

1 **Impact of ice sheet meltwater fluxes on the climate**
2 **evolution at the onset of the Last Interglacial**

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14 **Abstract**

15 Large climate perturbations occurred during the transition between the penultimate glacial
16 period and the Last Interglacial (Termination II), when the ice sheets retreated from their
17 glacial configuration. Here we investigate the impact of ice sheet changes and associated
18 freshwater fluxes on the climate evolution at the onset of the Last Interglacial. The period
19 from 135 to 120 kyr BP is simulated with the Earth system model of intermediate complexity
20 LOVECLIM v.1.3 with prescribed evolution of the Antarctic ice sheet, the Greenland ice
21 sheet and the other Northern Hemisphere ice sheets. Variations in meltwater fluxes from the
22 Northern Hemisphere ice sheets lead to North Atlantic temperature changes and modifications
23 of the strength of the Atlantic meridional overturning circulation. By means of the
24 interhemispheric see-saw effect, variations in the Atlantic meridional overturning circulation
25 also give rise to temperature changes in the Southern Hemisphere, which are additionally
26 modulated by the direct impact of Antarctic meltwater fluxes into the Southern Ocean.
27 Freshwater fluxes from the melting Antarctic ice sheet lead to a millennial time scale oceanic

28 cold event in the Southern Ocean with expanded sea ice as evidenced in some ocean sediment
29 cores, which may be used to constrain the timing of ice sheet retreat.

30

31 **1 Introduction**

32 Understanding the climate and ice sheet evolution during past warm periods in the history of
33 the Earth may provide important insights for projections of future climate and sea-level
34 changes. The growing amount of paleo-reconstructions for the Last Interglacial period (e.g.
35 Govin et al., 2012; Capron et al., 2014) in combination with improved model simulations of
36 this most recent warm period (e.g. Bakker et al., 2013; Lunt et al., 2013, Langebroek and
37 Nisancioglu, 2014; Loutre et al., 2014) make it an interesting target for studying the coupled
38 climate-ice sheet system.

39 According to reconstructions, the Last Interglacial (LIG, from ~130-115 kyr BP) was
40 characterised by a global annual mean surface temperature of up to 2° C above the pre-
41 industrial (e.g. Turney and Jones, 2010; Capron et al., 2014) and a sea-level high stand of 6-9
42 m above the present day (Kopp et al., 2009; Dutton and Lambeck, 2012). As the penultimate
43 glacial maximum was at least as severe as the Last Glacial Maximum (LGM) in both
44 hemispheres (EPICA community members, 2004; Svendsen et al., 2004), this implies a large
45 amplitude glacial-interglacial transition in terms of temperature and ice sheet configuration.
46 At the onset of the LIG, a rapid warming of ~10°C from the preceding cold state is recorded
47 in deep Antarctic ice cores (Masson-Delmotte et al., 2011) to have occurred between ~135
48 kyr BP and 130 kyr BP. Current ice core records from the Greenland ice sheet (GrIS) do not
49 extend long enough back in time to cover the entire penultimate deglaciation and associated
50 warming (NEEM community members, 2013), but a similar timing and magnitude of
51 warming compared to the Antarctic can be reconstructed for sea surface temperatures off the
52 West European margin (Sánchez Goñi et al., 2012). The warming is closely related with an
53 ice sheet retreat in both hemispheres. Despite large uncertainties in reconstructions, the global
54 sea-level stand at 135 kyr BP of as low as -80 m (Grant et al., 2012) is indicative of the large
55 amount of freshwater that entered the ocean in the form of meltwater from the retreating ice
56 sheets during Termination II. Aside from determining the amplitude of sea-level changes,
57 which is the focus of many studies (e.g. Robinson et al., 2011; Stone et al., 2013), the
58 associated climate impacts and possible feedbacks on the ice sheet evolution of this

59 freshwater forcing are an important element for a process understanding of the coupled
60 climate-ice sheet changes at that time.

61 A climatic mechanism that is thought to be directly related to changes in the NH ice sheet
62 freshwater fluxes (FWF) is the interhemispheric see-saw effect (Stocker, 1998) that links SH
63 warming to a weakening of the Atlantic meridional overturning circulation (AMOC). If the
64 see-saw effect was active during the onset of the LIG, NH ice sheet melting during
65 Termination II would have been the cause for a substantial AMOC weakening and NH
66 cooling, while reduced interhemispheric heat transport would have caused a gradual SH
67 warming (Stocker and Johnson 2003). The see-saw mechanism was evoked to explain part of
68 the peak Antarctic warming during the LIG (e.g. Holden et al., 2010; Marino et al., 2015),
69 even though some Southern Ocean (SO) warming was shown by Langebroek and Nisancioglu
70 (2014) to be possible with orbital forcing alone (without NH freshwater forcing). The see-saw
71 mechanism has been speculated to have caused increased Antarctic ice shelf melting and
72 West Antarctic ice sheet (WAIS) retreat (Duplessy et al., 2007). The retreat of the WAIS,
73 which is believed to have been grounded at the edge of the continental shelf during the
74 penultimate glaciation, generated a large anomalous flux of freshwater into the SO. Such
75 freshwater forcing could have had a substantial influence on the SO configuration in terms of
76 sea ice extent and ocean circulation as shown in model experiments for the last deglaciation
77 (Menviel et al., 2011), for future global warming scenarios (Swingedouw et al., 2008) and for
78 the present day (Bintanja et al., 2013). The impact of increased Antarctic FWF is thought to
79 consist of a surface ocean freshening, stratification of the surface ocean and cooling, in turn
80 promoting sea ice growth (e.g. Bintanja et al., 2013) and reduced Antarctic Bottom Water
81 (AABW) formation (Menviel et al., 2011). Recently, Golledge et al. (2014) suggested that
82 such a mechanism might also have provided a feedback on Antarctic ice sheet (AIS) retreat
83 for meltwater pulse 1A during the last glacial-interglacial transition (Termination I), by
84 promoting warming of mid-depth ocean waters that provide additional heat for melting ice
85 shelves.

86 In the present work, we study the effect of evolving ice sheet boundary conditions on the
87 climate, by simulating the climate evolution at the onset and over the course of the LIG with
88 an Earth system model of intermediate complexity (EMIC). The model is forced with realistic
89 ice sheet boundary conditions from offline simulations of ice dynamic models of the AIS and
90 GrIS and reconstructions of the other NH ice sheets. With this study we extend the work of

91 Loutre et al. (2014) by additionally including dynamic ice sheet changes of the GrIS and AIS
92 and focusing on the effect of ice sheet freshwater fluxes on the climate, particularly in the
93 Southern Hemisphere (SH). The model and experimental setup are described in section 2 and
94 3, respectively, followed by results (Sect. 4, 5 and 6), their discussion in section 7 and
95 conclusions (Sect. 8).

96

97 **2 Model description**

98 We use the EMIC LOVECLIM version 1.3, which includes components representing the
99 atmosphere, the ocean and sea ice, the terrestrial biosphere and the ice sheets (cf. Figure 1).
100 The model has been utilised in a large number of coupled climate-ice sheet studies (e.g.
101 Driesschaert et al., 2007; Swingedouw et al., 2008; Goelzer et al., 2011; 2012a; Loutre et al.,
102 2014) and is described in detail in Goosse et al. (2010).

103 In this study, the climate components are forced by time-evolving ice sheet boundary
104 conditions, which are calculated off-line, i.e. uncoupled from the climate evolution. Our
105 modelling approach for the ice sheets consists of a combination of reconstructed NH ice
106 sheets (except the GrIS) based on geomorphological data (Sec. 2.1) and of standalone ice
107 dynamic simulations of the GrIS and AIS (Sec. 2.2). In either case, the boundary conditions
108 provide time evolving topography, ice sheet extent (albedo) and spatially and temporally
109 variable FWF to the climate model. As common practice for simulations with large amounts
110 of freshwater input, we conserve global salinity and global volume in the ocean model to
111 avoid numerical problems.

