How might the North American ice sheet influence the Northwestern Eurasian climate?

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1 Abstract

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It is now widely acknowledged that past Northern hemisphere ice sheets covering Canada and 3 northern Europe at the last glacial maximum (LGM) exerted a strong influence on climate by 4 causing changes in atmospheric and oceanic circulations. In turn, these changes may have 5 impacted the development of the ice sheets themselves through a combination of different 6 7 feedback mechanisms. The present study is designed to investigate the potential impact of the 8 North American ice sheet on the surface mass balance (SMB) of the Eurasian ice sheet driven by simulated changes in the past glacial atmospheric circulation. Using the LMDZ5 9 atmospheric circulation model, we carried out twelve experiments under constant LGM 10 conditions for insolation, greenhouse gases and ocean. In these experiments, the Eurasian ice 11 sheet is removed. The twelve experiments differ in the North American ice-sheet topography, 12 ranging from a white and flat (present-day topography) ice sheet to a full-size LGM ice sheet. 13 This experimental design allows the albedo and the topographic impacts of the North American 14 15 ice sheet onto the climate to be disentangled. The results are compared to our baseline 16 experiment where both the North American and the Eurasian ice sheets have been removed. In summer, the sole albedo effect of the American ice sheet modifies the pattern of planetary 17 waves with respect to the no-ice sheet case, resulting in a cooling of the northwestern Eurasian 18 region. By contrast, the atmospheric circulation changes induced by the topography of the 19 North American ice sheet lead to a strong decrease of this cooling. In winter, the Scandinavian 20 and the Barents-Kara regions respond differently to the American ice sheet albedo effect: in 21 22 response to atmospheric circulation changes, Scandinavia warms and total precipitation is more abundant whereas the Barents-Kara area cools with decrease of convective processes causing a 23 decrease of total precipitation. The gradual increase of the altitude of the American ice sheet 24 leads to less total precipitation and snowfall and to colder temperatures over both the 25 Scandinavian and the Barents-Kara Seas sectors. We then compute the resulting annual surface 26 mass balance over the Fennoscandian region from the simulated temperature and precipitation 27 fields used to force an ice-sheet model. It clearly appears that the SMB is dominated by the 28 ablation signal. In response to the summer cooling induced by the American ice-sheet albedo, 29 high positive SMB values are obtained over the Eurasian region, leading thus to the growth of 30 an ice sheet. On the contrary, the gradual increase of the American ice sheet altitude induces 31 32 more ablation over the Eurasian sector, hence limiting the growth of Fennoscandia. To test the robustness of our results w.r.t the Eurasian ice sheet state, we carried out two additional LMDZ 33 34 experiments with new boundary conditions involving both the American (flat or full-LGM) and a high Eurasian ice sheets. The most striking result is that the Eurasian ice sheet is maintained 35 under full-LGM North American ice-sheet conditions, but loses ~10 % of its mass compared to 36 the case in which the North American ice sheet is flat. These new findings qualitatively confirm 37 the conclusions from our first series of experiments and suggest that the development of the 38 Eurasian ice sheet may have been slowed down by the growth of the American ice sheet, 39 offering thereby a new understanding of the evolution of Northern hemisphere ice sheets 40 throughout glacial-interglacial cycles. 41

42 **1. Introduction**

The last million years is characterized by alternating glacial and interglacial phases. During 43 glacial periods, large ice sheets covering present-day Canada and northwestern Eurasia (Dyke 44 45 and Prest 1987; Clark et al. 1993; Peltier 2004; Svendsen et al. 2004; Lambeck et al. 2006; Tarasov et al. 2012) exerted a strong influence on climate. Several studies highlighted the 46 47 importance the climatic changes induced by continental-scale ice sheets (e.g. Clark 1999). It has been recognized that the ice-sheet topography is likely to be the main factor altering the 48 49 atmospheric circulation in the Northern hemisphere (Broccoli and Manabe 1987; Pausata et al. 2011). As an example, the pioneering study carried out by Manabe and Broccoli (1985) with an 50 51 atmospheric general circulation model shows that the North American ice sheet caused a split 52 of the westerlies. The authors also highlight a larger amplitude of the planetary waves due to 53 the presence of ice sheets. Recently, Ullman et al. (2014) tested the influence of different North American ice-sheet reconstructions on the climate. They showed that a higher American ice 54 sheet leads to a more zonal Atlantic jet, thereby confirming the key role of ice-sheet topography 55 on atmospheric circulation. Using model outputs from the Paleo Modelling Intercomparison 56 Project 2 (PMIP2) (Braconnot et al. 2007), Laîné et al. (2009) and Rivière et al. (2010) 57 highlight a strengthening and an equatorward displacement of the sub-tropical jet-stream during 58 the Last Glacial Maximum (LGM) w.r.t. the pre-industrial period. Changes in the position and 59 the strength of the North Atlantic jet stream induce changes in the storm tracks, and therefore in 60 precipitation (Hall et al. 1996; Kageyama and Valdes 2000; Laîné et al. 2009; Rivière et al. 61 2010). These changes also have an influence on the energy transport and can therefore modify 62 63 the temperature.

The climatic changes induced by large-scale ice sheets exert an influence on both temperature 64 and precipitation that drive the ice-sheet surface mass balance (SMB), defined as the sum of 65 snow accumulation and snow and ice ablation. Using a simple ice-sheet model based on an 66 idealized geometry coupled to a stationary-wave model, Roe and Lindzen (2001a,b) underline 67 the importance of accounting for the feedbacks between ice sheets and temperatures induced by 68 69 changes in stationary waves to properly simulate the evolution of an ice sheet. They show that the self-induced temperature anomaly due to an ice sheet leads to a warming over the ice-sheet 70 71 western part. This may explain the absence of ice over Alaska at the LGM. They also suggest 72 that the stationary waves excited by the North American ice sheet may have contributed to a warming over Europe, influencing the development of the Eurasian ice sheet. In the same way, 73 with a three-dimensional stationary wave model, Liakka et al. (2011) showed that the position 74

of the southern margin of ice sheets strongly depends on the temperature anomalies due to 75 stationary waves, which are modified by the ice sheet itself. More recently, with the use of the 76 77 CAM3 atmospheric model run under four different climatic contexts (last interglacial, marine isotope stages 5b and 4 (MIS5b and MIS4), and LGM periods), Löfverström et al. (2014) 78 79 showed how the atmospheric circulation changes induced by the ice sheets could have influenced the growth of the ice sheets themselves. Similarly to Roe and Lindzen (2001a), they 80 show that the summer atmospheric circulation change due to the presence of ice sheets may 81 cause a sufficient warming over Siberia and Alaska to inhibit ice growth. They also perform 82 83 two experiments under MIS4 conditions to test the influence of one ice sheet on the other. They conclude that the summer temperature anomaly induced by the presence of the MIS4 American 84 85 ice sheet is too weak to explain the small size of the European ice sheet at that period, but may have contributed to the westward shift of the ice-sheet mass centre. Although several studies 86 87 have been devoted to the mutual influences of the ice sheets (Roe and Lindzen 2001a; Beghin et al. 2014; Löfverström et al. 2014), no study has specifically investigated the mechanisms 88 89 through which the American ice sheet may have influenced the European climate, and therefore the European ice-sheet surface mass balance. 90

The aim of this paper is to investigate the atmospheric response to changes in the North 91 American ice sheet boundary conditions and to examine how the related atmospheric 92 circulation changes impact the climate over the Scandinavian and Barents-Kara Seas regions. 93 To achieve this goal, we use the LMDZ5 atmospheric general circulation model run with 94 different thicknesses of the North American ice sheet taken as boundary conditions, with no ice 95 at all over Eurasia. We investigate the mechanisms by which the American ice sheet may 96 change the surface mass balance of the Eurasian ice sheet. Besides gradually increasing the 97 altitude of the North American ice sheet, the originality of the present study also relies on ice-98 sheet model simulations. We use climatic fields simulated by the LMDZ5 model as inputs to a 99 three-dimensional ice-sheet model to compute the surface mass balance of the Eurasian ice 100 101 sheet. The description of the climate and ice-sheet models is given in section 2 as well as the experimental design. The model results, presented in sections 3 and 4 are followed by a 102 103 discussion (section 5). Section 6 summarizes the main findings of our study.

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107 **2. Model and experiment**

108 2.1 The atmospheric model

109 The LMDZ5tmospheric model used in this study belongs to the LMDZ model family developed at Laboratoire de Météorologie Dynamique (Sadourny and Laval 1984; Le Treut et al. 1994; Le 110 Treut et al. 1998, Hourdin et al., 2006). LMDZ is the atmospheric component of the IPSL-111 CM5A coupled ocean-atmosphere model (Dufresne et al. 2013). The dynamical equations are 112 113 discretized on a longitude-latitude-staggered Arakawa C-grid (Kasahara 1977). The model ensures the conservation of both enstrophy (square of wind rotational) for baroclinic flows 114 (Sadourny 1975b; Sadourny 1975a) and the axi-symmetric component of the angular 115 momentum. The model version used in this study has 39 vertical levels and runs on a 96 x 95 116 model grid resolution (3.75° x 1.9°). A complete description of the model can be found in 117 (Hourdin et al. 2013). 118

119 2.2 The ice-sheet model

120 The GRISLI ice-sheet model is a three-dimensional thermo-mechanical model which simulates the evolution of ice-sheet geometry (extension and thickness) and the coupled temperature-121 122 velocity fields in response to climate forcing. A comprehensive description of the model can be found in (Ritz et al. 2001) and (Peyaud et al. 2007). Here, we only summarize the main 123 124 characteristics of this model. The equations are solved on a cartesian grid (40 km x 40 km). Over the grounded part of the ice sheet, the ice flow resulting from internal deformation is 125 126 governed by the shallow-ice approximation (Hutter 1983; Morland et al. 1984). The model deals with ice flow through ice shelves using the shallow-shelf approximation (MacAyeal 127 1989). It also predicts the large-scale characteristics of the ice streams using criteria based on 128 the effective pressure and hydraulic load. At each time step, the velocity and vertical profiles of 129 temperature in the ice are computed as well as the new geometry of the ice sheet. The isostatic 130 adjustment of bedrock in response to ice load is governed by the flow of the asthenosphere, 131 with a characteristic time constant of 3000 years, and by the rigidity of the lithosphere. The 132 temperature field is computed both in the ice and in the bedrock by solving a time-dependent 133 134 heat equation. The surface mass balance is defined as the sum between accumulation and ablation computed by the empirical positive degree-day (PDD) method (Reeh 1991; Fausto et 135 al. 2009). This method assumes that melt rates of snow and ice are linearly related to the 136 number of PDD through degree-day factors for snow and ice materials (Braithwaite 1984; 137 1995). 138

139 **2.3 Experimental set-up**

In order to investigate the albedo and topography effects induced by the North American ice 140 sheet on the Eurasian climate at the LGM, we carried out 12 simulations under LGM conditions 141 142 for greenhouse gases (GHG), insolation, sea-surface temperatures and sea-ice). In this series of simulations, the altitude of the North American ice sheet surface ranges from that of the 143 144 present-day surface to 100 % of that used in the PMIP3 LGM experiments. The PMIP3 LGM ice sheets result from a combination of three reconstructions, namely ICE-6G v2.0 (Toscano et 145 146 al. 2011), GLAC-1 (Tarasov et al. 2012) and ANU (Lambeck 2001). The way this new reconstruction has been obtained is explained in detail on the PMIP3 website 147 148 (https://wiki.lsce.ipsl.fr/pmip3/doku.php/pmip3:design:pi:final:icesheet) and in Abe-Ouchi et al. (2015). In our baseline experiment (noIS), the land-ice mask is modified (w.r.t. PMIP3) to 149 150 remove both the European and the American ice sheets and the altitude is that of the presentday (Fig. 1). In the other simulations, we only remove the European ice-sheet mask. These 151 simulations are referred to as xxdhL, where dhL represents the surface height difference 152 between the PMIP3 LGM Laurentide ice sheet and the present-day surface, and xx represents 153 the percentage of dhL taken into account. To simplify the writing, the American and Eurasian 154 ice sheets are respectively referred to as "Laurentide ice sheet" (LIS) and the "Fennoscandian 155 ice sheet" (FIS) in the following. 156

157 The topography of the 00dhL experiment is therefore the same as today, but the land-ice mask is set to the LGM one (Figs. 1 and 2). Greenland and Antarctic ice sheets are the same as in the 158 PMIP3 experimental set-up. Both insolation and GHG boundary conditions are similar to those 159 160 defined in the PMIP3 protocol: the orbital forcing is taken at 21 kyr BP from (Berger 1978), while atmospheric GHG concentrations are those recorded in Antarctic ice cores ($CO_2 = 185$ 161 162 ppm, (Lüthi et al. 2008); $CH_4 = 350$ ppb, (Loulergue et al. 2008); $N_2O = 200$ ppb, (Spahni et al. 2005)). The LGM land-sea mask is also taken into account, with closed Bering Strait, and land 163 in Hudson Bay and Barents Sea (Fig. 1). The sea-surface temperatures (SST) and sea-ice cover 164 come from 50 consecutive years of the IPSL PMIP3 LGM run outputs (Kageyama et al. 2013). 165 These oceanic surface conditions are depicted in Figure 1. In the North Atlantic area, the 166 temperature meridional gradient is steep in the western part of the basin and weakens in its 167 168 eastern part, as for the pre-industrial period. The sea-ice edge is located South of Iceland in 169 winter and retreats to the northern limits of the Nordic Seas in summer. These rather warm North Atlantic / Nordic Seas conditions are due to a strong Atlantic meridional overturning 170 circulation (Kageyama et al. 2013). Using these boundary conditions, the LMDZ5 simulations 171

- are run for 50 years, and we study the last 40 years to be at equilibrium. Two-sided Student t-
- tests with a p-value = 0.05 (von Storch and Zwiers, 2003) are performed on key diagnostics.

174 **3. AGCM results**

175 In this section, we focus on how the LIS influences the climatic fields which can have an 176 impact on the surface mass balance of an ice sheet. The key climatic variables used to compute 177 the amount of ablation with the PDD method are the monthly temperatures and the monthly 178 total (solid plus liquid) precipitation.

179 **3.1** The northwestern Eurasian temperature

180 **3.1.1 Impact of the Laurentide ice sheet**

Using the comparison between the 00dhL and the noIS experiments, we first consider the 181 impact of the LIS albedo on the 2-meter summer air temperature over the Scandinavian-182 Barents-Kara (SBK) region (Fig. 3, 00dhL experiment, left column). The most striking feature 183 is a cooling over the whole Eurasian continent with a maximum (more than 10 °C) centered 184 over the Barents-Kara Seas and a moderate cooling (2 to 4°C) over Scandinavia. The area 185 experiencing the largest cooling also corresponds to the area with the largest albedo (orange 186 contours). This is a direct consequence of the LGM land-sea mask used in this study for which 187 188 the Barents-Kara Seas are represented as continental areas. This favors snow accumulation and therefore higher albedo values in this area, causing in turn a positive feedback on temperatures. 189 190 Figure 3 also displays the 2m mean summer temperature over the SBK region for different 191 altitudes of the LIS. Here, we only give results of selected experiments to simplify the presentation of our results. It appears that the SBK cooling gets smaller when the LIS gets 192 193 higher. Overall, the simulated summer surface air temperature evolves together with the LIS altitude, with a progressive reduction of the zonal and meridional extent of the largest cooling 194 195 zone. When the LIS reaches its full LGM size (100dhL- noIS), the summer cooling over Barents-Kara Seas is reduced to 4-6°C and to 2°C at most over Scandinavia. The topography of 196 197 the LIS has therefore a warming effect on the SBK region for the summer season.

