka	9.5 ka	pre-industrial	No	calculated
IPSL-115 ka	115 ka	pre-industrial	No	calculated
IPSL-122 ka	122 ka	pre-industrial	No	calculated
IPSL-126 ka	126 ka	pre-industrial	No	calculated
IPSL-126 ka GM	126 ka	pre-industrial	Yes	calculated
IPSL-2 × CO <sub>2</sub>	0 ka	CMIP3	No	calculated
IPSL-4 × CO <sub>2</sub>	0 ka	CMIP3	No	calculated
LMDZiso-ctrl	0 ka	348 ppmv	No	Prescribed from AMIP
LMDZiso-6 ky	6 ka	280 ppmv	No	Prescribed as AMIP+ (IPSL 6 ky–IPSL 0 ky)
LMDZiso-126 ky	126 ka	280 ppmv	No	Prescribed as AMIP+ (IPSL 126 ky–IPSL 0 ky)
LMDZiso-126 ky GM	126 ka	280 ppmv	Prescribed from IPSL-126 ky GM	Prescribed as AMIP+ (IPSL-126 ky GM–IPSL 0 ky)
LMDZisoSST	0 ka	280 ppmv	No	AMIP+4 °C
LMDZiso 4 × C	0 ka	2 × 348 ppmv	No	IPSL 2 × CO <sub>2</sub>
LMDZiso 4 × C	0 ka	3 × 348 ppmv	No	IPSL 4 × CO <sub>2</sub>

[Title Page](#)
[Abstract](#)
[Introduction](#)
[Conclusions](#)
[References](#)
[Tables](#)
[Figures](#)

[Back](#)
[Close](#)
[Full Screen / Esc](#)
[Printer-friendly Version](#)
[Interactive Discussion](#)

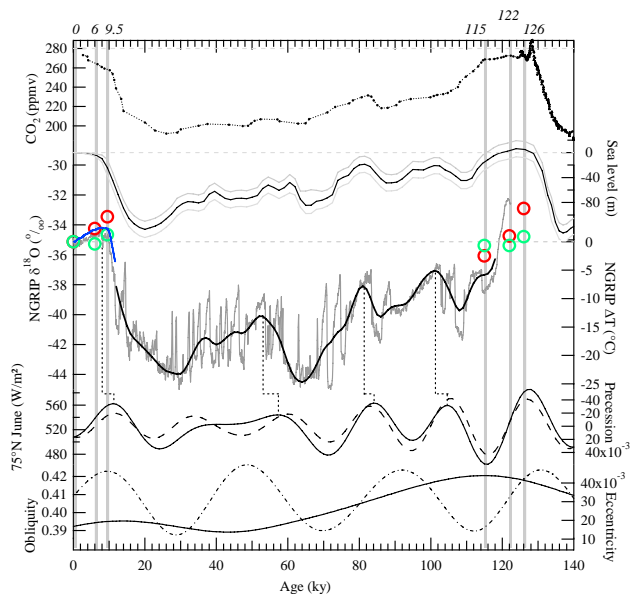

## Sensitivity of interglacial Greenland temperature

V. Masson-Delmotte et al.

**Table 2.** IPSL-CM4 results for Greenland (from grid points located above 1300 m elevation): annual mean, July and precipitation weighted temperature ( $^{\circ}\text{C}$ ) as well as deposition effect (difference between precipitation weighted and annual mean temperature) and ratio of summer (April–September) to annual precipitation. Results are given for the different simulations in response to orbital forcing only. Absolute values are given as well as anomalies with respect to the control simulation (numbers shown between parentheses).