112 **2.1 Northern Hemisphere ice sheet forcing**

113 We have little geomorphological evidence for Northern Hemisphere (NH) ice sheet evolution
114 during Termination II since it was mostly destroyed by the re-advance leading to the LGM.
115 Therefore, the reconstruction of NH ice sheet evolution for the period of interest is made
116 based on information from the last deglaciation. The method was already described in some
117 detail in Loutre et al. (2014). Nevertheless, we include a more thorough description here
118 (Appendix A). The resulting boundary conditions used to force the climate model consist of a
119 chronology of ice mask and surface elevation changes (Figure 2) and freshwater fluxes
120 (Figure 3b) over the entire LIG period. Support for the derived chronology of NH ice sheet
121 evolution and their FWF can be found in records of ice-rafted detritus (IRD) from the

122 subpolar North Atlantic (Kandiano et al., 2004; Oppo et al., 2006). These show variability of
123 similar signature during the deglaciation and in particular a last IRD peak at ~128 kyr BP
124 preceding low IRD levels throughout the LIG.

125 **2.2 Simulations of the Greenland and Antarctic ice sheets**

126 For the present study, the climate components are partially forced by results from stand-alone
127 simulations of the GrIS and AIS, which have been adapted from existing ice sheet model
128 experiments (Huybrechts 2002). The configuration of both ice sheet models and the forcing
129 interface follows the description in Goosse et al. (2010) with the following exceptions.
130 Forcing for the ice sheet models is derived from scaling present-day observations of
131 temperature and precipitation with indices based on ice core records, as often done for long-
132 term paleo ice sheet modelling (e.g. Huybrechts, 1990; Letréguilly et al., 1991; Zweck and
133 Huybrechts, 2005; Greve et al., 2011). For the GrIS the forcing record was created following
134 Fürst et al. (2015). We combine a synthesised Greenland $\delta^{18}\text{O}$ record derived from Antarctica
135 Dome C using a bipolar seesaw model (Barker et al., 2011) with the NEEM temperature
136 reconstruction (NEEM community members, 2013) between 128.44 kyr BP and 120 kyr BP.
137 The Barker $\delta^{18}\text{O}$ record is converted to a spatially uniform temperature anomaly with a
138 constant temperature/ isotope factor $\Delta T = 2.4 \text{ }^\circ\text{C}/\text{‰} * (\delta^{18}\text{O} + 34.83)$ as in Huybrechts (2002).
139 Positive temperature anomalies of the NEEM record are scaled by a factor 0.6 to fulfil
140 constraints on maximal ice sheet retreat from Camp Century and Dye3 ice core locations that
141 are assumed to have been ice covered during the LIG. This places the GrIS evolution in the
142 range of former model estimates during that period (e.g. Robinson et al., 2011; Born and
143 Nisancioglu, 2012; Stone et al., 2013). Such scaling is in line with recent studies (e.g. van de
144 Berg et al., 2013; Merz et al., 2014; Sjolte et al., 2014; Steen-Larsen et al., 2014) that put in
145 question the high temperature of the central estimate reconstructed from the NEEM record.
146 Precipitation rates for ice sheet forcing vary percentagewise as a function of the $\delta^{18}\text{O}$ record.
147 The AIS forcing is derived directly from the Antarctica Dome C record (EPICA community
148 members, 2004), following again procedures described by Huybrechts (2002). Here
149 precipitation changes are assumed proportional to the saturated water vapour pressure
150 gradient relative to the temperature above the surface inversion layer. Furthermore, both ice
151 sheet models are forced by changes in global sea-level stand based on the benthic deep-sea
152 record of Lisiecki and Raymo (2005) for the GrIS and on a more recent sea-level
153 reconstruction using Red Sea data (Grant et al., 2012) for the AIS, where the sea-level

154 changes are the dominant forcing. The chronology of the Red Sea record is expected to be
155 more accurate since new dating techniques are applied (Grant et al., 2012). The impact of
156 using another sea-level record for the GrIS simulation over the LIG is small, because of the
157 largely land-based character of the ice sheet during that period. The AIS model is run at a
158 horizontal resolution of 20 x 20 km instead of 10 km x 10 km (as in the standard LOVECLIM
159 configuration and for the GrIS model) due to computational constraints for the relatively long
160 duration of the LIG simulation.

161 To embed the dynamic GrIS simulation in the other NH boundary conditions, the geometric
162 evolution of the GrIS overrides prescribed changes where Greenland ice is present. Therefore,
163 the prescribed ice sheet evolution and associated FWF are not limited by the present-day
164 configuration of the GrIS as in Loutre et al. (2014). The ice sheet evolution is illustrated in
165 Figure 2 for the modelled GrIS embedded in the NH reconstruction (top) and for the modelled
166 AIS (bottom). Ice volume evolution for the NH ice sheets and GrIS and AIS are given in
167 Figure 3a and Figure 3c, respectively. The FWF from the dynamic GrIS and AIS (Figure 3d)
168 replace the background freshwater flux from runoff over land calculated by the land model.

169 In our setup, the combined sea-level contributions from Antarctica and the NH (including
170 Greenland) fall within the 67% confidence interval of probabilistic sea-level reconstructions
171 (Kopp et al., 2009) for the first peak in sea-level contributions and the following period
172 (~124-120 kyr BP). For both hemispheres, the final 20 m rise in sea-level at the onset of the
173 LIG is however steeper and occurs 1~2 kyr earlier as compared to the reconstructions. When
174 assuming a maximum contribution from glaciers (0.42 ± 0.11) and an additional estimate for
175 thermal expansion of the ocean (0.4 ± 0.3) as given by Masson-Delmotte et al. (2013), the
176 assumed ice sheet evolution in our setup reproduces well the average sea-level contribution
177 between 125 and 120 kyr BP from the best estimate of Kopp et al. (2009), but it does not
178 represent the multi-peak structure of global sea-level contribution during the LIG as suggested
179 by Kopp et al. (2009, 2013). More details about the ice sheet and sea-level evolution can be
180 found in a companion paper (Goelzer et al., 2016) that specifically deals with the sea-level
181 contribution of the ice sheets during the LIG in a fully coupled model set-up.

182 **2.3 Initialisation**

183 The goal of our initialisation technique is to prepare a climate model state for the transient
184 simulations starting at 135 kyr BP that exhibits a minimal coupling drift. Both the GrIS and

185 AIS models are integrated over the preceding glacial cycles and the entire LIG in stand-alone
186 mode. The climate model is then initialized to a steady state with ice sheet boundary
187 conditions, greenhouse gas (GHG) forcing and orbital parameters for the time of coupling
188 (135 kyr BP). In this way, when LOVECLIM is integrated forward in time for transient
189 experiments, the climate component is already relaxed to the ice sheet boundary conditions
190 and exhibits a minimal model drift in unforced control experiments (not shown).

191

192 **3 Experimental setup**

193 All simulations are forced by time-dependent changes in GHG concentrations and insolation
194 running from 135 kyr BP until 120 kyr BP (Figure 4). The radiative forcing associated with
195 the reconstructed GHG levels (Petit et al., 1999; Pépin et al., 2001; Raynaud et al., 2005;
196 Loulergue et al., 2008; Spahni et al., 2005) is below preindustrial values for most of this
197 period and barely exceeds it at ~128 kyr BP. The changes in the distribution of insolation
198 received by the Earth are dynamically computed from the changes in the orbital configuration
199 (Berger, 1978) and represent the governing NH forcing during peak LIG conditions aside
200 from evolving ice sheet boundary conditions. In the following, we will compare results of the
201 reference experiment with all ice sheet boundary conditions evolving in time (Reference) to
202 experiments in which the ice sheet boundary conditions are partially fixed to the pre-industrial
203 configuration (Table 1). To disentangle the effects of the individual ice sheets, the
204 experiments noGfwf (suppressed GrIS freshwater fluxes) and noAGfwf (suppressed FWF
205 from both AIS and GrIS) are complemented by two predecessor experiments with fixed AIS
206 and GrIS and evolving NH boundary conditions (noAG), as well as a climate experiment
207 forced by insolation and GHG changes only with all ice sheet boundary conditions fixed
208 (noIS). The latter two experiments correspond to the allLR and IGonly experiments from
209 Loutre et al. (2014).

210

211 **4 Effect of GrIS and AIS on the temperature evolution at the onset of the LIG**

212 Including the forcing from the NH ice sheets in terms of configuration and FWF has been
213 shown by Loutre et al. (2014) to be crucial to simulate the onset of the LIG temperature
214 increase and its amplitude variations more in line with proxy records. This helps to partially
215 overcome problems of EMICs (and general circulation models) to simulate the strong

216 temperature contrasts inferred from proxy reconstructions (Bakker et al., 2013; Lunt et al.,
217 2013). The increased amplitude of temperature changes in our simulations is due to albedo
218 and elevation changes in addition to the larger effect of the implied freshwater forcing from
219 the NH ice sheets (Loutre et al., 2014). Here the Loutre et al. (2014) experiments are
220 complemented with runs that additionally include changes in ice sheet configuration and FWF
221 from the GrIS and AIS. We first discuss the effect of including these additional ice sheet
222 boundary conditions. A specific focus on the FWF follows in section 5.