In winter, the response of the temperature to the LIS albedo and its evolution with the LIS altitude are different from the summer response (Fig. 3, right panels). First, the LIS albedo induces warmer temperatures (+3 °C) over South Scandinavia, the British Isles and present-day North and Baltic Seas, and cooler temperatures over an area in the Barents Sea expanding from Svalbard to Novaya Zemlya (3 °C). Surprisingly, in the 30dhL experiment, the amplitude of these temperature changes (w.r.t. noIS) is reduced, showing that the response of winter temperatures to changes in LIS altitude is far from being linear. Beyond 60dhL, the cooling zone progressively expands, and spreads across the entire British Isles and Scandinavia-Barents-Kara area when the LIS has its full LGM size. In contrast to the climatic response simulated for the summer season, the LIS topography leads to colder surface air temperatures in winter.

To understand the origin of these contrasted responses, we investigate the processes through which the changes in the altitude of the Laurentide ice sheet modify the atmospheric circulation.

212 **3.1.2** Atmospheric circulation processes

The comparison of the zonal anomaly (00dhL vs. noIS) of the 500 hPa summer geopotential 213 height shows that the sole albedo effect of the LIS is sufficient to drastically change the 214 atmospheric circulation (Fig. 4, 00dhL and noIS, left column). The ridge over the Rockies, 215 clearly visible in noIS, disappears in 00dhL, and the trough over the Labrador sector is more 216 217 extensive. Ridges over the North Atlantic, the Greenland ice sheet and northern Europe are more developed in 00dhL (w.r.t. noIS). By contrast, the troughs over Iceland, the Norvegian 218 219 and Barents Seas appearing in noIS are weaker or even vanish in 00dhL. In noIS, the trough over Svalbard implies southerlies over the Barents Sea (Fig. 4, right column) and therefore, 220 221 warm temperatures in this area (Fig. 3). The weakening of this trough in 00dhL as well as the stronger ridge over Greenland induces northerlies over Svalbard and the Barents Sea (Fig. 4) 222 and therefore colder temperatures (Fig. 3). 223

The Labrador trough and the North Atlantic ridge become stronger when the LIS gets higher, 224 until it reaches 60 % of its full size. Beyond 60dhL, the Labrador trough and the Atlantic ridge 225 keep more or less the same amplitude. When the LIS gets higher, the Greenland ridge gets 226 weaker, the European ridge remains almost unchanged (Fig. 4), and there is a return of the 227 Iceland-Svalbard-Barents trough. This new trough, centered between Greenland and Svalbard, 228 along with the weakening of the Greenland ridge, brings southerlies again over the northern 229 Barents Sea. The Barents winds shift from northerlies to southerlies when the LIS is 60 % of its 230 231 full size. The southerlies bring warmer temperatures, explaining the SBK warming with the increase of the LIS height (Figs 3 and 4). 232

In winter (Fig. 5), the sole albedo effect of the LIS does not drastically change the geopotential 233 anomaly. However, the meridional wind field shows a slight decrease of the southerlies over 234 the Barents Sea, explaining the slight cooling over this region. There is also a slight decrease of 235 236 the northerlies simulated over the British Isles, the eastern North Atlantic Ocean and southern Scandinavia, consistent with the moderate warming simulated over this region. In the 30dhL 237 experiment, a strengthening of the southerlies is observed (Fig. 5) inducing a warming (w.r.t. 238 00dhL) over the SBK region (Fig. 3) and explaining the non-linear behavior on winter 239 temperatures. When the altitude of the LIS further increases, there is a progressive emergence 240 241 of a trough over the Scandinavian-Barents sector. This trough comes with northerlies over the Norvegian Sea and weaker southerlies over the Barents sea (Fig. 5, right column), consistently 242 243 with colder temperatures over this area, as seen in Fig. 3.

This analysis explains why opposite sign temperature responses are obtained for the summer and winter seasons. Since ablation is rather sensitive to the summer season, more ice is expected over the Fennoscandian area when the sole albedo effect of the LIS is operating (i.e. 00dhL experiment). However, the development of an ice sheet in this area also depends on the amount of snow accumulation. Therefore, we next examine hereafter the LIS impact on the SBK total precipitation and snowfall.

250 **3.2 Total and solid precipitation over northwestern Eurasia**

251 **3.2.1 Impact of the Laurentide ice sheet**

In this section, we investigate the impact of the LIS on the mean annual precipitation and 252 snowfall over northwestern Eurasia (Fig. 6). In response to the LIS albedo effect (00dhL-noIS 253 panel) a slight excess of precipitation is simulated over the northeastern North Atlantic and the 254 southern and western parts of Scandinavia, but less precipitation is simulated over the Barents-255 Kara Sea area. However, over the whole SBK region, the snowfall anomaly is positive. This 256 257 increase in snowfall is due to the summer signal (not shown) related to the cooling observed over this region (Fig. 3). This cooling is strong enough to allow for an increase of the fraction 258 259 of solid precipitation and to maintain the snow cover.

When the LIS reaches its full LGM altitude (100dhL-noIS), the positive anomaly of total precipitation over Scandinavia simulated in the 00dhL experiment shifts southward, reaching even the French and the Iberian Peninsula Atlantic coasts. This results in dry conditions over the western flank of the Scandinavian region, both in terms of total and solid precipitation. Concomitantly, the negative anomaly of total precipitation over the Barents-Kara Sea region simulated in 00dhL expands westward leading to a drier climate and to a decrease of snowfall over Scandinavia. These results suggest that the Barents-Kara Seas and the Scandinavian regions are sensitive to different characteristics of the LIS: whereas the precipitation over Scandinavia appears to be sensitive to the LIS height, the precipitation anomaly over the Barents-Kara Seas mainly appears to result from the LIS albedo effect and is rather insensitive to height.

271 **3.2.2 Large-scale vs convective precipitation**

To disentangle the mechanisms responsible for precipitation changes in these regions we split 272 273 up the total precipitation into its large-scale and convective components (Fig. 7). The largescale component behaves similarly to the total precipitation shown in Figure 6 with a positive 274 275 anomaly in the 00dhL experiment simulated over Scandinavia and the eastern North Atlantic 276 sector and the southward shift of this pattern simulated when the LIS gets higher. Hence the drier signal over the western flank of Scandinavia obtained with the full-LGM LIS (100dhL-277 noIS) can also be found in the large-scale component of the total precipitation. This strongly 278 suggests that the Scandinavian precipitation, as expected, is driven by large-scale processes. 279

The negative anomaly of precipitation appearing in the Arctic Ocean north of the Barents-Kara 280 281 area (Fig. 6) is also clearly visible in the large-scale component (Fig. 7); by contrast, the large negative precipitation anomaly south of the Novaya Zemlya clearly arises from convective 282 283 processes. Over the Barents-Kara area, the origin of the precipitation signal is less clear but seems to be a combination of both convective and large-scale processes. The decrease of 284 convective precipitation negative anomaly (w.r.t noIS) over the entire SBK sector when the LIS 285 is higher can be easily explained by the weakening of the summer cooling discussed in 286 section 3.1. 287

288 To further investigate the mechanisms governing precipitation over Scandinavia and the Barents-Kara area, we examine the seasonality of precipitation. We first compute the average 289 precipitation above Scandinavia (55-70°N: 0-20°E) (Fig. 8a) and the Barents-Kara region (65-290 90°N : 20-100°E) (Fig 8b), then we compare summer and winter precipitation anomalies with 291 the annual precipitation anomaly. Results displayed in Figure 8 show that over Scandinavia 292 293 winter precipitation changes are clearly correlated with annual precipitation (r = 0.99). This means that the annual signal observed over Scandinavia (Figs 6 and 7) is mainly due to winter 294 295 precipitation. Conversely, the Barents-Kara sector is rather governed by summer precipitation

296 (at least for LIS altitudes up to 60 % of the PMIP3 one), as shown by the positive correlation 297 (r = 0.98) obtained for the summer season (Fig. 8b).

298 **3.2.3** Precipitation changes over Scandinavia: implication of the jet stream shift

We therefore look for a mechanism prevailing in winter to explain the Scandinavian precipitation differences. Using the PMIP3 model outputs for the LGM and the PI time slices, Beghin et al. (accepted) found a correlation between the southward shift of the North Atlantic jet stream from the pre-industrial to the LGM periods and the winter precipitation changes over the Iberian Peninsula. Following the same approach, we examine whether a similar mechanism to the one found in (Beghin et al. accepted) may explain the precipitation changes over Scandinavia when the LIS characteristics (albedo and topography) are modified.

306 Similarly to Chavaillaz et al. (2013), we use the 850 hPa zonal wind to define the jet stream. The position of the jet displayed in Fig. 9 corresponds to the position of the 850 hPa meridional 307 308 wind maximum. The LIS albedo effect induces a slight northward shift (w.r.t. noIS) of the North Atlantic jet stream (00dhL-noIS panel). As the LIS becomes higher, the jet moves 309 310 progressively southward and shifts southward of its noIS position between 30dhL and 60dhL. From 60dhL and beyond, the Scandinavian winter precipitation anomaly becomes more and 311 312 more negative. This suggests a relationship between the shift of the North Atlantic jet stream and the precipitation over the Scandinavian region. 313

314 To confirm this assumption, we plot the shift of the jet $(\Delta \phi)$ as a function of the precipitation difference between the noIS and the xxdhL experiments over Scandinavia (Fig. 10). $\Delta \phi$ is 315 defined as the difference of latitudes between noIS and xxdhL experiments where the 850 hPa 316 317 zonal wind is at its highest value. The latitude of this maximum has been found by computing the zonal mean of the zonal wind over the Atlantic basin ($50^{\circ}W - 10^{\circ}W$). To minimize the 318 uncertainties due to model discretization, we used a parabolic interpolation around the grid 319 point where the mean zonal wind is largest. The maximum of the corresponding parabola yields 320 a better approximation of the exact latitude of the maximum jet wind. This method allows to 321 find more accurately the latitude of the jet maximum (Chavaillaz et al. 2013; Beghin et al. 322 accepted). We find a good linear correlation between the shift of the North Atlantic jet and the 323 Scandinavian precipitation anomaly during the winter season (Fig. 10). The quantification of 324 the range of the precipitation anomalies and of the jet shift is obtained by bootstrapping: a 325 sample of thousand values is obtained by calculating the average of 1000 randomly picked 326 327 samples of 40 year-duration in the noIS and the xxdhL original samples. Linear regressions of 328 the 5th and 95th percentiles of the bootstrapping sample confirm the close link between the 329 shift of the North Atlantic jet stream and the winter precipitation changes over Scandinavia due 330 to the effect of LIS.

331 The 500 hPa geopotential height (z500) zonal anomaly over the East Laurentide sector (90°W-60°W) does not change significantly when a flat LIS is imposed (Figs 5, 11). As the LIS gets 332 333 higher, the Labrador trough expands southward: figure 11 shows that the zero isoline is located at ~43°N in the noIS and 00dhL experiments and progressively reaches 37°N for a full-LGM 334 LIS. Over the Atlantic sector (50°W-10°W) the Icelandic low / Azores high pattern remain 335 similar in all experiments up to 60 dhL and only show a significant change in the full LIS 336 experiment (Figs 5, 11). Indeed, in 100dhL, the Icelandic low's amplitude weakens but 337 simultaneously expands southward (Fig. 11). Consistently with these changes in the z500 zonal 338 339 anomaly, the 500hPa westerlies shift southward over the whole Atlantic sector (Fig; 11, bottom). Again, the largest southward shift for this variable is between the 60dhL and 100dhL 340 experiments, showing the non-linearity of the atmospheric circulation to the LIS height. 341

342 **3.2.4 Precipitation changes over the Barents-Kara area**

As previously mentioned in section 3.2.2, the mean annual precipitation changes over the 343 344 Barents-Kara area are dominated by the summer signal when the LIS remains at a relatively low level (i.e. up to 60dhL). This result is confirmed by the comparison of the patterns of 345 346 annual (Fig. 6), summer and winter (Fig 9) precipitation anomalies. The amount of summer 347 precipitation is at its lowest level when the sole albedo effect of the LIS is accounted for. It increases with the LIS altitude in response to warmer temperatures (Fig; 3). Conversely, the 348 summer snowfall amount decreases from 00dhL to 100dhL due to the progressive lowering of 349 the fraction of solid vs liquid precipitation. 350

351 **3.3 Summary**

To summarize, the LIS albedo effect induces colder summer temperatures over the entire northwestern Eurasian region (SBK area) w.r.t. the noIS experiment. Over South Scandinavia, the winter temperatures (Fig. 3) are warmer due to weaker northerlies (Fig. 5). More precipitation is also simulated in winter in response to the poleward shift of the North Atlantic jet stream (Fig. 6). As the Barents-Kara sector is dominated by the summer signal, less precipitation is simulated in response to colder summer temperatures (w.r.t. noIS), while a larger snowfall amount is obtained. When the LIS gets higher the negative anomaly of summer temperatures over the Barents-Kara region becomes weaker due to changes in the atmospheric circulation. In addition, the higher the LIS, the smaller the amounts of annual precipitation and snowfall over the Scandinavian and the Barents-Kara areas, in response to the equatorward shift of the jet.

4. What is the influence of the LIS on the growth of the FIS ?

364 4.1 In case of no initial Fennoscandian ice sheet

What are the impacts of the LIS on the surface mass balance of both Scandinavian and Barents-365 366 Kara areas? To investigate this question, we use the LMDZ5 monthly temperatures and precipitation fields of each experiment to force the ice-sheet model GRISLI. LMDZ5 climatic 367 fields are downscaled over the GRISLI grid using a bilinear interpolation. Due to the difference 368 of resolution between the atmospheric and the ice-sheet models, temperature is vertically 369 corrected using a linear vertical gradient of 6 K/km. Precipitation is also vertically corrected 370 using an exponential function of the temperature (Charbit et al. 2002; Charbit et al. 2007). 371 Snowfall is recalculated using the downscaled precipitation and temperature because the more 372 detailed topography of GRISLI allows snowfall when LMDZ5 provides only liquid 373 precipitation. 374

Figure 12 displays the snow accumulation and the ablation fields computed by GRISLI for the five selected experiments. As expected from the LMDZ5 results of annual precipitation and snowfall, snow accumulation over Scandinavia is larger when the LIS is flat and decreases as the LIS is higher. A similar observation can be made over the Barents-Kara region. Consistently with the simulated summer temperatures (Fig. 3), when the LIS is flat, the ablation is weaker over the entire SBK region, and increases as the LIS gets higher. Nevertheless, even with the full-LGM LIS, the ablation remains weaker than in the noIS experiment.