Simulation	Annual mean Greenland temperature (anomaly) ( $^{\circ}\text{C}$ )	July Greenland temperature (anomaly) ( $^{\circ}\text{C}$ )	Ratio of summer half year (April–September) to annual precipitation (percentage of change)	Precipitation weighted Greenland temperature (anomaly) ( $^{\circ}\text{C}$ )	Deposition effect (anomaly) ( $^{\circ}\text{C}$ )
Control simulation	-28.3	-12.9	0.60	-25.9	2.4
6 ka	-27.9 (+0.4)	-10.6 (+2.3)	0.62 (+3%)	-24.7 (+1.2)	3.2 (+0.7)
9.5 ka	-27.5 (+0.8)	-8.5 (+4.4)	0.65 (+8%)	-23.3 (+2.6)	4.2 (+1.8)
115 ka	-28.8 (-0.5)	-15.4 (-2.5)	0.58 (-3%)	-27.0 (-1.1)	1.8 (-0.7)
122 ka	-28.1 (+0.2)	-11.9 (+1.0)	0.62 (+3%)	-25.1 (+0.8)	3.0 (+0.6)
126 ka	-27.4 (+0.9)	-7.1 (+5.8)	0.68 (+13%)	-22.3 (+3.6)	5.2 (+2.7)

[Title Page](#)
[Abstract](#)
[Introduction](#)
[Conclusions](#)
[References](#)
[Tables](#)
[Figures](#)
[Back](#)
[Close](#)
[Full Screen / Esc](#)
[Printer-friendly Version](#)
[Interactive Discussion](#)



**Fig. 1.** From top to bottom: black dots, atmospheric CO<sub>2</sub> concentration (Vostok and EDC ice cores) on the EDC3 age scale (Barnola et al., 1987; Laurantou et al., 2010); solid black line with uncertainties, estimation of eustatic sea level (Waelbroeck et al., 2002); NorthGRIP ice core  $\delta^{18}\text{O}$  data on a 20 yr resolution on GICC05 and EDC3 age scale (Capron et al., 2010a). The orbital component of the record is displayed (thick black line) and was calculated using the first three components of a singular spectrum analysis. A tentative estimate of the temperature change is also displayed, following (Masson-Delmotte et al., 2005b) (right axis). The reconstruction of the Holocene Greenland temperature (at fixed elevation) (Vinther et al., 2009) is displayed as a bold blue line. Summer (red) and annual mean (green) temperature anomalies simulated by the IPSL model are displayed as open circles for 6, 9.5, 115, 122, and 126 ka. The 75° N June insolation (black line,  $\text{W m}^{-2}$ ) and orbital parameters (precession parameter – long dashed line, obliquity – short dashed line, and eccentricity – solid black line) are displayed in the two lowest panels.