223 The temperature evolution (Figure 5) before 127 kyr BP is in both hemispheres strongly
224 influenced by the ice sheet boundary conditions and in particular by the freshwater forcing
225 from the ice sheets. The experiments including FWF from the NH ice sheets (Reference and
226 noAG) clearly show temperature variations on the multi-millennial time scale in both
227 hemispheres following variations in ice sheet freshwater input (cf. Figure 3). Differences in
228 the temperature evolution between noAG and the reference experiment are small in the NH,
229 where the additional freshwater flux from Greenland is small compared to the other sources.
230 In the SH, by contrast, a large perturbation arises around 130 kyr BP, when FWF from the
231 AIS peak. Global mean and hemispheric mean temperatures are similar in all runs after ~127
232 kyr BP, when the ice sheets have largely reached their interglacial configuration and FWF are
233 similar between the different experiments. An exception is the GrIS, which is retreating until
234 ~120 kyr BP but accounts for only a small FWF contribution. The similarity of the results in
235 the runs after ~127 kyr BP implies that the temporal memory of the response to ice sheet
236 changes in the system is limited to the multi-centennial time scale, at least for the surface
237 climate. The location of largest freshwater induced temperature variations in the NH is the
238 North Atlantic between 40° N and 80° N. Here changes in the AMOC cause a perturbation of
239 the northward oceanic heat transport and temperature changes, which are further amplified by
240 sea ice-albedo and insulation feedbacks. Greenland experiences maximum warming in the
241 reference experiment around 125 kyr BP of up to 2.7°C in the annual mean compared to the
242 pre-industrial over remaining ice covered central Greenland. Here the temperature evolution
243 is largely similar to the experiment with GrIS changes not accounted for (noAG), which
244 exhibits a maximum warming of 2.4°C (Loutre et al., 2014). However, the summer warming
245 reaches up to 10°C at the northern margin and even up to 14°C over southern margins over a
246 then ice-free tundra (not shown). The strong warming in the ice sheet periphery is due to a
247 combination of elevation changes and local albedo changes, confined to the immediate region
248 of ice sheet lowering and retreat. In the SH, the largest temperature perturbations linked to

249 both NH and SH freshwater fluxes occur in the SO. The largest warming over the ice sheet
250 itself is simulated over the WAIS (not shown) and is mainly a consequence of the local
251 elevation changes as the ice sheet retreats. However, mainly due to the marine based character
252 of the WAIS, albedo changes are much more limited compared to Greenland as the retreating
253 ice sheet surface is mostly replaced by sea ice. Modelled temperature changes over the East
254 Antarctic ice sheet (EAIS) have been compared to temperature reconstructions for four ice
255 core locations (Figure 6). The reference experiment shows a more pronounced warming
256 between 135 and 129.5 kyr BP compared to the experiments excluding Antarctic ice sheet
257 changes (noAG and noIS). While the modelled warming still appears to be underestimated
258 and delayed compared to the reconstructions, the reference simulation clearly improves the
259 representation of the EAIS temperature evolution compared to experiments with fixed
260 Antarctic boundary conditions.

261

262 **5 Role of ice sheet meltwater fluxes**

263 To study the role of the different freshwater contributions from the ice sheets in more detail
264 and evaluate their importance for the climate evolution, we compare additional simulations
265 where FWF from the GrIS and AIS are suppressed relative to the reference experiment
266 (Figure 7). The ice sheet configuration (topography and albedo) remains unchanged in these
267 experiments. The effect of AIS FWF can therefore be evaluated as the difference between
268 noGfwf and noAGfwf, whereas the effect of GrIS FWF becomes apparent from comparing
269 the reference simulation with noGfwf. The AIS FWF (Figure 7f) leads to considerable
270 changes in the SH, but has very little impact on the NH temperature evolution (cf. Figure 5b).
271 Conversely, variations in the NH (Figure 7a) and GrIS freshwater forcing on millennial time-
272 scales imply temperature changes in the SH on a background of general LIG warming.

273 Differences between the experiments in the AMOC evolution (Figure 7b) are largely
274 explained by whether FWF from the NH ice sheets and the GrIS are included or not. Here
275 AMOC strength is calculated as the maximum value of the meridional overturning stream
276 function below the Ekman layer in the Atlantic Ocean between 45° and 65° N. The effect of
277 the FWF from the GrIS (cf. Reference and noGfwf in Figure 7b) is limited compared to the
278 large impact of the general NH ice sheet forcing and consists of an additional weakening of
279 the AMOC. It is most pronounced during periods of AMOC recovery and after 130 kyr BP,
280 when melting of the GrIS beyond its present-day configuration sets in. Note that the simulated

281 evolution of AMOC strength in the reference experiment is in good agreement with paleo
282 evidence based on $\delta^{13}\text{C}$ data (Bauch et al., 2012) and in particular with a recent reconstruction
283 based on chemical water tracers (Böhm et al., 2015). The timing of Heinrich Stadial 11 (~132
284 kyr BP) and the variations in AMOC strength after that are well captured by our reference
285 simulation, which gives independent credibility to our NH ice sheet reconstructions.

286 The evolution of NH sea ice area (Figure 7c) generally shows maxima at times of AMOC
287 minima and vice versa and is closely linked to NH surface temperature variations (cf. Figure
288 5b) by modifying the heat exchange between ocean and atmosphere. The largest sea ice area
289 between 135 and 130 kyr BP is simulated in the reference experiment, which also exhibits the
290 lowest AMOC strength of all experiments.

291 The situation in the SH is more complex as surface temperature and sea ice evolution are
292 influenced both by freshwater forcing from the AIS as by the FWF in the NH. The AMOC
293 variability gives rise to changes in the SH through the so-called interhemispheric see-saw
294 effect (Stocker 1998). The SH begins to warm as the NH cools due to modified oceanic heat
295 transport across the equator. Minima in SH temperature (cf. Figure 5c) and maxima in SH sea
296 ice area (Figure 7d) are therefore associated with maxima in AMOC strength. The additional
297 effect of including GrIS freshwater forcing is consequently also felt in a warmer SH with less
298 sea ice formation. However, the influence of GrIS freshwater fluxes and consequential
299 AMOC variations on the SH temperature appears to be mostly limited to the beginning of the
300 experiment between ~135 and 131 kyr BP. It could be speculated that this is related to the
301 larger extent of the SH sea ice in a colder climate, making the system more sensitive due to an
302 increased potential for sea ice-albedo and insulation feedbacks. We also note that modelled
303 periods of increased NH freshwater fluxes, reduced AMOC strength and higher SH
304 temperatures are roughly in phase with periods of steeper increase in GHG concentrations (cf.
305 Figure 4b), in line with evidence from marine sediment proxies that indicate that CO_2
306 concentration rose most rapidly when North Atlantic Deep Water shoaled (Ahn and Brook,
307 2008). Since GHGs and NH freshwater fluxes are (independently) prescribed in our
308 experiments, the described in-phase relationship lends further credibility to our NH ice sheet
309 reconstruction.

310 The FWF from AIS melting (Figure 7f) increases the SO sea ice area (Figure 7d) by
311 freshening and stratifying the upper ocean waters, which in turn leads to lower surface
312 temperatures. In our experiments, the increased freshwater flux from the retreating AIS (cf.

313 noGfwf versus noAGfwf) between 131 and 129 kyr BP is in phase with a period of transient
314 AMOC strengthening (Figure 7b), which leads to a combined effect of surface cooling and
315 sea ice expansion in the SO.

316 The formation of AABW is strongly controlled by salinity and sea ice area (and therefore
317 temperature) of the polar surface waters and hence by both Antarctic and indirectly by NH
318 freshwater fluxes. Here the strength of AABW formation is calculated as the minimum value
319 of the global meridional overturning stream function below the Ekman layer south of 60° S.
320 The AABW formation (Figure 7e) is stronger for saltier and colder surface conditions and
321 therefore strongest in case noAGfwf, where FWF are suppressed from the AIS (saltier) and
322 the GrIS (colder). For a similar Antarctic freshwater forcing, the AABW formation is stronger
323 for a larger SH sea ice area. Including Antarctic FWF leads to a generally weaker AABW
324 formation as surface waters become fresher (cf. noGfwf versus noAGfwf). These
325 relationships imply also that a stronger decrease in AABW formation, associated with
326 decreased CO₂ uptake by the ocean can be found for periods of steeper increase in prescribed
327 radiative forcing. Again, this appears to support consistency in timing between prescribed
328 radiative and NH ice sheet forcing in our modelling.