The resulting surface mass balance is shown in Fig. 13a. The similarities between the ablation 382 383 (Fig. 12b) and the surface mass balance patterns indicate that the surface mass balance is 384 dominated by the ablation. In the absence of LIS, the surface mass balance is positive over only a small part of the SBK area, namely over Svalbard. Under these conditions, no ice sheet grows 385 386 (Fig. 13b). When the LIS is flat, the surface mass balance is positive over the Barents-Kara Seas and over the northern part of Scandinavia, allowing the growth of ice. The resulting FIS is 387 less extended than those provided by the ICE-6G (Toscano et al. 2011) and the ANU (Lambeck 388 2001) reconstructions. This is likely due to the absence of the albedo feedback, since our 389 390 approach is not based on a full coupling method and does not allow to account for the impact of ice-sheet changes onto the climate. As the LIS gets higher, the limit of the positive surface mass balance shifts westward and northward, excluding the Kara Sea and Scandinavia from the positive surface-mass balance area. When the LIS has its full LGM size, the surface mass balance is positive only over the Svalbard. As a consequence, the simulated FIS is smaller.

4.2 What if the Fennoscandian ice sheet is initially high?

The aim of this study is not to provide a fully realistic scenario of the evolution of the ice sheets 396 397 throughout an entire glacial cycle. Rather, based on idealized cases, it is to investigate the impact of the Laurentide ice sheet on the Eurasian ice sheet surface mass balance. In the first 398 399 part of this study, we have shown that the growth of the LIS prevents the development of the 400 FIS. However, back to a more realistic context in which both ice sheets co-evolved, the 401 relevance of this conclusion may be questioned and the issue of the FIS inception under full 402 LGM LIS conditions may appear somehow artificial. Thus, to make our previous conclusions more robust, we perform two additional LMDZ experiments with new ice-sheet boundary 403 conditions. In both simulations, the FIS has the elevation computed by the ice-sheet model for 404 the 00dhL run (i.e. the largest FIS elevation of all the GRISLI experiments), while the LIS is 405 either "flat" (00dhLFIS) or has its full LGM altitude (100dhLFIS). Here, our objective is to test 406 407 the self-consistency of the overall study, by checking that the growth of the FIS is slowed down when the LIS gets higher. We first examine how a high FIS modifies the northwestern Eurasian 408 climate under 00dhL and 100dhL conditions (i.e. with flat or full LGM LIS ice sheet), in terms 409 of the temperature and precipitation fields used in our modelling approach to compute the ice-410 sheet surface mass balance. We then use the climatic outputs to force the ice-sheet model. 411

412 **4.2.1 Response of the atmosphere**

413 Figure 14 displays the summer surface air temperature, which is the main driver of the ablation, for the 00dhL (flat LIS, no FIS), 00dhLFIS (flat LIS, high FIS), 100dhL (high LIS, noIS) and 414 415 100dhLFIS (high LIS, high FIS) experiments. The comparison between Figs 14a and 14b and Figs 14c and 14d illustrate the impact of the Fennoscandian ice sheet on the 2m summer 416 temperature. As expected, the presence of Fennoscandia induces a large cooling over the ice 417 sheet itself, due to the temperature-altitude relationship and to the albedo effect. In previous 418 sections, we have shown that the cooling observed over the Barents-Kara region is weakened 419 when the LIS gets higher due to a modification of the large-scale atmospheric circulation (Figs 420 3-4). A similar mechanism is observed when the Fennoscandian ice sheet is high (100dhLFIS 421 experiment): the cooling over the ice sheet is reduced and the warming along the southern 422

423 margin is amplified (Fig. 14d vs Fig. 14b). These features result from an increase of warm air 424 advection in the Barents-Kara sector and to a decrease of northerly heat flux in the south. The 425 presence of the FIS also induces a split of the North Atlantic jet stream in two branches 426 resulting in northward deflected heat fluxes along the western margin of the ice sheet and the 427 Barents Sea sector (Figs 15b,d and 16b,d). The heat advection modulates the ice-sheet altitude 428 and albedo effects but is not sufficient to cancel out the cooling observed over the ice sheet.

Figure 17 displays maps of snow accumulation differences (after downscaling on the GRISLI 429 grid) between the four experiments computed by GRISLI after downscaling. The presence of 430 the Fennoscandian ice sheet results in an increase of precipitation and snow accumulation over 431 its western flank, along the Norwegian Sea coast (Figs 17a,c). This is due to the barrier effect of 432 the ice sheet and to the subsequent increase of the vertical wind. This effect is less pronounced 433 434 when the LIS is high (Fig. 17c) as a result of the southward displacement of the North Atlantic jet stream (Figs 6,11). A second striking feature is the increase of snow accumulation in the 435 436 southern part of the Eurasian ice sheet with respect to the no FIS situation. This is mainly due to colder temperatures over the ice sheet, increasing the fraction of total precipitation which falls 437 438 as snow. In case of a flat LIS, this accumulation zone is confined to the southern margin and drier conditions (w.r.t no FIS case) in the interior of the Eurasian ice sheet are observed, mainly 439 located over the Kara Sea (Fig. 17a), and likely due to the strong barrier effect. By contrast, 440 when the LIS is high (Fig. 17c), differences between FIS and no FIS experiments are larger and 441 expand over more than the southern half of the FIS. This coincides with a larger temperature 442 difference over the FIS area when the LIS is high (Figs 14c,d) compared to the low LIS 443 situation (Figs 14a,b). 444

445 As mentioned above (section 3.2.1 and Fig. 6), in the absence of FIS, the total precipitation is 446 driven by large-scale processes over the Scandinavian region and by convection over the 447 Barents-Kara sector. When the LIS is high, the southward shift of the winter jet stream (Fig. 9) and the increase of summer temperature lead to an overall decrease of snowfall over the 448 Eurasian region. When the FIS is present, the situation is rather different and convective 449 precipitation is negligible (not shown) due to the cold temperatures induced by the presence of 450 the FIS. The differences between the 00dhLFIS and 100dhLFIS experiments are mainly located 451 over the western and southeastern parts of the ice sheet with respectively larger and smaller 452 snowfall amounts simulated in the 00dhLFIS experiment (w.r.t 100dhLFIS). The excess of 453 454 precipitation (total and solid) in the flat LIS/high FIS case is a direct consequence of the

455 increase of both zonal and meridional humidity transports in winter (not shown) favoured by456 the enhanced strength of the 850 hPa horizontal winds (Fig 18).

457 **4.2.2** Consequences on the simulated Eurasian ice sheets

The climatic outputs simulated in the 00dhLFIS and 100dhLFIS experiments have been used to 458 compute the surface mass balance of the Eurasian ice sheet at the GRISLI resolution and to 459 force the ice-sheet model with the downscaling procedure described in section 4.1. Starting 460 from the initial state provided by the 00dhL GRISLI run with flat (00dhLFIS) or full LGM LIS 461 (100dhLFIS), GRISLI has been integrated during 100 kyr. At the end of the simulation, the 462 simulated ice volume and ice covered area simulated in the 00dhLFIS experiment are 463 respectively 10 and 11 % greater than those obtained with the 100dhLFIS simulation. The 464 differences in simulated ice thickness between both simulations are displayed in figure 19 465 466 (00dhLFIS - 100dhLFIS). In most parts of the simulated ice covered areas, ice thickness differences are fully consistent with the snow accumulation signal (Fig. 17). In particular, large 467 positive anomalies (> 500 m) are found in the southwestern sector, that are greater than 1000 m 468 near the ice sheet margin. In the northern part, lower ice thicknesses obtained in the 00dhLFIS 469 experiment correspond to the area of lower snow accumulation (w.r.t 100dhLFIS), as ablation 470 is very small due to summer temperatures largely below 0°C. Surprisingly, another large 471 472 positive difference is found in the eastern part which is characterized by a lower amount of 473 snow accumulation in case of flat LIS. (w.r.t full LIS). This pattern results from warmer temperatures simulated in the southern and eastern parts of the 100dhLFIS Eurasian ice sheet in 474 response to the decrease of cold air advection (see section 4.2.1). Warmer temperatures found 475 476 in the south and east of the FIS when the LIS has its full LGM size favour the westward and the northward retreats of the eastern and southern margins of the ice sheet. Since our approach does 477 478 not account for the feedback of the simulated ice sheets on the climate, the amplitude of the different responses are likely to be underestimated due to an underestimation of the albedo 479 effect. However, this analysis confirms that the growth of the FIS is slowed down when the LIS 480 is high and demonstrates the self-consistency of the overall study. 481

482 **5. Discussion**

483 **5.1. Glaciation scenario**

Geomorphological reconstructions revealed that the Fennoscandian ice sheet eastern limit
reached the Putorana Mountains in Siberia during the Early Weichselian (90-80 kyr BP)
(Mangerud et al. 1998; Svendsen et al. 2004). According to Svendsen et al. (2004), this eastern

487 limit moved back to the Kara sea around 60 kyr BP, along with the southward extension of the Scandinavian ice sheet (Andersen and Mangerud 1989; Lundqvist 1992; Mangerud et al. 1998; 488 489 Houmark-Nielsen 1999). At the LGM, the ice sheet finally reached the British Isles, and the 490 Kara ice shelf did not cover the entire Kara sea anymore (Landvik et al. 1998; Ehlers et al. 491 2004; Svendsen et al. 2004). More recently, Kleman et al. (2013) combined geological evidences and numerical modelling to reconstruct the paleo-topography of past Northern 492 hemisphere ice sheets during MIS5b (~ 86 kyr), and MIS4 (~64 kyr) and 2 (LGM). They found 493 that the LIS and the FIS grow in concert in the early part of the last glacial cycle (i.e. until 494 495 MIS5b), while after MIS4, their evolution is radically different with a further growth of the LIS and a significant shrinkage of the FIS due to its southwestward retreat towards the main 496 497 precipitation sources.

498 Löfverström et al. (2014) propose to attribute the westward shift of the FIS and the decrease of the Kara ice sheet to warm temperatures over the southeastern part of the Fennoscandian 499 500 region, induced by the FIS itself, along with the upslope precipitation effect proposed by Sanberg and Oerlemans (1983). In their study, the LIS exerts only a small influence on the 501 502 retreat of the Kara ice sheet. While Löfverström et al. (2014) used realistic LIS topographies provided by the Kleman's et al. (2013) reconstructions, we used idealized geometries to study 503 separately the impact of the LIS albedo and topography on climate. The results of our first 504 series of xxdhL experiments clearly show that a large but low LIS favors a positive SMB over 505 the SBK sector (Fig. 13, left panels) and therefore the FIS maintenance (Fig.13, right panels). 506 By contrast, as the LIS gets higher, ablation increases leading to a smaller FIS. Back to a more 507 508 realistic glacial context in which both North American and Eurasian ice sheets exist (xxdhLFIS experiments), we also simulate a net decrease of snow accumulation under full LGM LIS 509 conditions (w.r.t flat LIS) over the Kara sector combined with an increase over the western part 510 of FIS (Fig. 17d), consistently with the upslope precipitation effect proposed by Sandberg and 511 Oerlemans (1983). The large positive SMB values over the SBK sector when the LIS is low 512 513 combined with the westward displacement of the Eurasian ice sheet as the LIS gets higher are fully in line with the ice-sheet paleo-topography reconstructions provided by Kleman et al. 514 515 (2013) for the last glacial cycle. The remote interactions highlighted in the present study 516 between the LIS and the FIS systems also suggest that the ablation increase over the Barents-Kara Seas when the LIS has its LGM size may have contributed to the LGM FIS small size. 517 Conversely, the larger size of the FIS (Svendsen et al. 2004) during the late Saalian (around 518 519 140 ka) w.r.t. the LGM may have been related to a smaller LIS.

520 The remote interactions highlighted in the present study between the LIS and the FIS systems 521 also suggest that the larger size of the FIS (Svendsen et al. 2004) during the late Saalian 522 (around 140 ka) w.r.t. the LGM may have been related to a smaller LIS.

523 5.2 Oceanic effects

524 Our experimental setup has been designed to investigate how the atmospheric circulation 525 changes in response to the LIS albedo and topography effects may impact the climate over the 526 Eurasian region and therefore the evolution of the surface mass balance of the Fennoscandian 527 ice sheet. We have therefore neglected the potential impact of the ocean.

First of all, all experiments presented in this study have used the LGM land-sea mask and 528 529 therefore the Barents-Kara region is considered as a continental area. The consequences of this assumption on the simulated large-scale atmospheric circulation are likely to be negligible. 530 However, accounting for a marine Barents-Kara area would decrease the positive albedo 531 532 feedback in this region, and thus the FIS inception (see section 3.1.1). Our experimental setup uses LGM SST and sea ice boundary conditions provided by the coupled atmosphere-ocean 533 IPSL model (Kageyama et al. 2013). This prevents us from studying the conditions that could 534 have prevailed with a marine Barents-Kara area, because this sector was covered by the FIS in 535 536 this run, and future studies will be necessary to investigate further the relationships between the LIS and the FIS within more realistic contexts. 537

Secondly, the potential impacts of oceanic circulation changes have also been ignored. Yet, 538 539 changes in atmospheric circulation related to the presence of continental ice sheets may induce large modifications of the oceanic circulation. For example, Wunsch (2006) proposed a 540 541 scenario in which phase shifts of the stationary waves caused by the development of ice sheets 542 modify the wind direction and strength leading to abrupt changes in temperature and 543 precipitation fields (and thus to climatic events having similar features as Dansgaard-Oeschger 544 events). These changes eventually produce large-scale modifications of the oceanic circulation. In turn, any change in the large-scale oceanic circulation may influence the evolution of the 545 Northern hemisphere ice sheets as well as their mutual interactions through the atmosphere-546 ocean coupling. 547

Recently, using the Earth system model COSMOS, Zhang et al. (2014) found that gradual increase of the Northern hemisphere ice sheet height is able of causing abrupt transitions of the AMOC through a positive atmosphere-ocean-sea-ice feedback resulting in an adjustment of heat distribution and sea ice coverage in the North Atlantic ocean. In this study, abrupt AMOC transitions from a weak to a strong mode, occurring for intermediate ice sheet sizes, are

accompanied by a noticeable increase of surface air temperatures in the northeastern North 553 Atlantic. The warming effect simulated in our study over the Eurasian region when the LIS gets 554 higher could be therefore amplified by the North Atlantic warmer temperatures. However, the 555 mechanism highlighted in Zhang et al. (2014) is primarily driven by the simulated northward 556 shift of the westerlies when the Northern hemisphere ice sheet altitude gradually increases. On 557 558 the contrary, we simulate a southward displacement of the jet stream when the height of the Laurentide ice sheet increases. Also, using the IPSL-CM4 model, Pausata et al (2011) showed 559 that the position of the North Atlantic jet stream is not substantially influenced by sea-surface 560 561 temperatures or sea-ice coverage. In the same way, using the CAM3 atmospheric model coupled to a slab ocean, Löfverström et al. (2014) did not find any significant changes in the 562 563 simulated temperature patterns over Alaska, Asia and the North Atlantic sector when two different oceanic heat transports corresponding to interglacial and LGM conditions are used. 564