## Sensitivity of interglacial Greenland temperature

V. Masson-Delmotte et al.

Title Page

Abstract

Introduction

Conclusions

References

Tables

Figures



Back

Close

Full Screen / Esc

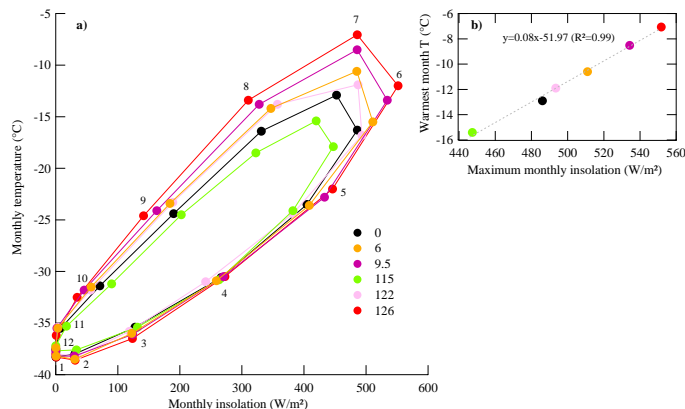
Printer-friendly Version

Interactive Discussion



## Sensitivity of interglacial Greenland temperature

V. Masson-Delmotte et al.

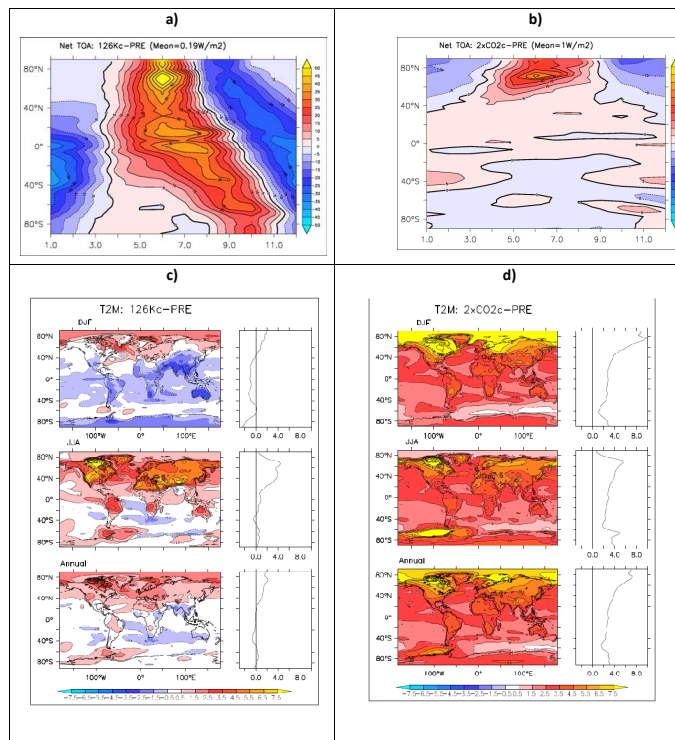


**Fig. 2.** (a) Seasonal cycle of IPSL model simulated central Greenland (>1300 m) temperature (°C) as a function of the seasonal cycle of TOA incoming solar radiation at 75° N ( $\text{W m}^{-2}$ ) for different orbital configurations (0, 6, 9.5, 115, 122 and 126 ka). For each period, the monthly data are displayed; black numbers indicate the number of the month (from 1 for January to 12 for December). The elliptic shape results from the phase lag between temperature and insolation. (b) Regression between maximum monthly insolation and the IPSL model central Greenland maximum monthly summer temperature (occurring one month after maximum insolation). A linear relationship is observed, with a slope of  $0.08 \text{ °C per } \text{W m}^{-2}$ . The same color code is used as in panel (a) for the various simulations.

[Title Page](#)
[Abstract](#)
[Introduction](#)
[Conclusions](#)
[References](#)
[Tables](#)
[Figures](#)
[⏪](#)
[⏩](#)
[◀](#)
[▶](#)
[Back](#)
[Close](#)
[Full Screen / Esc](#)
[Printer-friendly Version](#)
[Interactive Discussion](#)

## Sensitivity of interglacial Greenland temperature

V. Masson-Delmotte et al.



**Fig. 3.** Comparison of anomalies between the last interglacial and pre-industrial control IPSL model simulations (left) and  $2 \times CO_2$  and pre-industrial control simulation (right) for: **(a)** and **(b)** top of atmosphere net radiative budget ( $W m^{-2}$ ); **(c)** and **(d)** surface air temperature ( $^{\circ}C$ ) and **(e)** and **(f)** evaporation (mm/day). For panels **(a)** and **(b)**, anomalies as displayed as a function of month number (horizontal axis) and latitude (vertical axis). For panels **(c)** to **(f)**, anomalies are displayed as a function of longitude and latitude, for DJF (December-January-February), JJA (June-July-August) and for the annual mean. On the right side of each panel **(c)** to **(f)**, zonal mean anomalies are also displayed as a function of latitude.