329

330 **6 Temperature evolution in the Southern Hemisphere**

331 Millennial scale sea-surface temperature variations in the SH induced by NH freshwater
332 fluxes are the strongest in the SO, where anomalies can be amplified by sea ice-albedo and
333 insulation feedbacks. This is also the region that experiences the largest temperature change
334 due to FWF from the AIS itself (not shown).

335 In order to study the effect of Antarctic FWF in more detail, we also analysed the oceanic
336 temperature evolution south of 63°S (Figure 8). The effect of the AIS freshwater flux in the
337 reference experiment (compare noAGfwf with reference) becomes visible in the sea surface
338 temperature after 132 kyr BP (Figure 8a) as a cooling due to stratification and sea ice
339 expansion (Figure 8c). At the same time, the subsurface ocean warms (Figure 8b) as heat is
340 trapped under the stratified surface waters and expanding sea ice area. When the FWF decline
341 towards the end of the AIS retreat around 128 kyr BP, sea ice retreats again and the heat is
342 released to the atmosphere, where it generates an overshoot in SST compared to the
343 experiment with constant Antarctic freshwater fluxes (noAGfwf). The largest effect of this
344 heat buffering is found in winter in regions of strongest warming in the Bellingshausen Sea

345 and off the Gunnerus ridge adjacent to Dronning Maud Land. The maximum sea-ice extent in
346 the SH (Figure 8c) occurs at the time of largest surface cooling at 129.5 kyr BP. This
347 freshwater induced surface cooling at the onset of the LIG appears to be superficial and
348 relatively short lived and of clearly different signature compared to e.g. the Antarctic cold
349 reversal during the last deglaciation. The cooling event is indeed not recorded in our modelled
350 temperature evolution over central East Antarctica, in line with a lack of its signature in
351 Antarctic ice core records for that time period (Petit et al., 1999; EPICA community
352 members, 2004). A sea ice expansion during Termination II together with an *oceanic* cold
353 reversal around 129.5 kyr BP (Figure 8c) is however recorded in some deep-sea sediment
354 cores, where the composition of planktonic diatoms suggests meltwater as the primary cause
355 (Bianchi and Gersonde, 2002; Cortese and Abelmann, 2002).

356 As a further consequence, the timing of maximum annual mean surface air temperature
357 (defined as MWT for Maximum Warmth Timing; Bakker et al., 2013) in the SO differs by
358 several thousand years between experiments (Figure 9). Including Antarctic FWF leads to an
359 earlier MWT (by up to 2 kyr) in large parts of the SO south of 60°S and in the central and
360 eastern parts of the Atlantic sector of the SO up to 40°S (Figure 9d). Conversely, a later MWT
361 (by up to 3 kyr) is found in the Indian and Pacific sectors of the SO north of 60°S when
362 Antarctic FWF is accounted for (Figure 9d). In the reference experiment (Figure 9a) and
363 noGfwf (Figure 9b), the MWT lies relatively homogeneously between -129 kyr and -128 kyr
364 for the entire SO south of 45°S and coincides with the overshoot in SST after the peak input
365 of Antarctic FWF. The observed changes of the MWT in the SO due to the additional
366 Antarctic freshwater input can therefore in either way be understood as a shift towards the
367 time when heat from the mid-depth ocean buffer is released to the surface.

368

369 **7 Discussion**

370 Despite remaining uncertainties in the timing of ice sheet retreat during Termination II, we
371 find several lines of evidence in support of our ice sheet reconstructions and the associated
372 climatic signatures. The NH ice sheet reconstruction shows some similarity with the IRD
373 signal recorded in North Atlantic sediment cores (Kandiano et al., 2004; Oppo et al., 2006),
374 while the simulated evolution of the AMOC strength (Figure 7a) is in good agreement with a
375 recent reconstruction based on chemical water tracers (Böhm et al., 2015). The combination
376 of NH and SH sourced freshwater forcing variations produces a stronger decrease in AABW

377 formation, associated with decreased CO₂ uptake by the ocean for periods of steeper increase
378 in prescribed radiative forcing, in line with evidence from marine sediment proxies that
379 indicate that CO₂ concentration rose most rapidly when North Atlantic Deep Water shoaled
380 (Ahn and Brook, 2008). Reconstructing the NH ice sheet evolution during Termination II
381 with the same method but using the Grant et al. (2012) sea-level record for comparison with
382 Termination I has been shown to worsen agreement of the modelled climate with proxy
383 reconstructions (Loutre et al., 2014).

384 Our modelling results furthermore suggest that the major AIS retreat from its glacial
385 configuration could be constrained by an oceanic cold event recorded in several SO sediment
386 cores around Antarctica (Bianchi and Gersonde, 2002; Cortese and Abelmann, 2002). As a
387 schematic sensitivity test to uncertainties in the overall glacial AIS volume and retreat rate,
388 we have performed one more experiment identical to the reference experiment except for
389 Antarctic FWF scaled to 50% of their reference value. The resulting magnitude of the SO cold
390 event and overshoot is lower but exhibits the same timing and spatial expression as in the
391 reference case. The described mechanisms and effects can therefore be considered robust to
392 differences in the magnitude of the freshwater flux, resulting from uncertainties in glacial ice
393 volume or AIS retreat rate. Notably, the improved representation of the central East Antarctic
394 temperature evolution in the model when including Antarctic ice sheet changes (Figure 6) is
395 largely independent of the chosen freshwater forcing. This implies that changes in the
396 geometry of the ice sheet and modified atmospheric circulation patterns are the cause for the
397 stronger simulated temperature contrast.

398 The GrIS is generally assumed to have remained largely intact during the LIG (e.g. Robinson
399 et al., 2011; Colville et al., 2011; Stone et al., 2013; NEEM community members, 2013) and
400 indirect evidence of its freshwater contribution may be difficult to find due to the low
401 amplitude compared to the other NH ice sheets. However, recent ice core reconstructions of
402 the temperature evolution at the NEEM ice core site (NEEM community members, 2013)
403 point to a late retreat with a peak sea-level contribution close to 120 kyr BP. The GrIS can be
404 assumed to lose mass approximately as long as the temperature anomaly above the ice sheet
405 remains above zero. Based on the NEEM record, which has been used as forcing time series
406 in our stand-alone GrIS experiment, FWF from the GrIS peaks at ~125 kyr BP, but remains
407 elevated until around 120 kyr BP above the steady state background flux of an ice sheet in
408 equilibrium with the climate. The additional FWF from melting of the GrIS results in

409 relatively low temperatures over Southeast Greenland in response to a weakening of the
410 AMOC (not shown). The interaction between GrIS meltwater fluxes and oceanic circulation
411 hence give rise to a negative feedback on ice sheet retreat. This aspect could play an
412 important role for the stability of the southern dome of the ice sheet and should be examined
413 further with fully coupled climate-ice sheet simulations.

414 In general, the NH freshwater forcing leads to variations in the strength of the AMOC and
415 North Atlantic cooling and additionally through the bipolar see-saw effect, to temperature
416 changes in the SH. The only moment mid-depth ocean temperatures close to AIS grounding
417 lines are above pre-industrial values in our experiments is during the oceanic cold reversal
418 around 129.5 kyr BP, induced by anomalous FWF from the retreating AIS. During this
419 period, SO mid-depth temperature anomalies relative to the pre-industrial reach up to 0.3 K,
420 which could provide a positive but rather limited feedback on ice sheet retreat, similar to what
421 has been suggested by Golledge et al. (2014) for meltwater pulse 1A during Termination I.
422 However, the oceanic warming recorded in our model is not strong and the duration of the
423 perturbation does not appear to be long enough for a sustained impact on the retreat of the ice
424 sheet. Furthermore, the peak in freshwater flux appears when the ice sheet has already
425 retreated considerably and WAIS grounding lines are located mostly on the continental
426 shelves, more protected from the warm water build-up in the mid-depth ocean. A large-scale
427 marine ice sheet retreat of the likely less vulnerable EAIS sectors (Mengel and Levermann,
428 2014) appears particularly unlikely, given the atmospheric and oceanic forcing at the time
429 apparent in our modelling results. However, in-depth studies of these interactions require
430 detailed coupled simulations of the entire ocean-ice sheet system.