The diversity of these models' responses to changes in the LIS characteristics show that they are 565 566 highly model-dependent. This underlines the need for investigating more in-depth the links 567 between large-scale oceanic and atmospheric circulations with different models run under 568 climatic conditions rather different from those of the present-day periods. In fact, the recent inter-comparison study of the North Atlantic jet stream in LGM vs pre-industrial conditions in 569 570 the PMIP3 models (Braconnot et al. 2012) carried out by Beghin et al. (accepted) clearly shows 571 that the response of the jet stream to LGM conditions is highly model-dependent. The range of model results found in this latter study is confirmed by the range of various results found in the 572 recent literature. Using the outputs of four comprehensive atmosphere-ocean models, Laîné et 573 al. (2009) found a southward displacement of the North Atlantic jet stream under LGM 574 conditions. A similar conclusion can be drawn from the study of Ullman et al. (2014) carried 575 out with the GISS model and two different LIS topographies as LGM ice-sheet boundary 576 577 conditions: the difference of the simulated jet speed obtained with these different LIS 578 reconstructions (see their figure 2) shows a more southward position of the jet when the LIS is 579 higher. Similarly to our study, Löfverström et al. (2014) found an equatorward position of the winter zonal winds and a more zonal jet under LGM conditions (w.r.t. MIS5b and MIS4). On 580 581 the contrary, Roberts et al. (2014) found that a more zonal structure is simulated with the FAMOUS model after the occurrence of a Heinrich event, and thus after a decrease of the 582 583 Laurentide ice thickness. This indicates a southward shift of the North Atlantic jet stream when the LIS altitude is lower, and, by contrast, a northward displacement when it is higher. All these 584 585 contrasted responses call for further studies to determine why the results are so model-586 dependent.

587 6. Conclusions

The goal of this study was to investigate the atmospheric processes through which the LIS can 588 influence the northwestern Eurasian climate, and therefore how it can influence the FIS surface 589 590 mass balance. We performed twelve simulations under LGM conditions, but with different icesheet configurations. In our baseline experiment, the LIS and the FIS are removed. In the first 591 592 set of simulations, we only removed the FIS, and the LIS has different heights. In the final series of experiments, we considered the impact of the LIS height on a high FIS. We therefore 593 594 identified the effect of the LIS albedo and topography on the northwestern European climate. We focused on the fields having an impact on ice-sheet surface mass balance, i.e. the 595 596 temperature and the precipitation. We showed that:

597 1/ The LIS albedo acts to decrease the summer temperatures over the Fennoscandian sector. 598 This decrease is amplified by positive snow albedo feedback. This positive feedback is 599 somewhat related to our experimental setup, which imposes land over Barents-Kara instead of 600 water, allowing snow maintenance. The temperature decrease is weaker when the LIS is higher, 601 due to atmospheric circulation changes.

602 2/ In winter, the LIS albedo impact decreases the Barents-Kara temperature, but increases the
603 Scandinavian temperature. The higher the LIS, the colder the temperature over the whole
604 Fennoscandian sector.

605 3/ The effect of LIS albedo also tends to shift the North Atlantic jet stream poleward, and to 606 bring more precipitation and snowfall over Scandinavia. When the LIS is higher, the Atlantic 607 jet shifts equatorward, bringing less precipitation over northern Europe. As a consequence, 608 precipitation and snowfall decrease over Scandinavia and Barents-Kara sectors when the LIS is 609 higher.

610 4/ The analysis of the GRISLI simulations forced by the LMDZ precipitation and temperature 611 fields shows that the LIS albedo favors the growth of the FIS, essentially because of weak 612 ablation. As the LIS gets higher, the FIS gets smaller due to more ablation, in accordance with 613 the LMDZ5 summer temperature response. When the LIS reaches its full-LGM size, ice 614 remains over Svalbard only.

5/ Within the context of glacial-interglacial cycles, the FIS inception under full-LGM LIS
conditions appears to be unrealistic. Therefore, we performed additional simulations with a
high FIS and flat and full-LGM LIS. The results of these new sensitivity experiments show a

618 similar response to the first series of experiments: in the presence of a full LIS, the FIS 619 undergoes a significant mass loss. This confirms that a large LIS counteracts the growth of the 620 FIS although it is still maintained, such as it does within more realistic contexts, and highlights 621 therefore the self-consistency of the overall study.

This study highlights the mechanisms by which the LIS can influence the FIS surface mass 622 balance. The use of the ice-sheet model allows to quantify the impact of climatic changes due to 623 the LIS presence on the FIS. It shows that the albedo of the LIS favors the FIS growth whereas 624 the LIS topography acts against the FIS development. This conclusion is reinforced by the 625 results obtained with a high FIS. It is worth noting that without a transient simulation of the last 626 glacial period, it is not possible to accurately elaborate the lead-lag relationships of the 627 Northern hemisphere ice sheets during the last glacial period. However, based on our results, a 628 629 plausible glaciation scenario can be proposed with an early phase of FIS glacial inception favored by the areal expansion of the LIS, and followed by a limited growth of the FIS with a 630 progressive westward retreat of the FIS itself as the LIS altitude gradually increased. Such 631 proposed relationships between ice-sheet build-up demonstrate that this study also provides 632 633 new insights in the understanding of the various Northern hemisphere ice sheet configurations during the different glacial-interglacial cycles. 634

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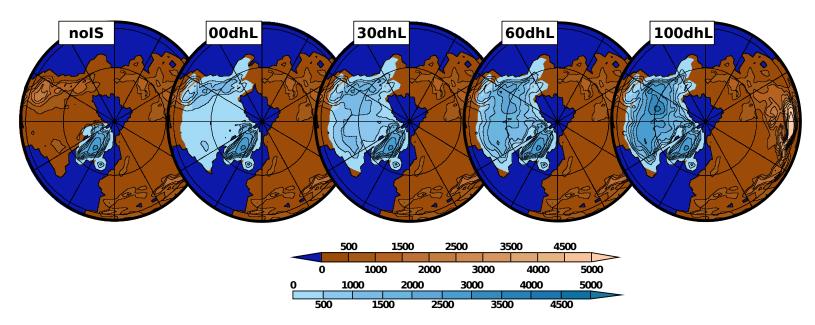
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820 Figure captions

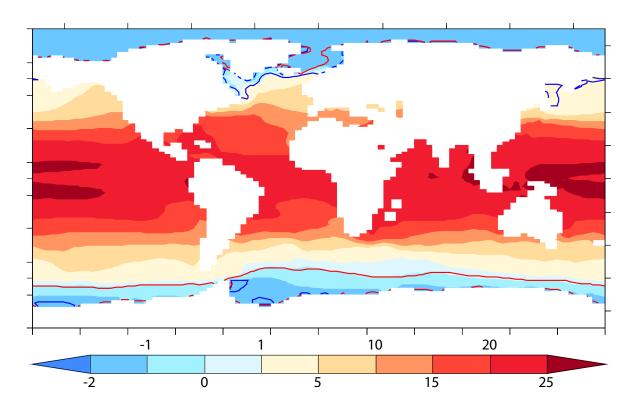
- Figure 1: Ice-sheet and oceanic boundary conditions used for the LMDZ experiments.
 a) Ice-sheet height (light blue scale) and land height (brown scale) given as boundary conditions
 for the five selected experiments out of twelves. Dark blue-brown limits represent the sea-land
- 824 mask, and brown-light blue limits represent the land-ice mask.
- b) Averaged annual sea-surface temperature (shading) and sea-ice extent for boreal summer
- 826 (red contours) and winter (blue contours) from the PMIP3 IPSL-CM5A-LR simulations. The
- averages are computed over 50 years which have been used to create the boundary conditions
- 828 for the LMDZ experiments.
- **Figure 2:** Laurentide ice sheet height for each LMDZ experiment (with name of the corresponding experiment). The height of the LIS is defined as a relative difference between PMIP3 Laurentide ice-sheet height and the present-day topography (Δh_{LIS}^{PMIP3}).
- **Figure3:** 2m summer (left) and winter (right) temperature without the LIS (noIS) and anomaly of the 2-meter temperature between the xxdhL experiments and the noIS experiment: we can therefore quantify the impact of the LIS presence, depending on its topography. The visible surface albedo difference between xxdhL and noIS experiments are represented by orange contours (albedo difference = 0.4) and yellow contours (albedo difference = 0.2). The statistical significance of the patterns of temperature anomalies has been checked using a two-sided
- 838 Student t-test with a p-value =0.05. The white shading indicates areas where the anomaly is not
- 839 significant.
- Figure 4: 500 hPa summer geopotential height anomaly from zonal mean (left panels) and 500
 hPa horizontal wind (arrows) and meridional wind (shaded, right panels).
- **Figure 5:** Same as Fig. 4 for the winter season
- **Figure 6:** Same as Fig. 3 for annual precipitation (left column) and annual snowfall (right column).
- Figure 7: Large-scale (left pannels) and convective (right pannels) annual precipitation without
 LIS (noIS experiment, top panels), and anomalies of large-scale and convective precipitation
 between the xxdhL and noLIS experiments. The statistical significance of the patterns of
 temperature anomalies has been checked using a two-sided Student t-test with a p-value = 0.05.
 The white shading indicates areas where the anomaly is not significant.
- Figure 8: (a) Scandinavian and (b) Barents-Kara precipitation anomalies (xxdhL noIS)
 during winter (circles) and summer (triangles) versus annual precipitation anomaly. The lower
 the Laurentide ice sheet altitude, the lighter grey the symbols are.
- **Figure 9:** Left column: same as Fig. 3 for winter precipitation (shaded) and position of the North Atlantic jet stream without LIS (dashed) and for the xxdhL experimenst (solid line).

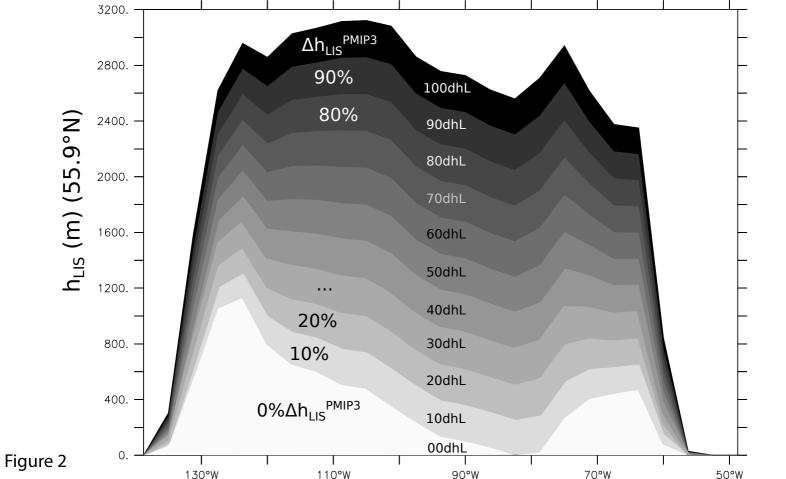
- Right column: summer precipitation (shaded) and snowfall (contour every 0.1 mm/day, blue contour for positive values).
- **Figure 10:** 850 hPa winter jet-stream shift (in latitude) versus winter precipitation anomaly over the Scandinavian region. The range of results has been obtained by bootstrapping. Up and down triangles represent the 5th precipitation percentile versus the 95th jet position percentile, and the 95th precipitation percentile versus the 5th jet position percentile respectively. Circles are the average values and crosses indicate the 5th and 95th percentiles.
- Figure 11: Meridional cross sections of the 500hPa geopotential height anomalie from zonal
 means over the East Laurentide (90°W-60°W, top panel) and the Atlantic sectors (50°W-10°W,
 middle) in meters, and of the 500hPa zonal wind (in m/s) for the Atlantic sector.
- **Figure 12:** (a) Snow accumulation computed by the GRISLI ice-sheet model for the noIS experiment (top panel), and differences of snow accumulation between the xxdhL and the noIS experiments (in m/yr). (b) same as (a) for the ablation (in m/yr).
- Figure 13: (a) Surface mass balance at the beginning of the simulation (in m/yr) and (b) surface elevation at the end of the simulation computed by the ice-sheet model for the selected experiments (in meters).
- Figure 14: Summer 2m temperatures for the a) 00dhL (flat LIS, no FIS) experiment, b)
 00dhLFIS (flat LIS, high FIS), c) 100dhL (high LIS, no FIS) and d)100dhLFIS (high LIS, high
 FIS). The extents of the ice sheets (blue contours) are superimposed on right panels for
 00dhLFIS and 100dhLFIS runs.
- **Figure 15:** Same as Fig. 14 for the vertically integrated meridional (shaded) and the horizontal heat transport (arrows) in summer. Numerical values are given in 10^{10} J m⁻¹ s⁻¹.
- Figure 16: Same as Fig. 14 for the 850hPa meridional wind (shaded) and the 850hPa horizontal
 winds (arrows) for the summer season (in m/s).
- Figure 17: Differences in snow accumulation (in m/yr) computed by the GRISLI ice-sheet
 model between a) 00dhLFIS and 00dhL experiments b) 00dhLFIS and 100dhL experiments, c)
 100dhLFIS and 100dhL experiments and d) 00dhLFIS and 100dhLFIS experiments.
- **Figure 18:** Same as Fig 14 for the precipitation (mm/day) and 850 hPa horizontal winds in winter.
- Figure 19: Difference of simulated ice thickness of the Eurasian ice sheet between 00dhLFISand 100dhLFIS after 100 kyr (in meters).