Title Page

Abstract

Introduction

Conclusions

References

Tables

Figures

◀

▶

◀

▶

Back

Close

Full Screen / Esc

Printer-friendly Version

Interactive Discussion

## Sensitivity of interglacial Greenland temperature

V. Masson-Delmotte et al.

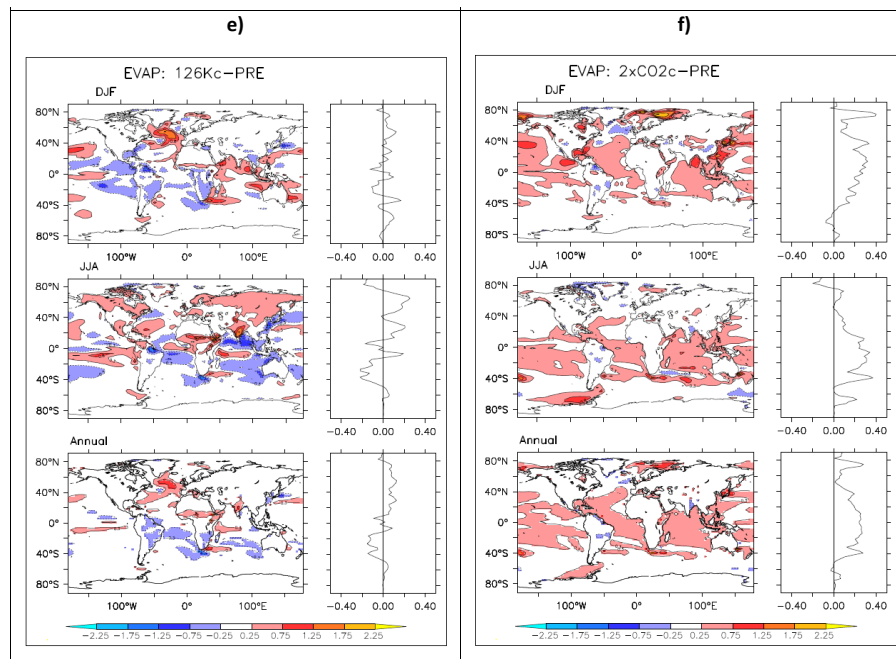


Fig. 3. Continued.

Title Page

Abstract

Introduction

Conclusions

References

Tables

Figures

⏪

⏩

◀

▶

Back

Close

Full Screen / Esc

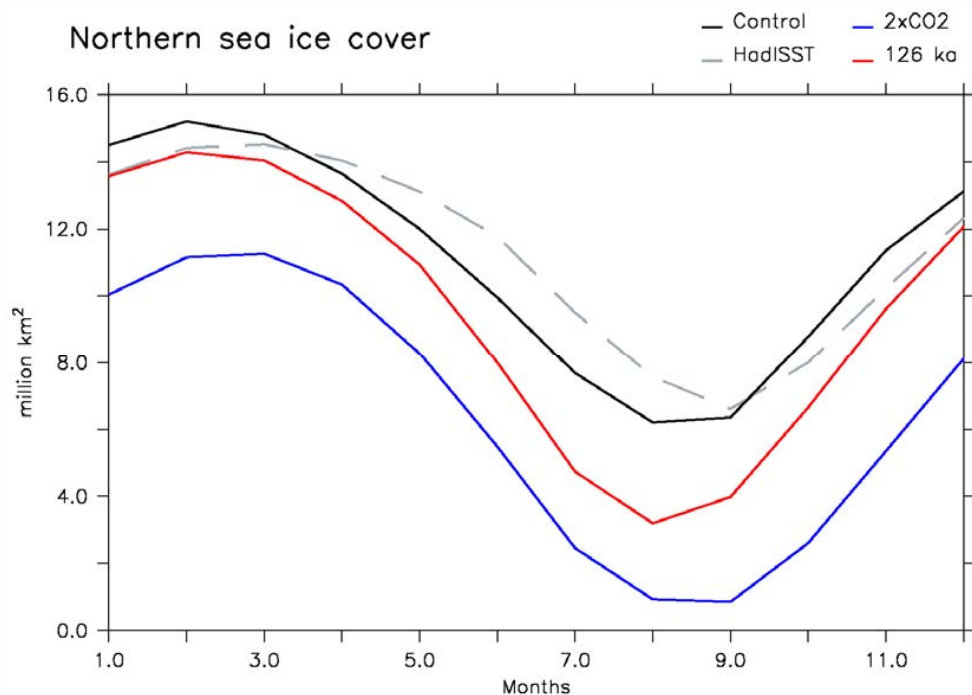
Printer-friendly Version

Interactive Discussion



**Sensitivity of  
interglacial  
Greenland  
temperature**

V. Masson-Delmotte et al.



**Fig. 4.** Monthly seasonal cycle of Northern Hemisphere sea ice extent for present day (black), 126 ka (red) and  $2 \times \text{CO}_2$  (blue) simulated by IPSL-CM4. Present day (1900–2010) climatological data are also displayed (dashed grey).