431 Despite aforementioned lines of evidence in support of the reconstructed NH ice sheet
432 evolution, a limitation to our modelling approach is the rescaling of ice sheet retreat during
433 Termination I, an attempt to address the sparseness of geomorphological field evidence for
434 Termination II. An alternative approach would be to physically model all ice sheets together
435 in one framework (e.g. de Boer et al. 2013), although spatial and temporal resolution of the
436 models is a limiting factor in that specific case. A rigorous modelling approach like the latter
437 could also help to prevent possible inconsistencies when combining ice sheet reconstructions
438 from different approaches. Nevertheless, any modelling approach will ultimately be
439 confronted with the same problem of scarce data for model validation during that period. The
440 exclusion of climate feedbacks on ice sheet evolution of our present one-way coupled

441 modelling approach is a general limitation, which we have addressed in a separate study with
442 a fully coupled model set-up (Goelzer et al., 2016).

443

444 **8 Conclusion**

445 We have presented a transient simulation of Termination II and the Last Interglacial period
446 with realistic ice sheet boundary conditions from reconstructed NH ice sheets and detailed
447 stand-alone simulations of the Greenland and Antarctic ice sheets. Our results show that the
448 temperature evolution at the onset of the Last Interglacial was in both hemispheres
449 considerably influenced by meltwater fluxes from the retreating ice sheets. While Antarctic
450 freshwater fluxes lead to strong perturbations of the Southern Ocean, NH freshwater fluxes
451 have an influence both on the Northern and SH temperature evolution through the oceanic
452 see-saw effect. The importance of additional freshwater input from the GrIS during
453 Termination II is small compared to the much larger fluxes from the other NH ice sheets and
454 becomes more important only later during the Interglacial when it is the only remaining ice
455 sheet contributing freshwater fluxes to the North Atlantic. In the SH, anomalous freshwater
456 input from the AIS leads to an episode of surface freshening, increased stratification and sea
457 ice cover and consequently reduced ocean heat loss to the atmosphere, with temporary heat
458 build-up in the mid-depth ocean. We argue that the surface ocean cooling associated with this
459 event may be used to constrain an early Antarctic retreat when matched with similar
460 signatures evident in some deep-sea sediment cores from the Southern Ocean.

461 Our transient simulations confirm results from earlier studies that stress the importance of ice
462 sheet boundary conditions for the climate evolution at the onset of the LIG. However, most of
463 the freshwater induced changes remain visible for at most 1-2 kyr after cessation of the
464 perturbations, indicative of a relative short memory of the (surface) climate system.
465 Additional effects may arise from climate-ice sheet feedbacks not considered in the present
466 model configuration, which should be investigated in fully-coupled experiments.

467

468 **Appendix A: Reconstruction of NH ice sheet forcing**

469 A direct reconstruction of NH ice sheet evolution during Termination II based on
470 geomorphological data is not possible, due to the scarcity of field evidence that was mostly
471 destroyed by the re-advancing ice sheets during the last glacial period. Therefore, a

472 reconstruction of Termination II is made by remapping the much better constrained ice sheet
473 retreat of Termination I.

474 **Ice extent during Termination I**

475 The evolution of the NH ice extent since the LGM was estimated based on published sources
476 (Table A1) dating back to the time of the NH ice sheet studies of Zweck and Huybrechts
477 (2003, 2005). For the large Laurentide and Eurasian ice sheets inferred ice extents are
478 relatively well determined from geomorphological data and the reconstruction remains in
479 good agreement with most recent sources (e.g. Hughes et al. 2016). For smaller ice sheets
480 such as the European Alps previous modelled ice extent was used (Zweck and Huybrechts,
481 2005).

482 For ice sheets with multiple sources of data the isochrones were merged using the most recent
483 source when conflicts occurred (e.g. Dyke et al. (2002) instead of Dyke and Prest (1987) for
484 the Innuitian ice sheet, Svendsen et al. (1999) instead of Andersen (1981) for the LGM
485 maximum of the Eurasian ice sheet). The most recent source was then used as a mask of
486 maximum ice extent for most recent isochrones of all sources. The only region, which
487 experienced an advance in ice extent using this technique was the southern Cordilleran ice
488 sheet according to the reconstruction of Clague and James (2002).

489 The INTCAL98 timescale of Stuiver et al. (1998) was used to convert radiocarbon dates to
490 calendar years for the sources in Table A1. The retreat of the ice sheets between the LGM and
491 PD was prescribed at 200 year resolution. Even for well-determined geomorphological
492 observations, uncertainties in dating and from the conversion of radiocarbon to calendar years
493 well exceed the 200 year temporal resolution used here. Figure A1 shows the deglaciation
494 chronology reconstructed in this manner.

495 **Ice sheet elevations during Termination I**

496 The NH ice sheets introduced significant changes to the surface topography of the region. As
497 LOVECLIM1.3 has only 3 atmospheric height levels, details regarding topography are not
498 strongly sensed. To include changes in surface topography in the model, parabolic profile ice
499 sheets are constructed using the extents shown in Figure A1, neglecting isostatic adjustment
500 (i.e. present-day surface elevation of the Earth's surface). The basal shear stress for the
501 parabolic profile reconstruction is chosen so that the difference in ice volume between LGM
502 and PD corresponds to 86 m of eustatic sea level change. With isostasy accounted for, a

503 similar elevation would result in an additional contribution of 24 m to a total equivalent
504 eustatic sea level change of 110 m (cf. Zweck and Huybrechts, 2005). Using this procedure
505 the maximum elevation of the Laurentide ice sheet is 3000 m near present-day Churchill in
506 Hudson Bay, and the maximum elevation of the Eurasian ice sheet is 2600 m 100 km west of
507 present-day Helsinki.

508 **Remapping Termination I to Termination II**

509 Remapping of the retreat during Termination I to Termination II is done using a benthic $\delta^{18}\text{O}$
510 record (Lisiecki and Raymo, 2005), assumed as an indicator of the global ice volume. In
511 practice, the NH ice sheet configuration for a given time (and $\delta^{18}\text{O}$ value) during Termination
512 II is taken from a time during Termination I when the $\delta^{18}\text{O}$ value had the same value. The
513 LGM sea-level contribution of the NH ice sheets relative to the present day of -110 m
514 translates into a similar magnitude for the penultimate glacial maximum (Lisiecki and Raymo,
515 2005). The resulting NH ice volume evolution for Termination II is shown in Figure 3a.
516 However, the method does not guarantee that the sea-level contribution of the reconstructed
517 NH ice sheets closely follows the global ice volume curve. This is generally due to the
518 mismatch between global ice volume and NH ice sheet reconstruction during Termination I,
519 and in part related to the unconstrained contribution of other components (AIS, thermal
520 expansion). Due to the assumed analogy, different configurations of the NH ice sheets (e.g.
521 Obrochta et al., 2014) and different relative timing of NH and SH deglaciation between last
522 and penultimate glaciation are not represented in these reconstructions. NH freshwater fluxes
523 were estimated from the same method by using derived volume changes as input to a
524 continental runoff routing model (Goelzer et al., 2012b) to identify the magnitude and
525 location of meltwater fluxes to the ocean.