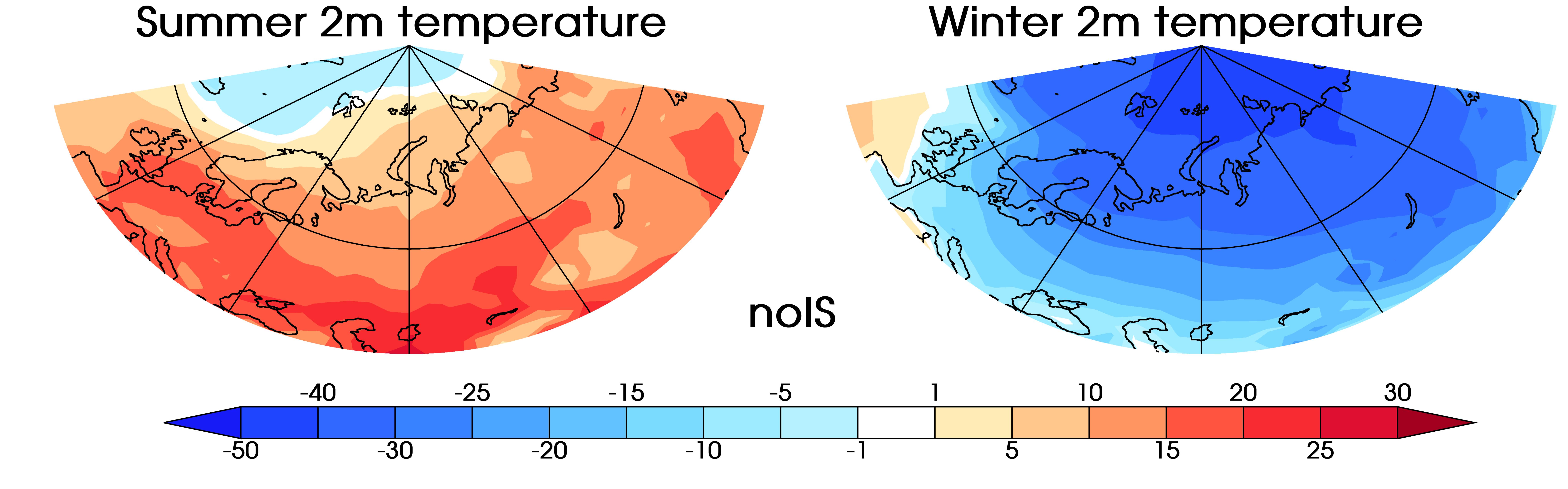
a) Ice-sheet boundary conditions



b) Annual SSTs and sea-ice extent averaged over 50 years








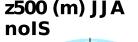
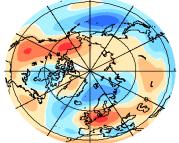
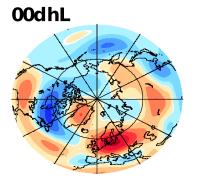
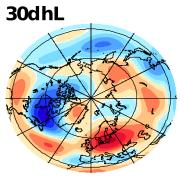


Figure 3

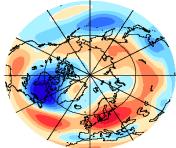








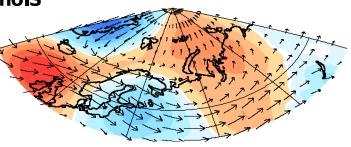
60dhL

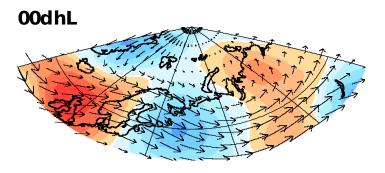


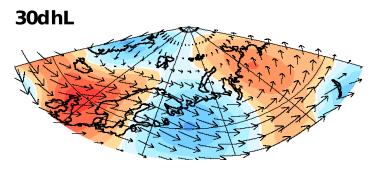
100dhL

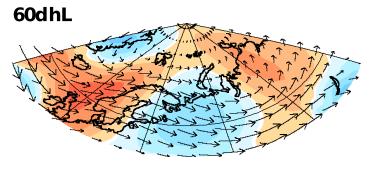


500hPa meridional wind (m/s) and 500 hPa winds noIS









100dhL

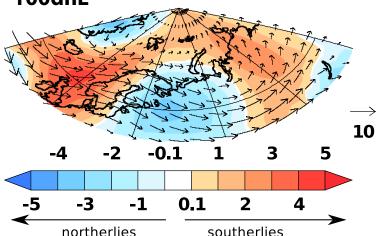
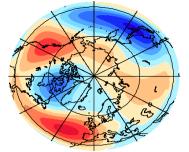
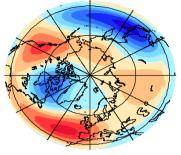


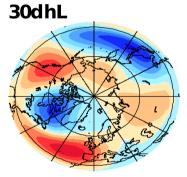
Figure 4

z500 (m) DJ F noIS

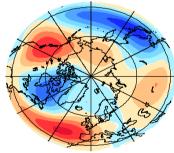


00dhL

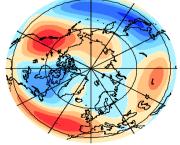


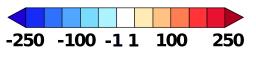


60dhL

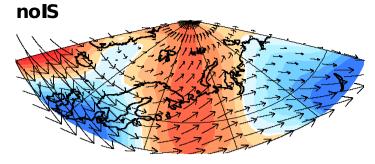


100dhL

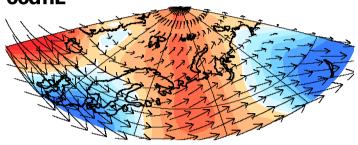


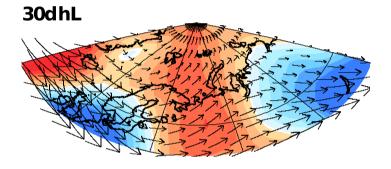


500hPa meridional wind (m/s) and 500 hPa winds

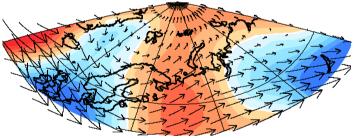


00dhL

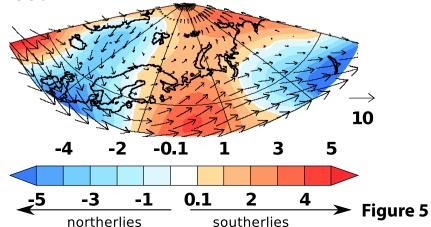


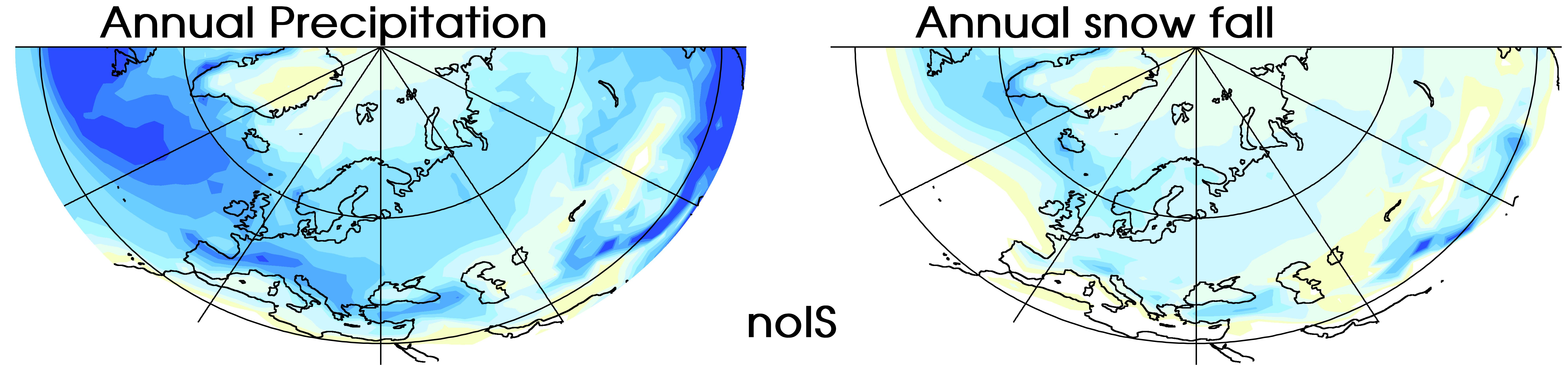


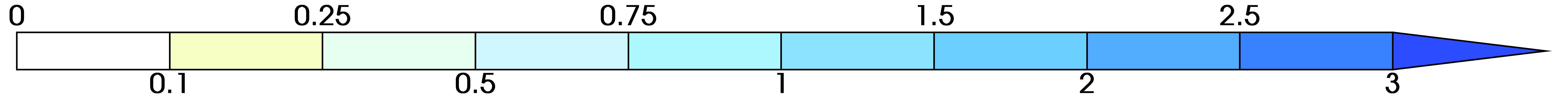
60dhL

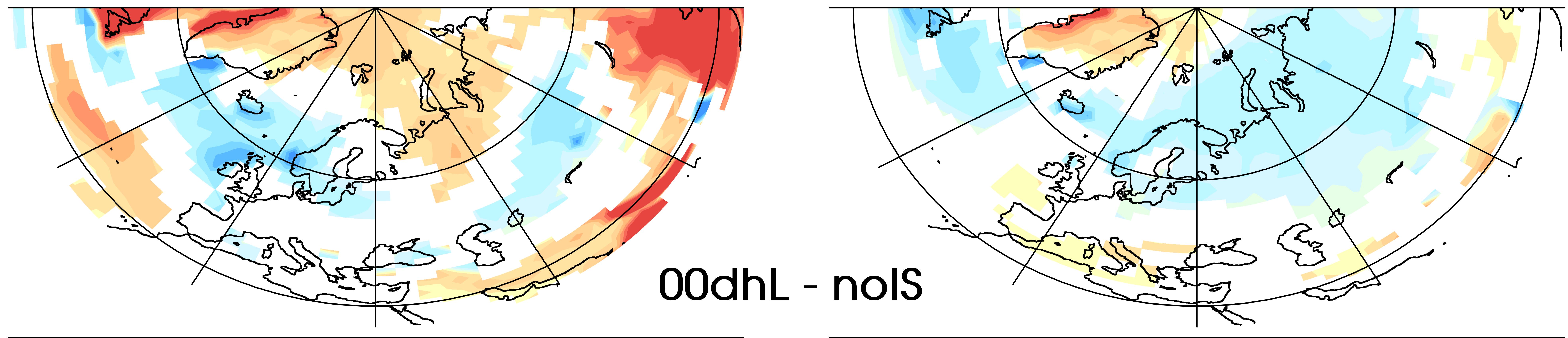



100dhL












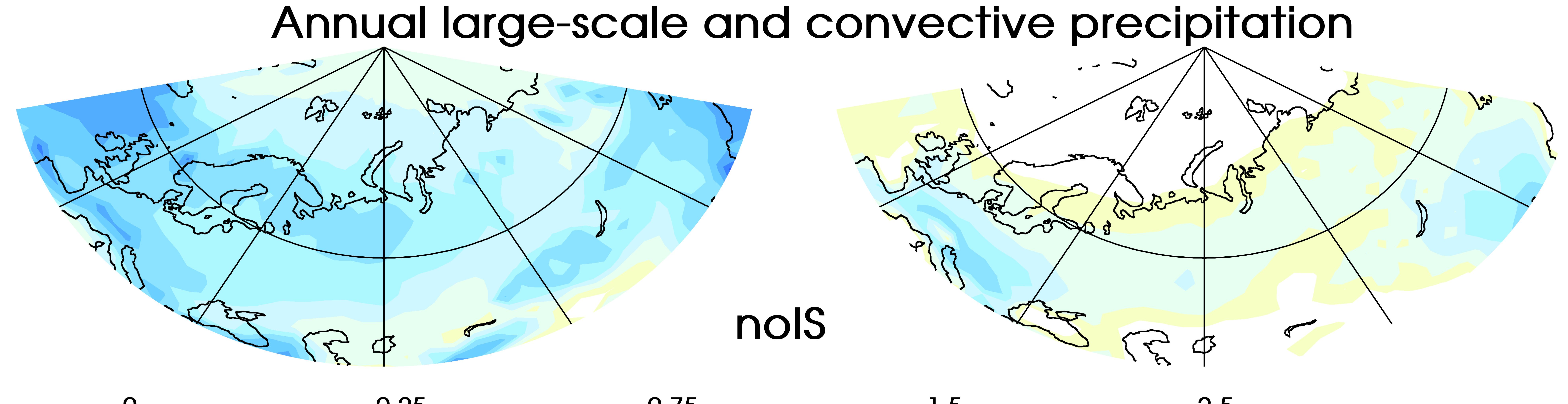


Figure 6



O	0.25		0.75	1.5	2	5
	0.1	0.5			2	3




Figure 7

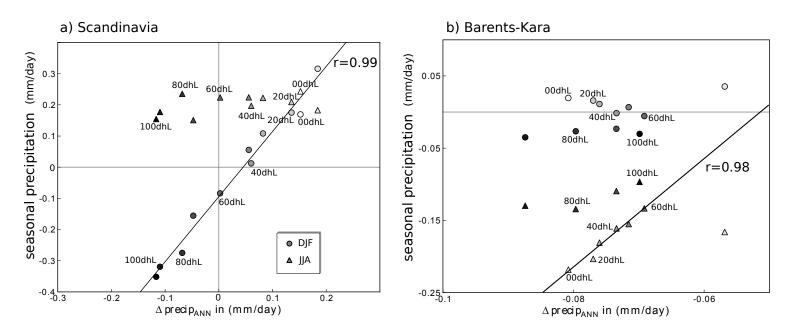
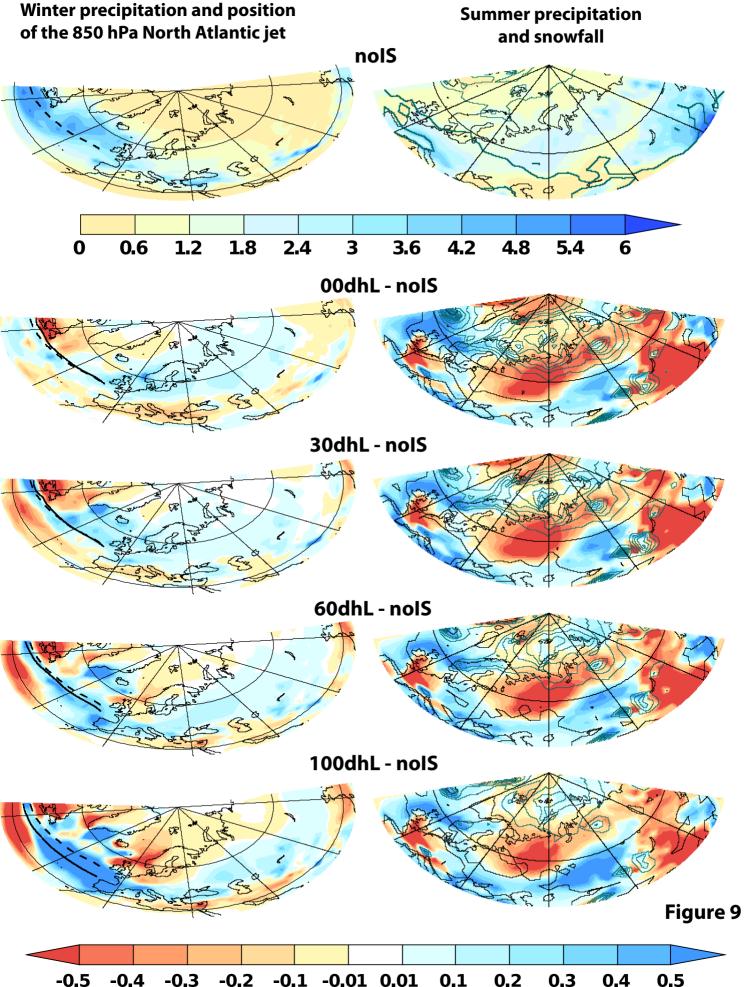
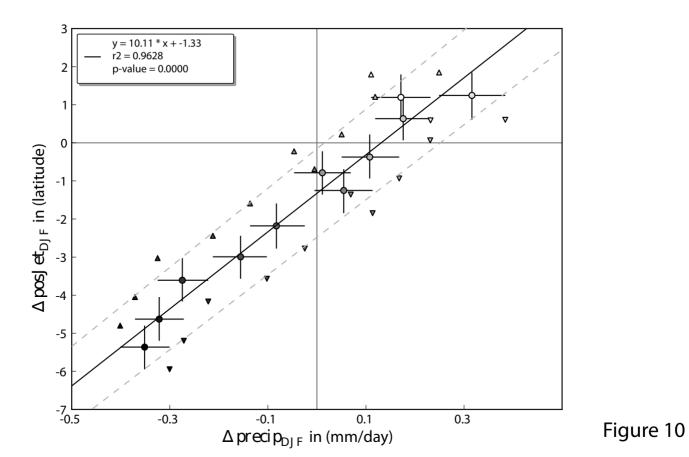