Title Page

Abstract

Introduction

Conclusions

References

Tables

Figures

◀

▶

◀

▶

Back

Close

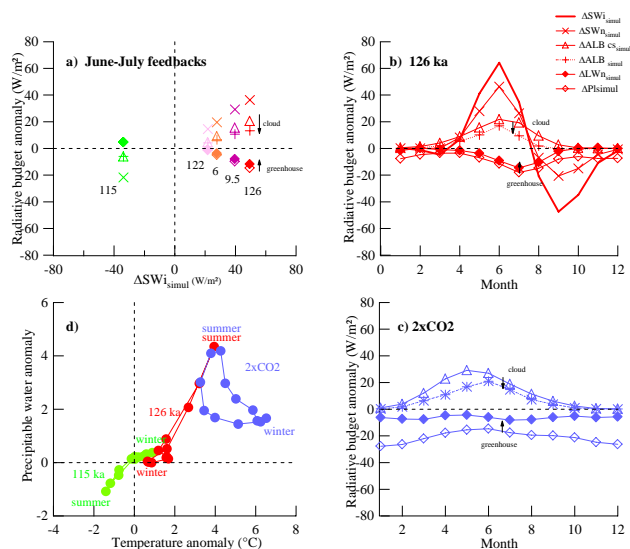
Full Screen / Esc

Printer-friendly Version

Interactive Discussion

Sensitivity of  
interglacial  
Greenland  
temperature

V. Masson-Delmotte et al.



**Fig. 5.** Analysis of atmospheric feedbacks affecting the top of atmosphere radiative budget. **(a)** June–July changes in radiative budget terms (see text and figure legend for details) as a function of June–July incoming solar radiation ( $W m^{-2}$ ), for different orbital contexts (115 ka, green; 122 ka, pink; 6 ka, orange; 9.5 ka, violet and 126 ka, red). Arrows depict the magnitude of albedo (difference between “+” and triangle symbols), cloud (difference between “+” and “x” symbols) and greenhouse (difference between open and filled diamonds) feedbacks for 126 ka. **(b)** Monthly values of the radiative budget 126 ka anomalies with respect to the control simulation (see text and legend for details) ( $W m^{-2}$ ). **(c)** Same as **(a)** but for  $2 \times CO_2$ . **(d)** Seasonal cycle of precipitable water anomaly as a function of temperature anomaly ( $^{\circ}C$ ) with respect to the control simulation, for 115 ka (green), 126 ka (red) and  $2 \times CO_2$  (blue).

Title Page

Abstract

Introduction

Conclusions

References

Tables

Figures

◀

▶

◀

▶

Back

Close

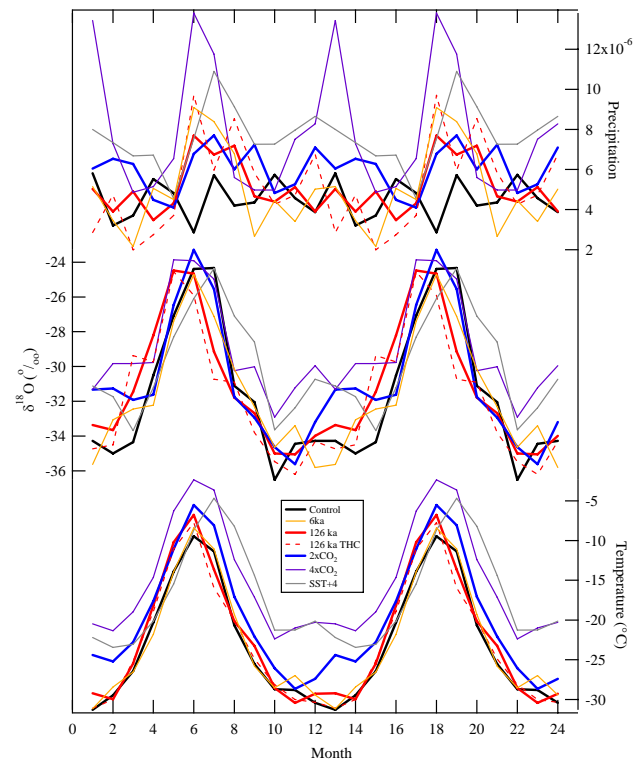
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Interactive Discussion