526

527

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536

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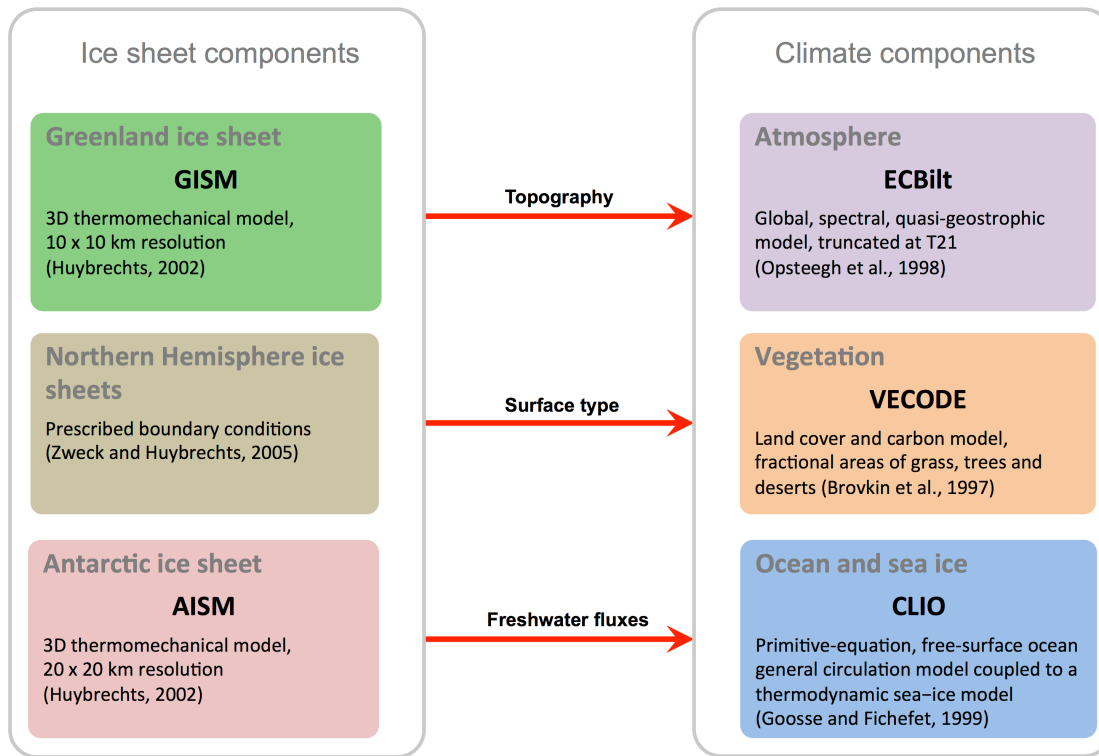
794 **11 Tables**

795

796 **Table 1: Matrix of all experiments and the respective ice sheet components that evolve in time (yes) or are**
 797 **fixed (no). In the latter case, freshwater fluxes (FWF, grey) are kept constant and topography and surface**
 798 **albedo are fixed to the preindustrial configuration.**

EXP	topo NH	FWF NH	topo GrIS	FWF GrIS	topo AIS	FWF AIS
Reference	yes	yes	yes	yes	yes	yes
noGfwf	yes	yes	yes	no	yes	yes
noAGfwf	yes	yes	yes	no	yes	no
noAG	yes	yes	no	no	no	no
noIS	no	no	no	no	no	no

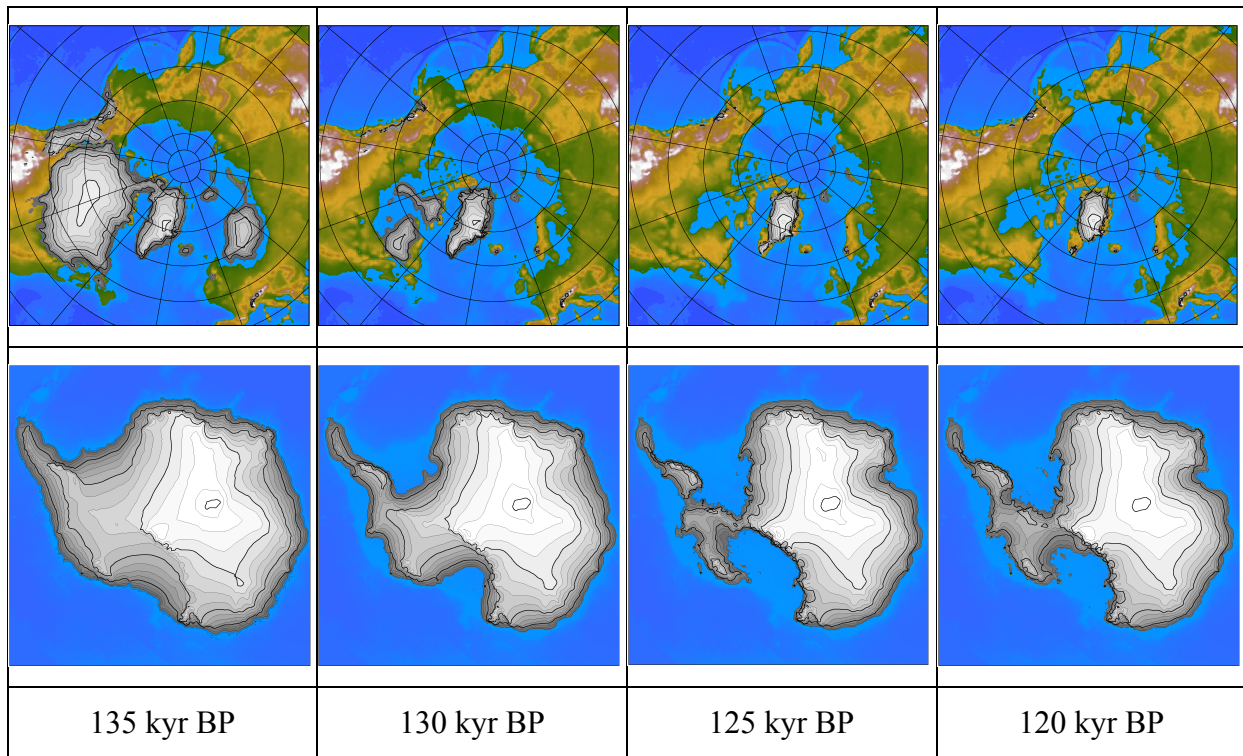
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802 **Figure 1: LOVECLIM model setup for the present study including prescribed ice sheet boundary**
 803 **conditions from the Northern Hemisphere, Greenland and Antarctic ice sheets.**

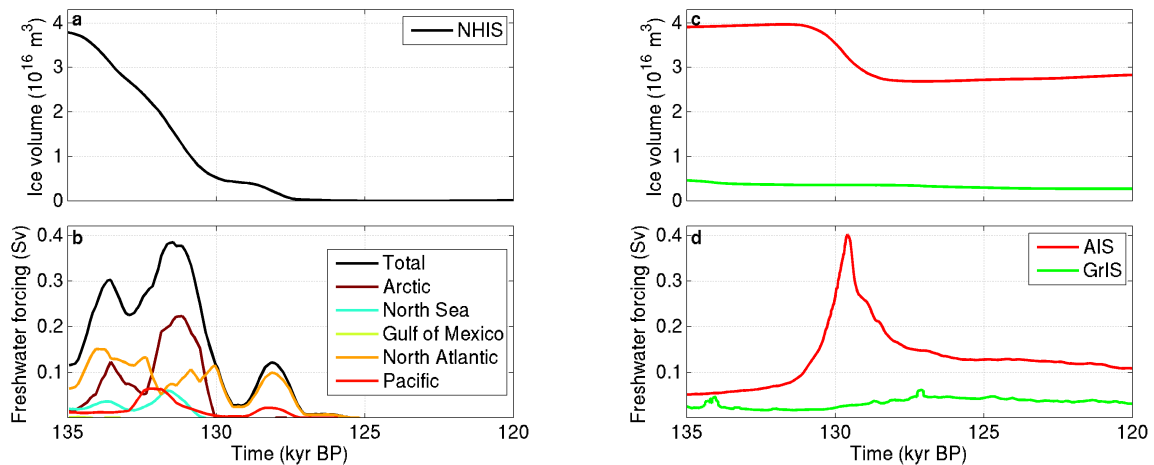
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806 **Figure 2: Evolution of reconstructed Northern Hemisphere ice sheets and embedded modelled GrIS (top)**
 807 **and modelled AIS (bottom) used as boundary conditions for the climate model.**