Figure 8





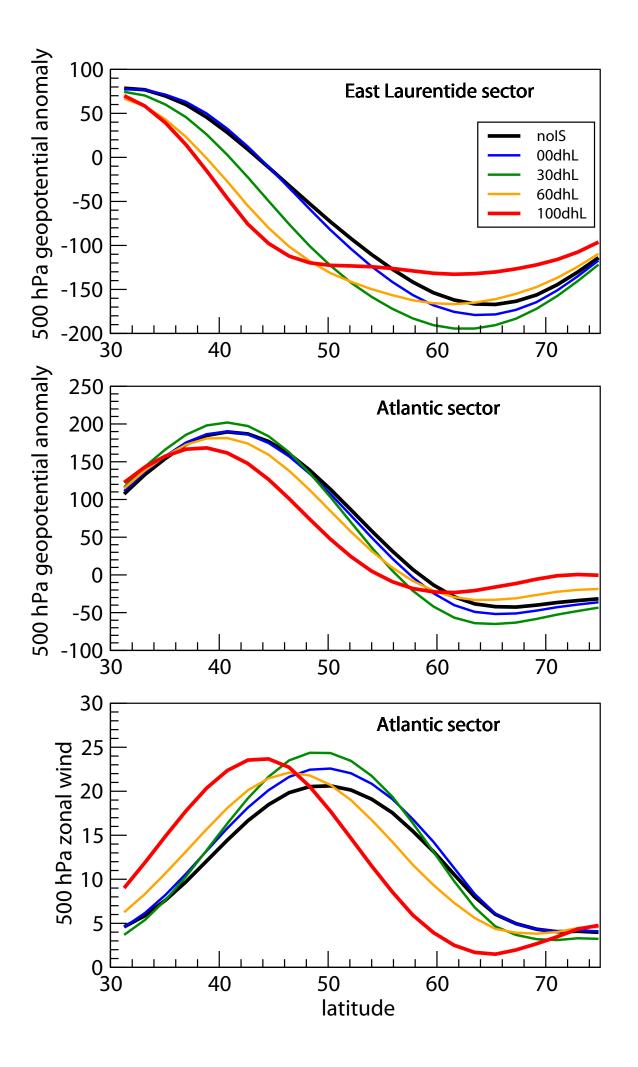
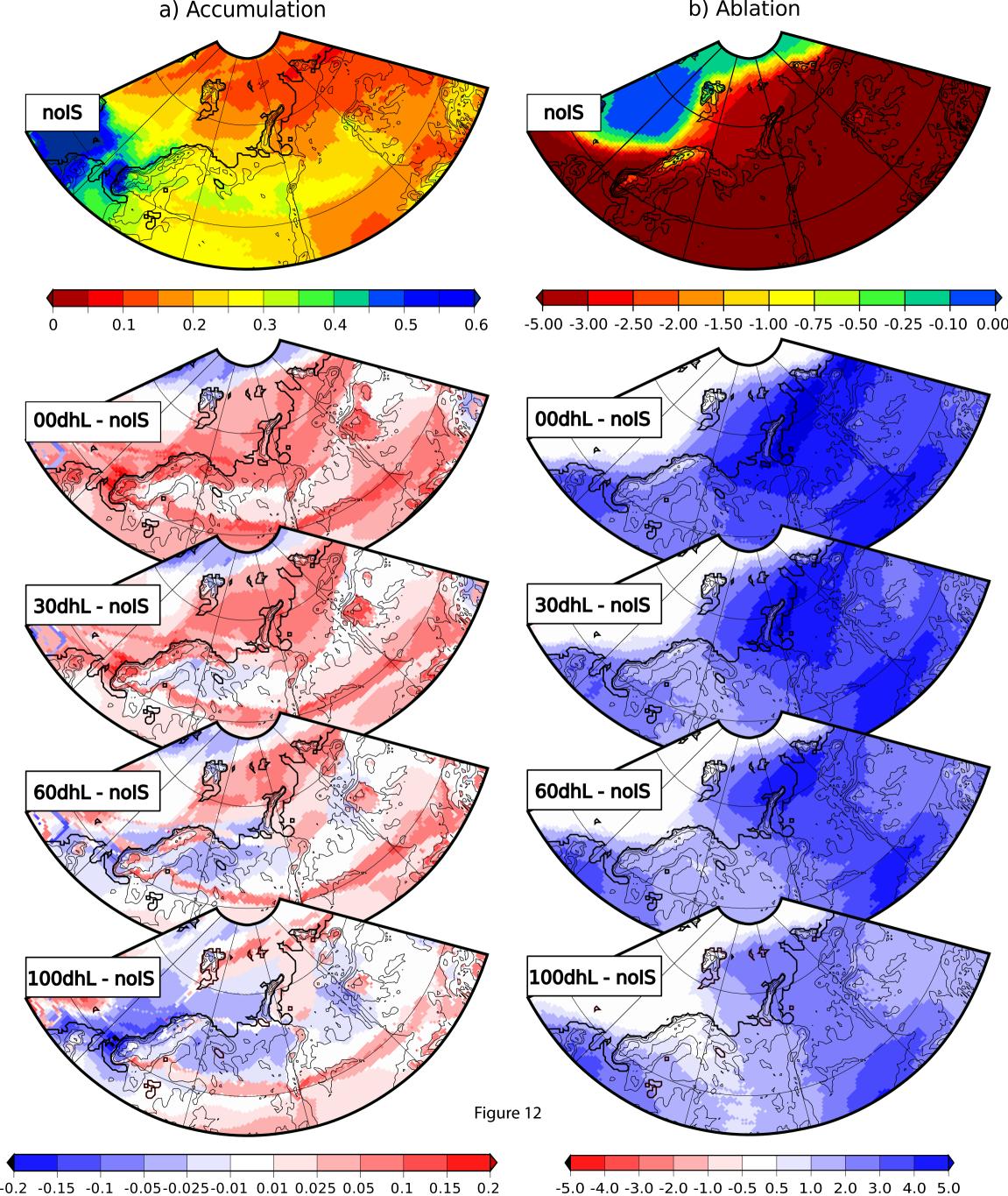
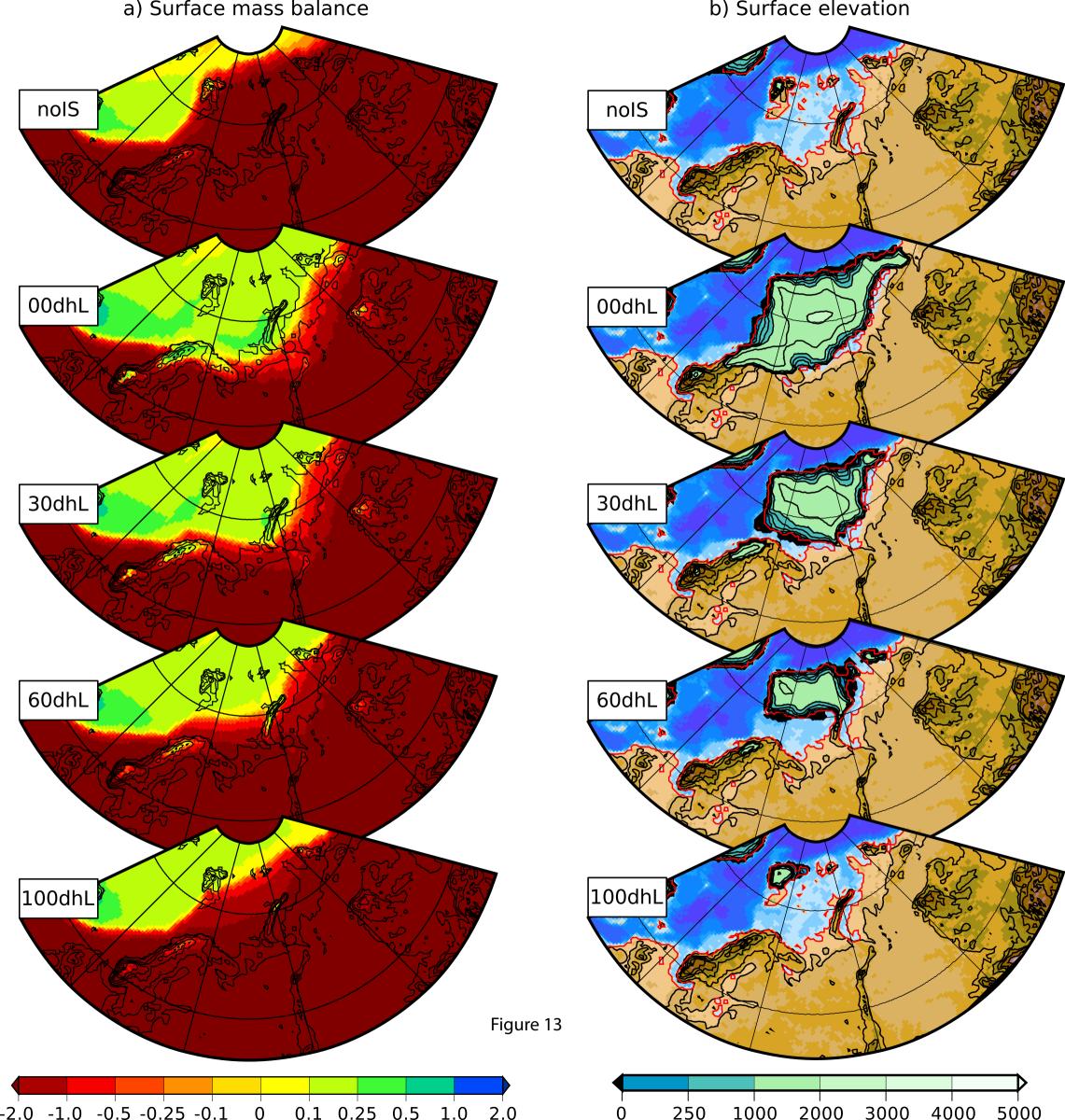
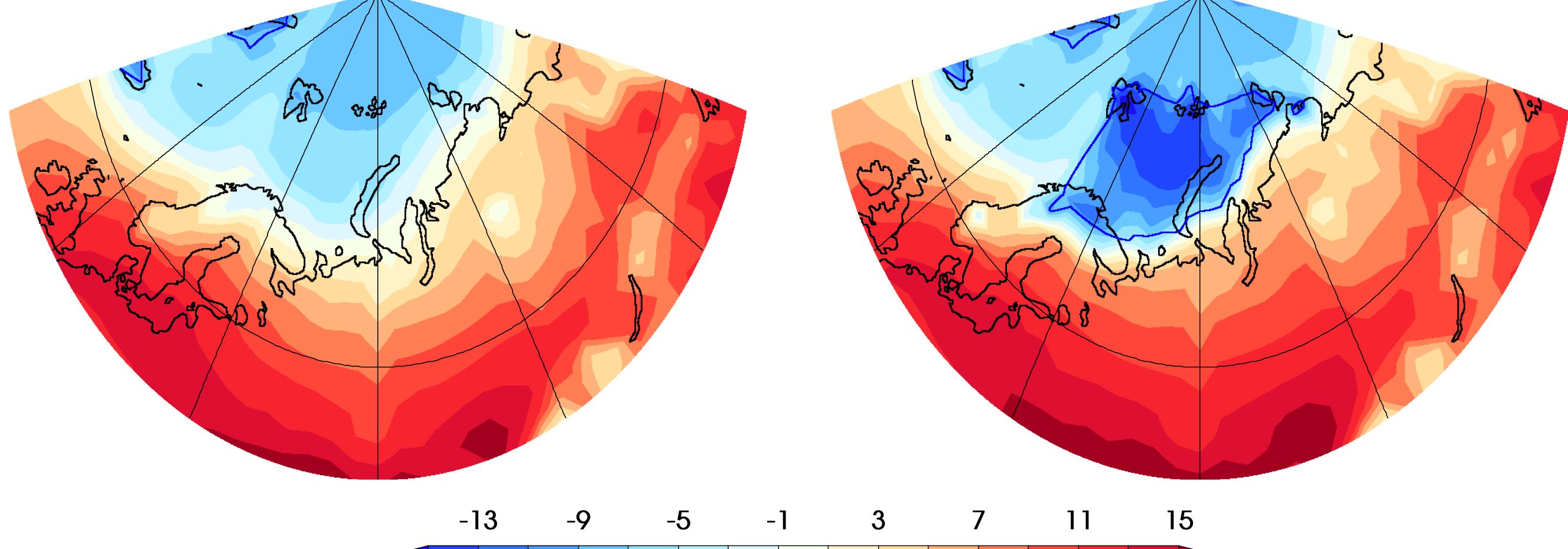