## Sensitivity of interglacial Greenland temperature

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**Fig. 6.** Monthly seasonal cycle of temperature, precipitation and precipitation d18O simulated by LMDZiso for different sets of boundary conditions (AMIP control, 6 ka, 126 ka, +4 °C SST, 2× and 4 × CO<sub>2</sub> concentrations) prescribed using the IPSL-CM4 sea surface conditions (see Table 1). For readability, the seasonal cycle has been repeated over 2 yr (24 months).

Title Page

Abstract

Introduction

Conclusions

References

Tables

Figures

◀

▶

◀

▶

Back

Close

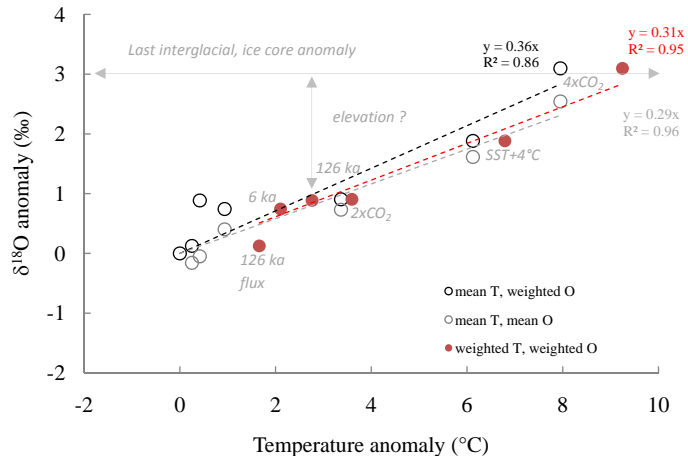
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Interactive Discussion

## Sensitivity of interglacial Greenland temperature

V. Masson-Delmotte et al.



**Fig. 7.** Simulated anomalies of Greenland precipitation weighted  $\delta^{18}\text{O}$  as a function of mean temperature (black open circles) and precipitation weighted temperature (red filled circles). Simulated anomalies of Greenland annual mean  $\delta^{18}\text{O}$  as a function of annual mean temperature (grey open circles) are also displayed for all the LMDZiso simulations. Anomalies are calculated with respect to the AMIP control simulation. Linear regressions are also displayed.