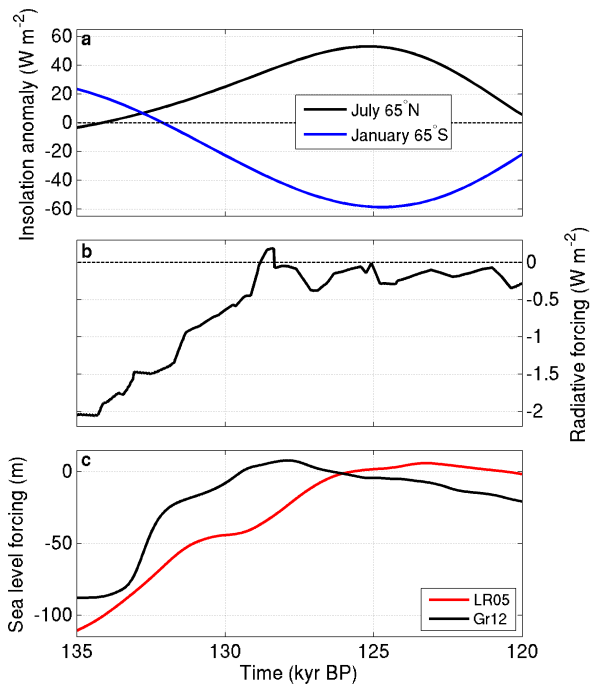
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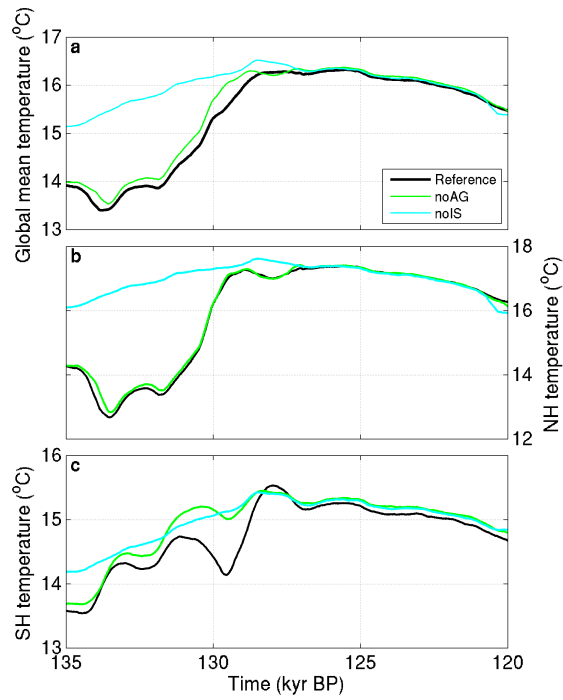


810 **Figure 3: Reconstructed ice volume (a, c) and freshwater forcing (b, d) from the NH ice sheets (left) and**
 811 **from the GrIS and AIS (right). See Goelzer et al. (2012b) for definition of oceanic basins in panel b.**

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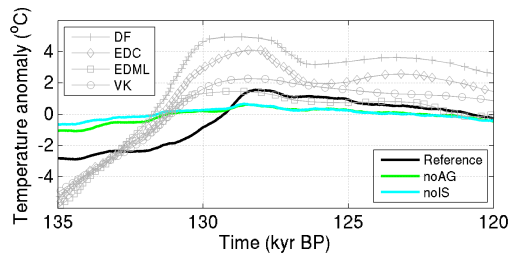


813 **Figure 4: Prescribed model forcings. (a) Average monthly insolation anomaly relative to the pre-industrial**
 814 **at 65° North in July (black) and 65° South in January (blue). (b) combined radiative forcing anomaly of**
 815 **prescribed greenhouse gas concentrations (CO₂, CH₄, N₂O) relative to the pre-industrial. (c) sea-level**
 816 **forcing for the ice sheet components derived from either oceanic δ¹⁸O data (Lisiecki and Raymo, 2005,**
 817 **red) scaled to a global sea-level contrast between LGM and present day of 130 m or derived from a Red**
 818 **Sea relative sea-level record (Grant et al. 2012, black).**
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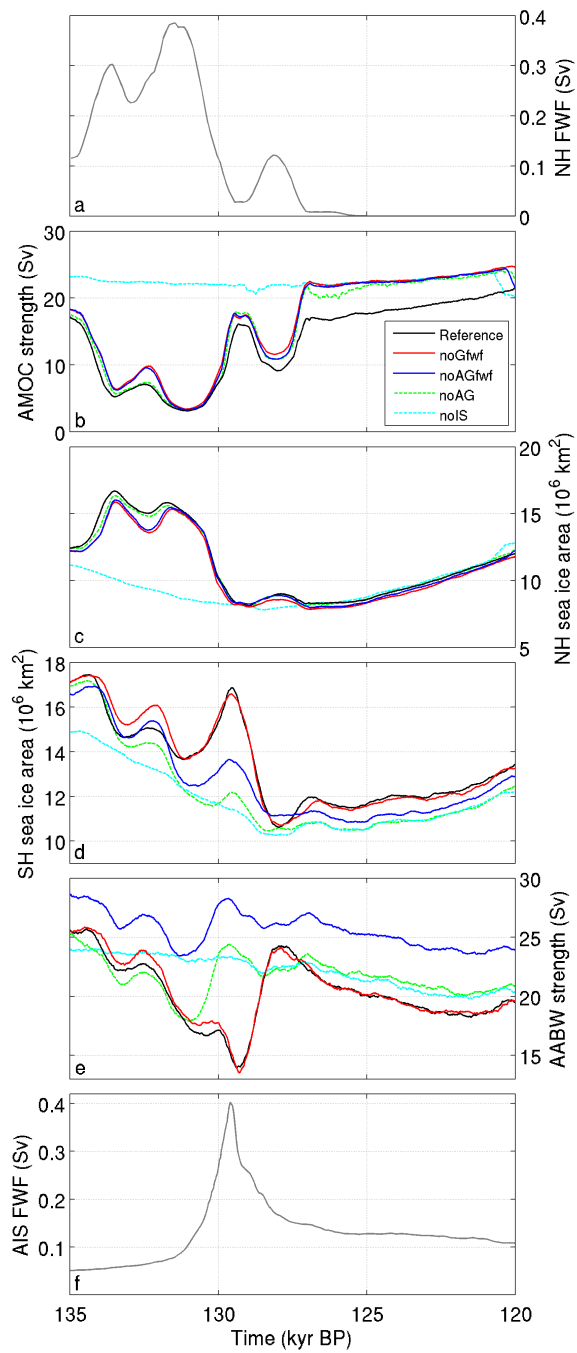
821 **Figure 5: Evolution of global mean (a), northern (b) and southern (c) hemispheric mean surface**
 822 **temperature for experiments with different ice sheet forcing included. Curves are smoothed with a**
 823 **running mean of 200 years for better comparison.**

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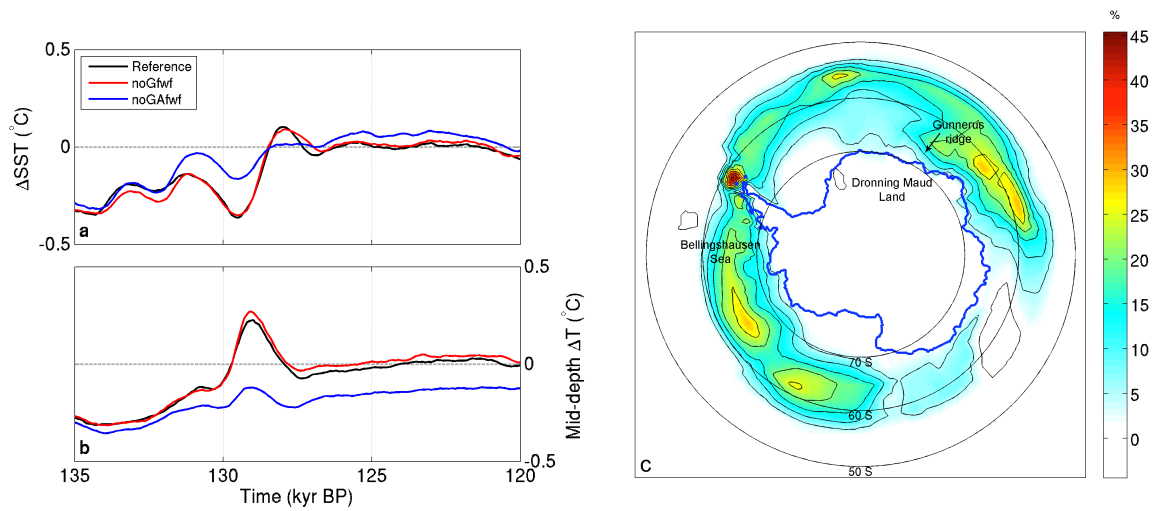


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826 **Figure 6: Comparison of modelled East Antarctic temperature evolution with reconstructed temperature**
 827 **changes at deep ice core sites. Modelled temperature anomalies are averaged over a region 72° - 90° S and**
 828 **0° - 150° E. Ice core temperature reconstructions for the sites EPICA Dronning Maud Land (EDML,**
 829 **75°00' S, 00°04' E), Dome Fuji (DF, 77°19' S, 39°40' E), Vostok (VK, 78°28' S, 106°48' E) and EPICA**
 830 **Dome C (EDC, 75°06' S, 123°21' E) are from Masson-Delmotte et al. (2011).**