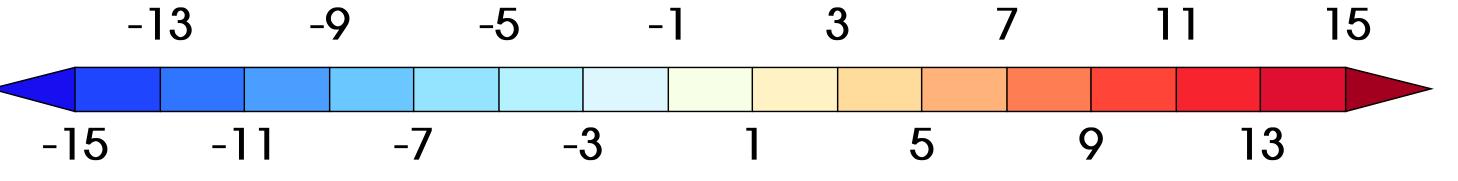
Figure 11





2m temperature a) 00dhL





c) 100dhL

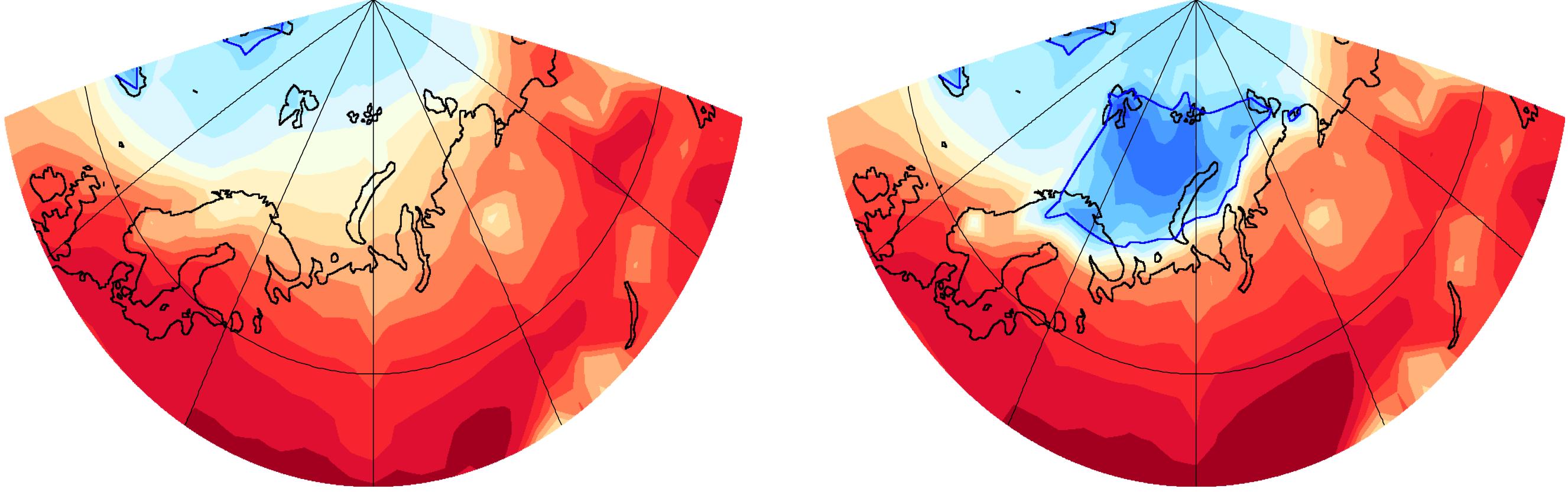
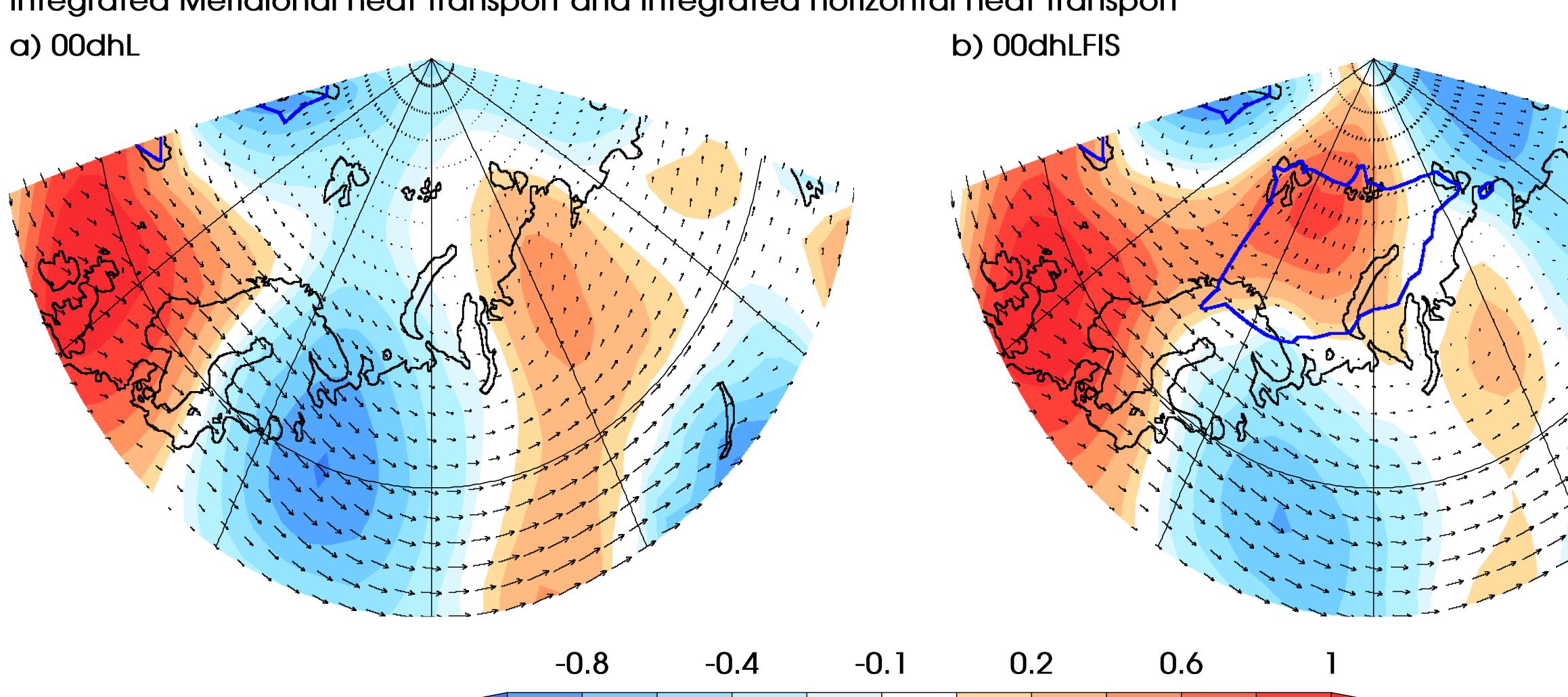