Title Page

Abstract

Introduction

Conclusions

References

Tables

Figures

◀

▶

◀

▶

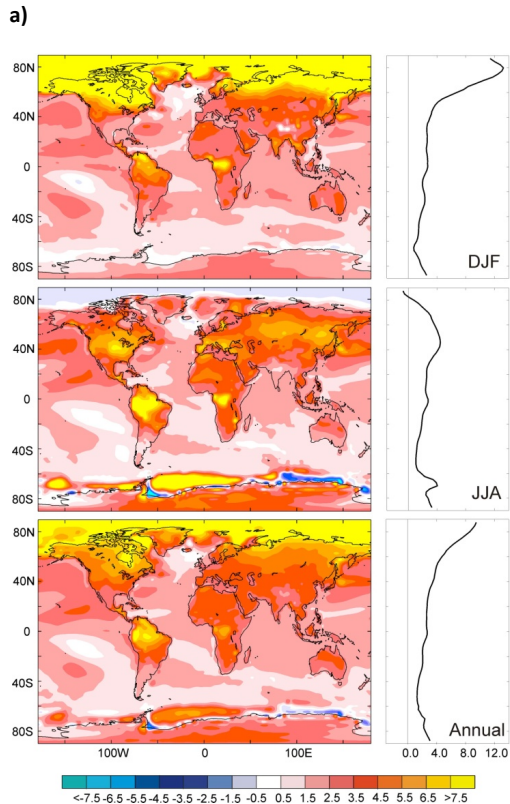
Back

Close

Full Screen / Esc

Printer-friendly Version

Interactive Discussion



**Fig. A1.** Anomalies between the 2100 and reference HadAM3iso model simulations for: **(a)** surface air temperature ( $^{\circ}\text{C}$ ) and **(b)** evaporation (mm day) and **(c)** sea ice extent changes. Anomalies as displayed as a function of longitude and latitude, for DJF (December-January-February), JJA (June-July-August) and for the annual mean. On the right side of each panel **(a)** to **(d)**, zonal mean anomalies are also displayed as a function of latitude.

## Sensitivity of interglacial Greenland temperature

V. Masson-Delmotte et al.

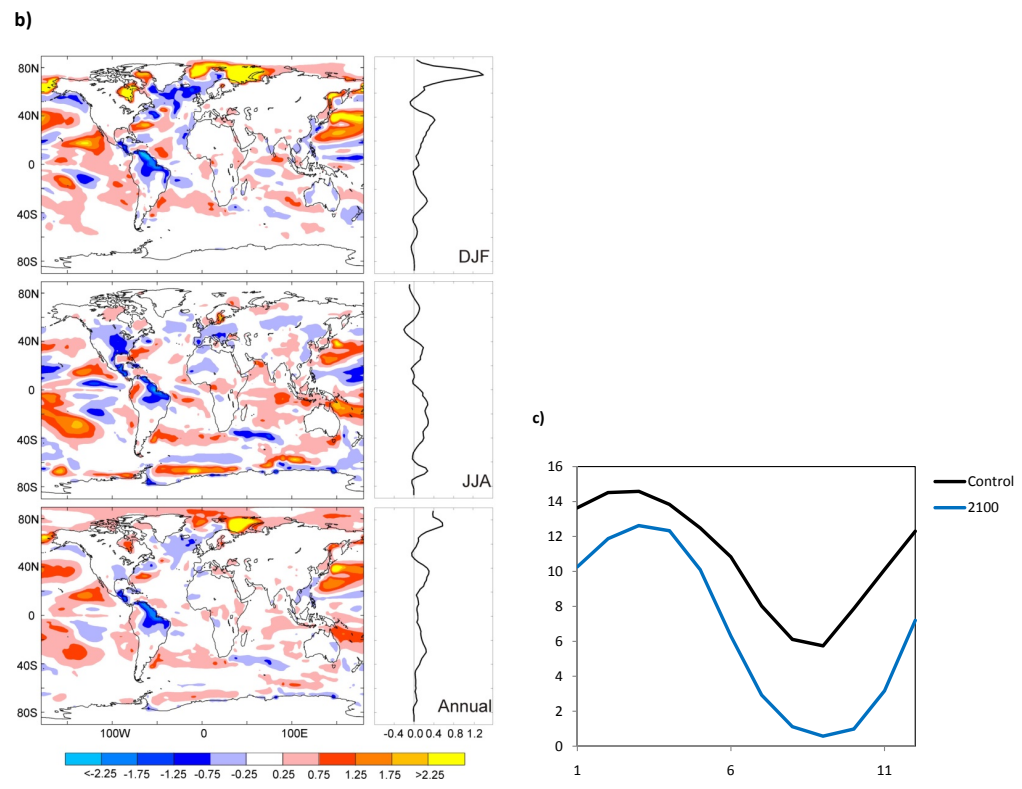


Fig. A1. Continued.

[Title Page](#)

[Abstract](#) | [Introduction](#)

[Conclusions](#) | [References](#)

[Tables](#) | [Figures](#)

[⏪](#) | [⏩](#)

[⏴](#) | [⏵](#)

[Back](#) | [Close](#)

[Full Screen / Esc](#)

[Printer-friendly Version](#)

[Interactive Discussion](#)