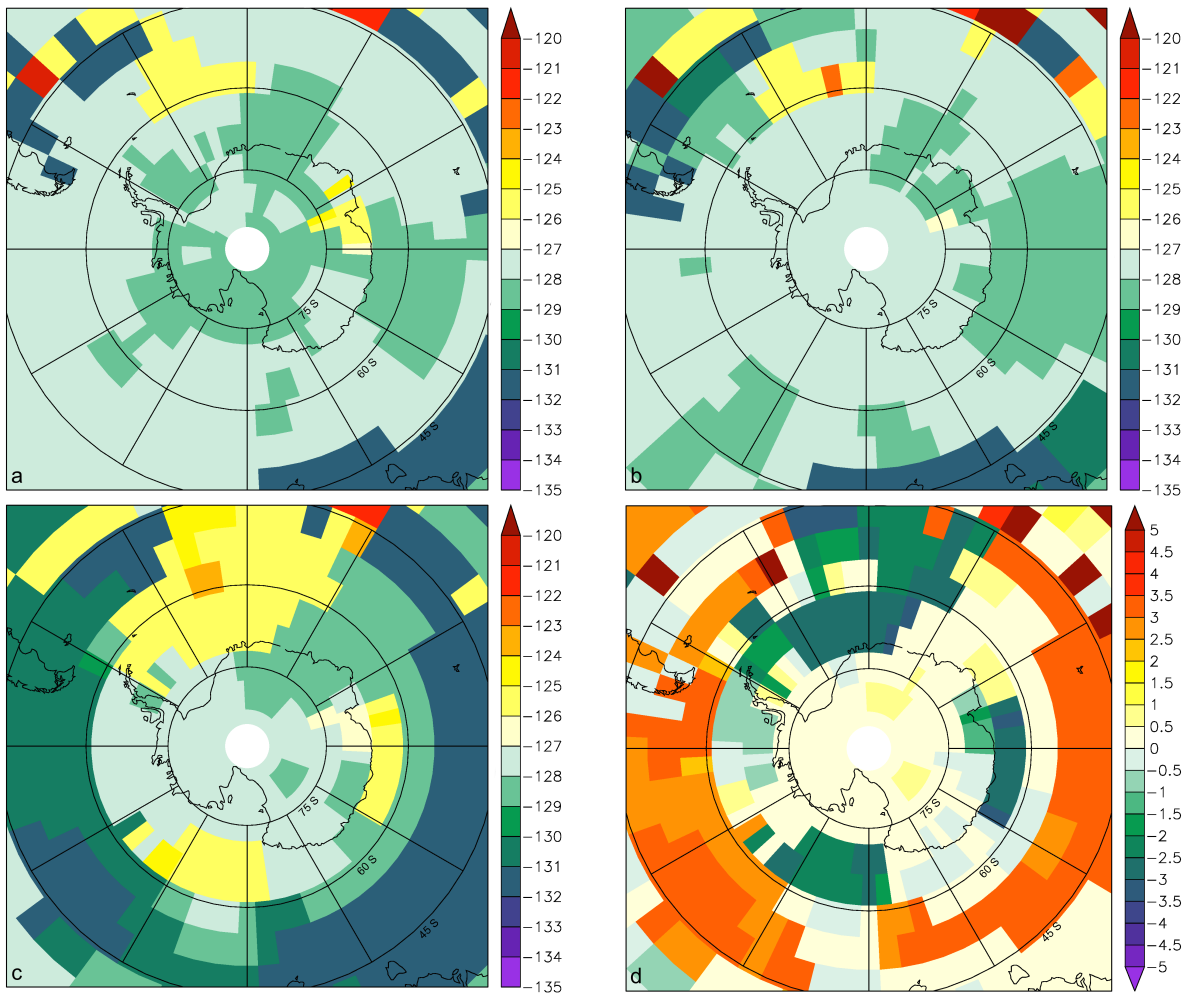
831 **Figure 7: Freshwater forcing and oceanic response characteristics. NH (a) and Antarctic ice sheet**
 832 **freshwater fluxes (f), strength of the AMOC (b), NH sea ice area (c), SH sea ice area (d) and strength of**
 833 **AABW formation (e) for the different experiments with and without freshwater forcing from Greenland,**
 834 **Antarctic and NH ice sheet melting. Curves are smoothed with a running mean of 200 years for better**
 835 **comparison.**
 836



837 **Figure 8: Evolution of annual mean sea surface temperature (a) and mid-depth (485-700 m) ocean**
 838 **temperature (b) anomalies relative to the pre-industrial in close proximity to the AIS (south of 63°S). (c)**
 839 **Meltwater related changes in annual mean sea ice area at 129.5 kyr BP from differences between**
 840 **experiments Reference and noAGfwf in per cent. The blue contour outlines the observed ice-shelf edge**
 841 **and grounded ice margin of the present-day AIS for illustration. All curves (a, b) are smoothed with a**
 842 **running mean of 200 years for better comparison.**

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Figure 9: Time of maximum surface air temperature (MWT) in kyr BP for experiments Reference (a), noGfwf (b) and noAGfwf (c) and difference in MWT between experiments noGfwf and noAGfwf (d) in kyr, showing the shift of the MWT when Antarctic freshwater fluxes are included.

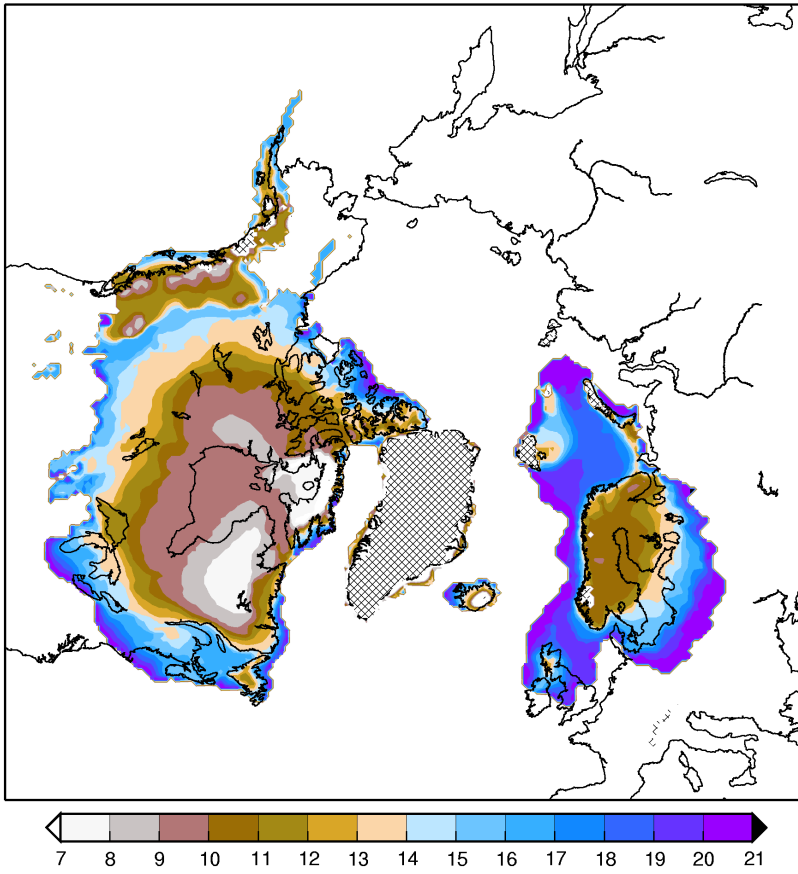
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851 **Table A1: Sources of geomorphological data or modelling results used to prescribe changes in Northern**
 852 **Hemisphere ice sheet extent for the retreat during Termination I.**

Ice Sheet	Source	Isochrone Time Period (kyr BP)
Laurentide	Dyke and Prest (1987)	18 – Present Day
Innuitian	Dyke et al. (2002)	18
Cordilleran	Clague and James (2002)	20 – Present Day (south)
	Dyke et al. (2002)	
	Mayewski et al. (1981)	18 (north) 21 – 7 (interior)
Iceland	Andersen (1981)	20 – Present Day
Eurasian	Andersen (1981)	20 – Present Day
	Landvik et al. (1998)	15 – 12 (Barents Sea)
	Mangerud et al. (2002)	18 (Southern Barents and Kara Seas)
	Svendsen et al. (1999)	18
European Alps	Zweck and Huybrechts (2003)	21 – Present Day (modelled ice extent)

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Figure A1: Interpolated ice sheet extent during the last deglaciation for the Northern Hemisphere ice sheets as a function of time (kyr BP). Hatched regions indicate present-day ice.