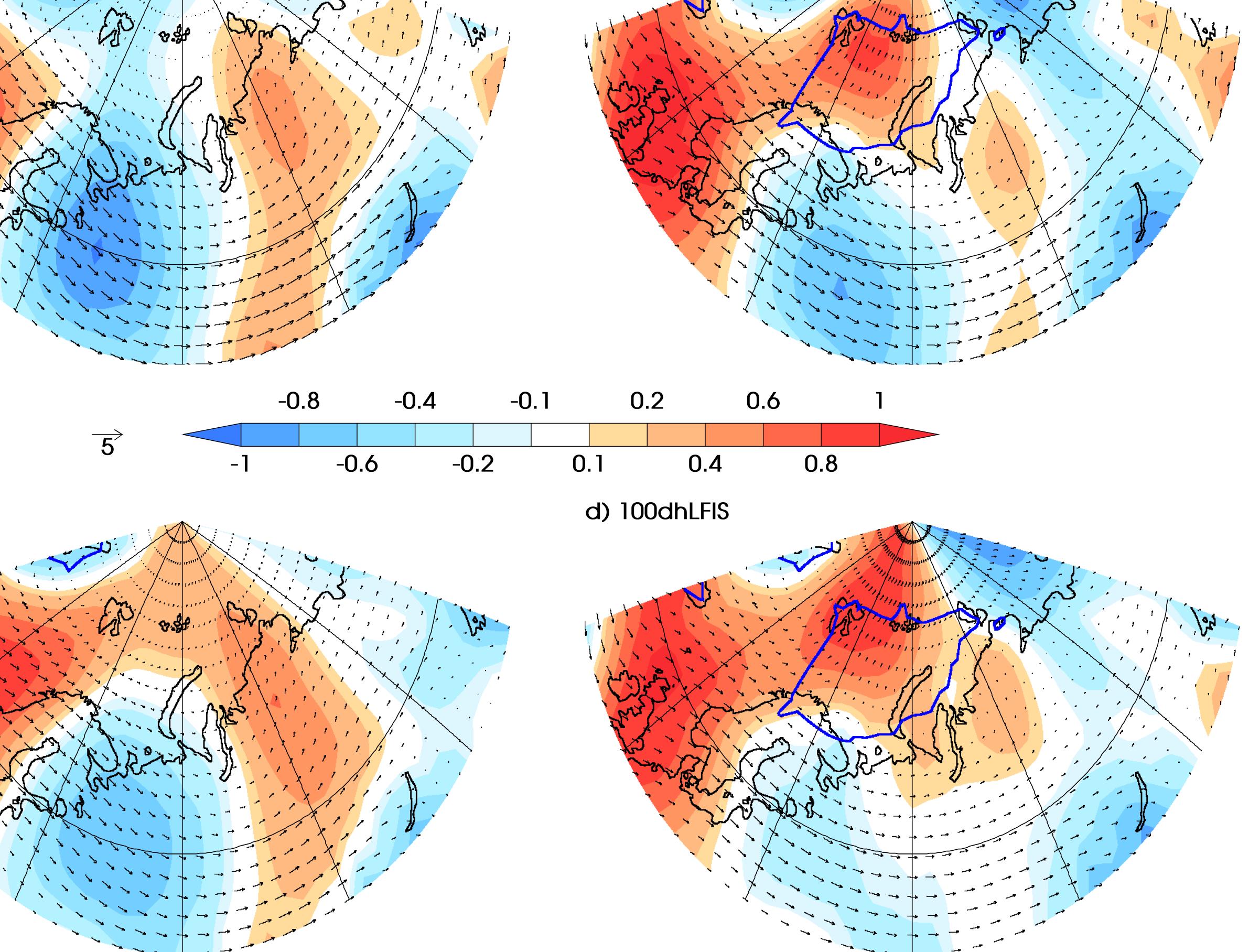
Figure 14

b) 00dhLFIS



Integrated Meridional heat transport and Integrated horizontal heat transport





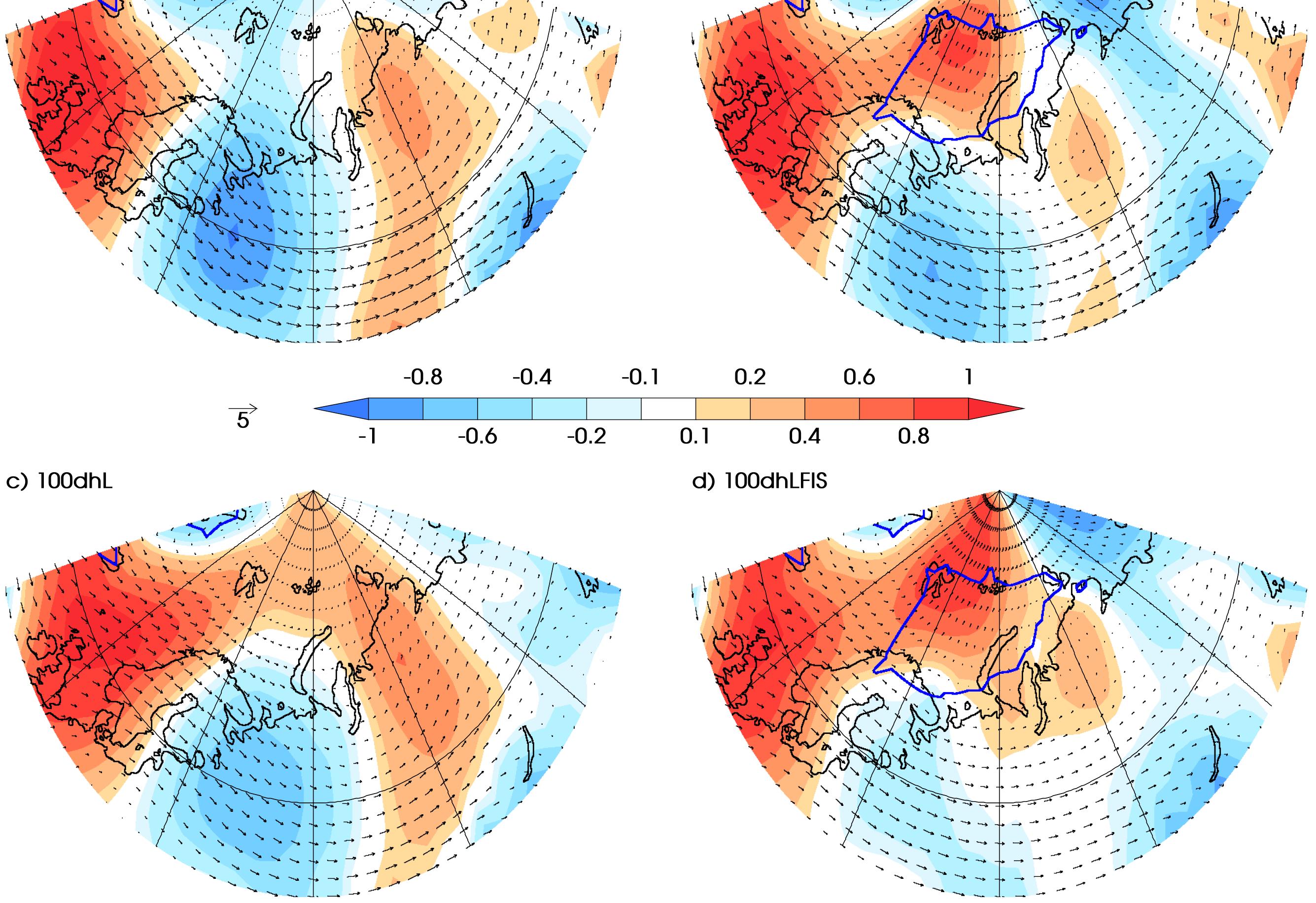
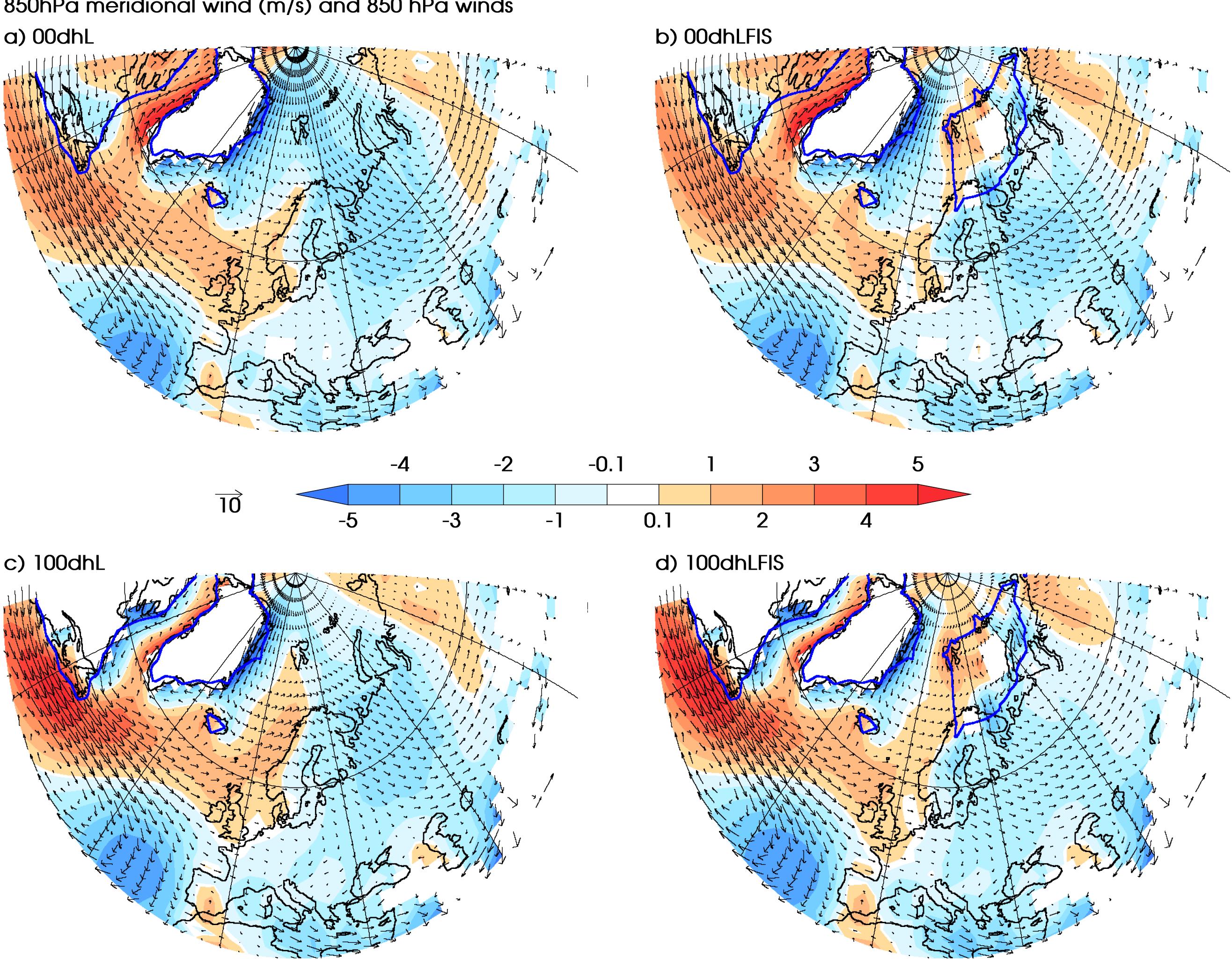
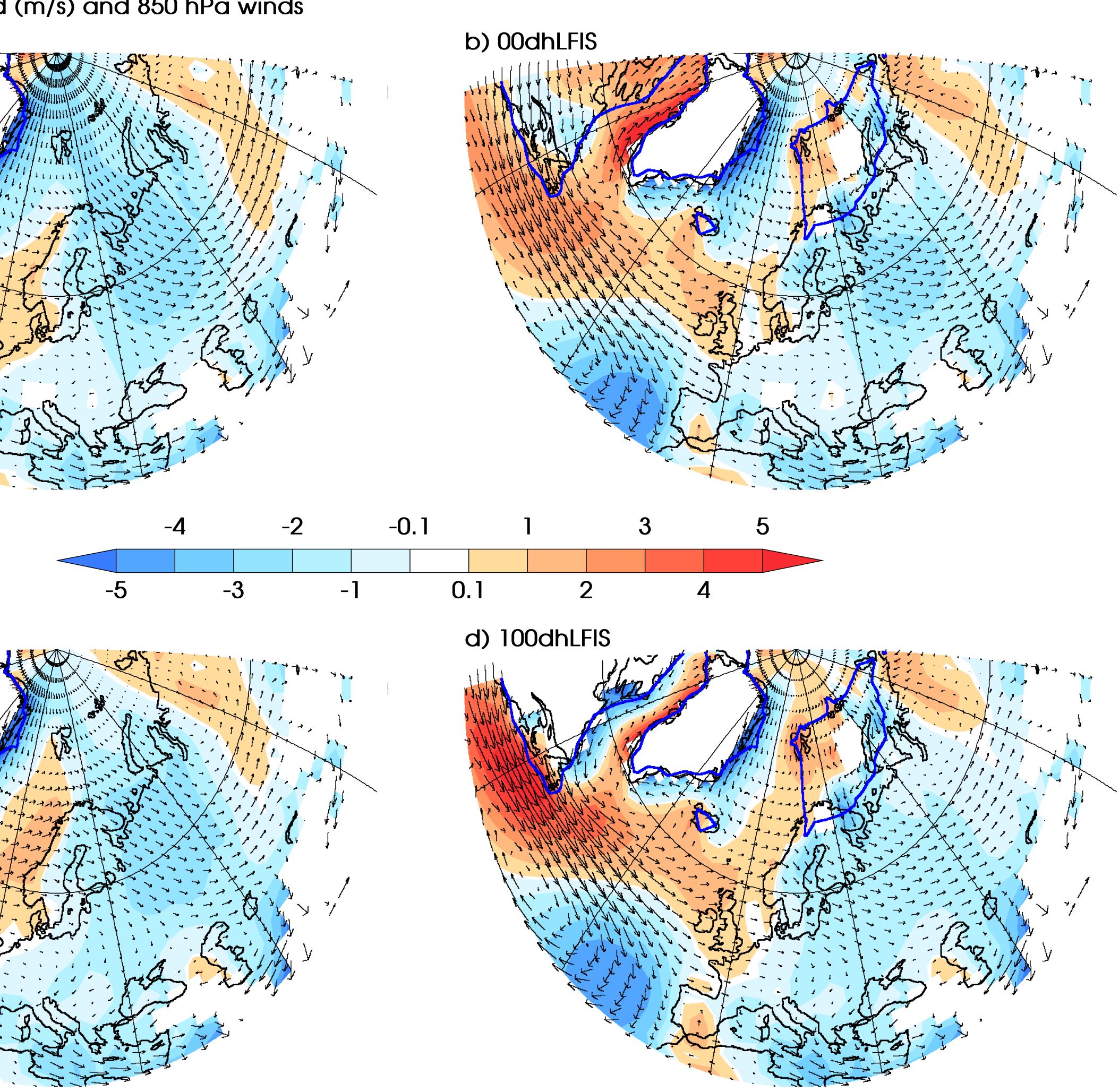


Figure 15

850hPa meridional wind (m/s) and 850 hPa winds





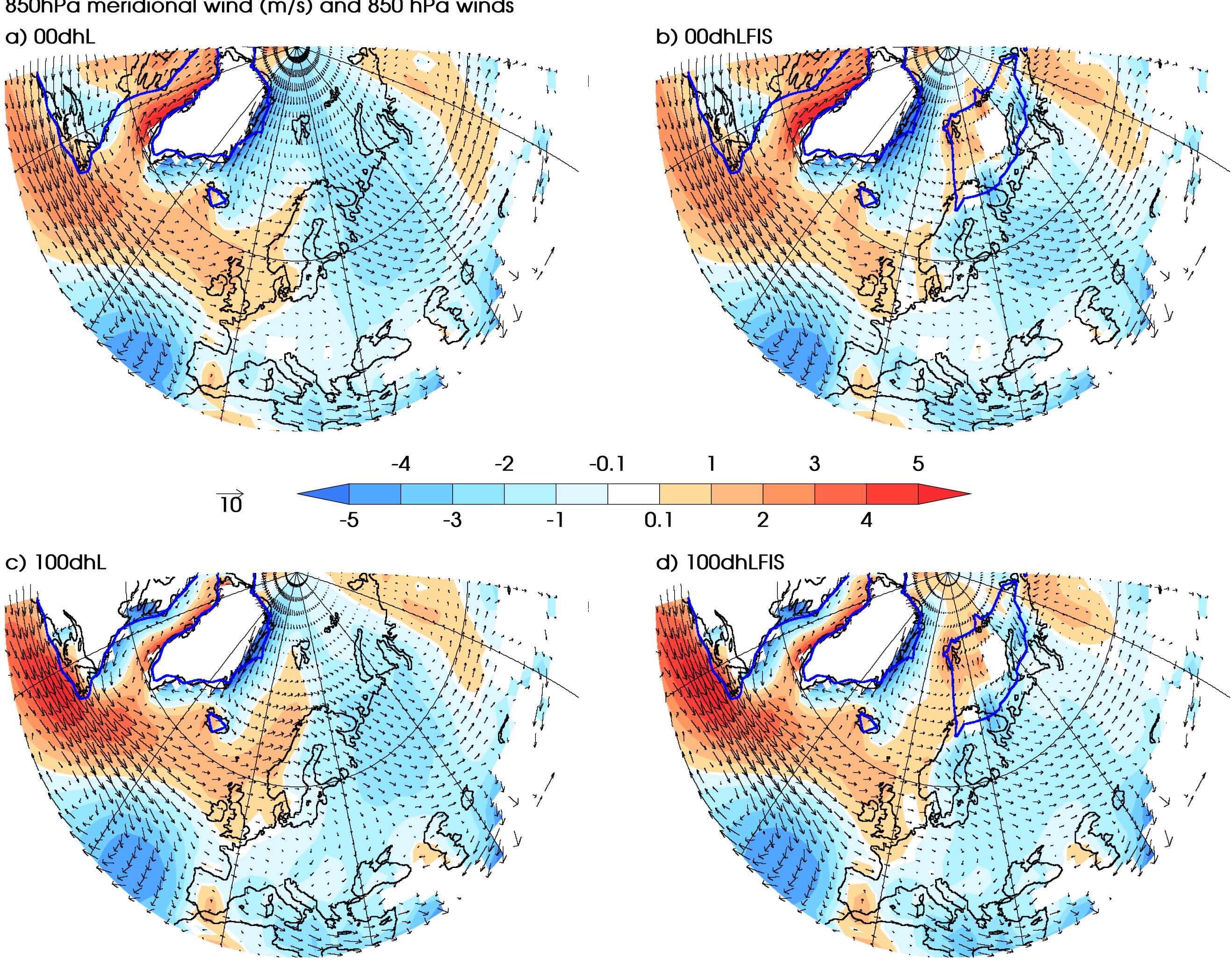


Figure 16

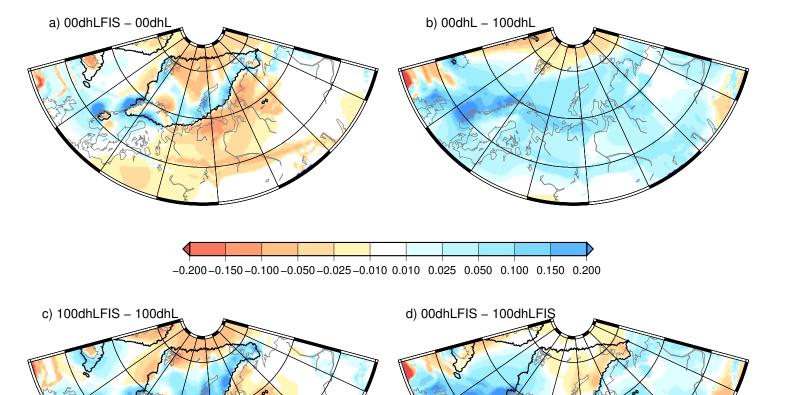
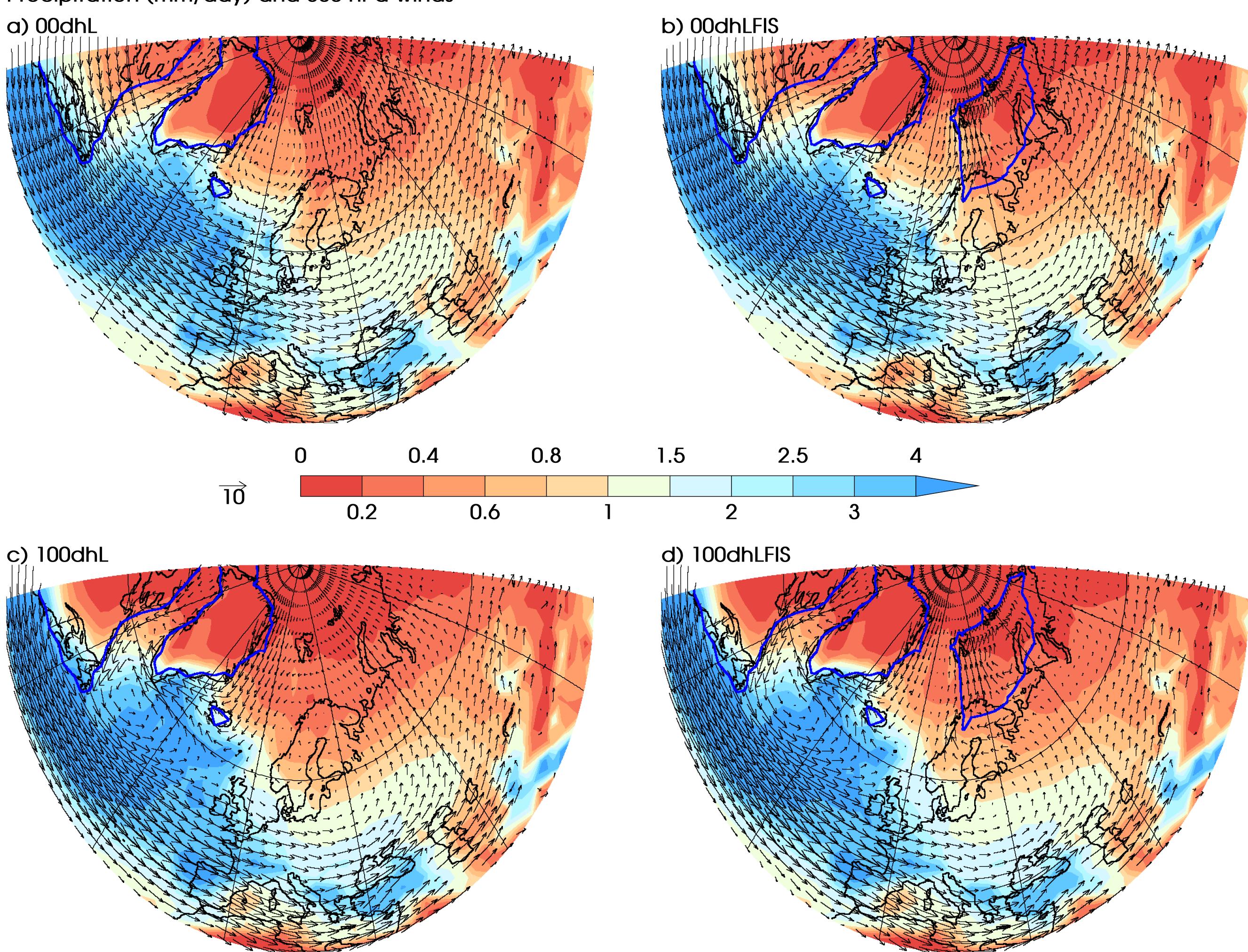
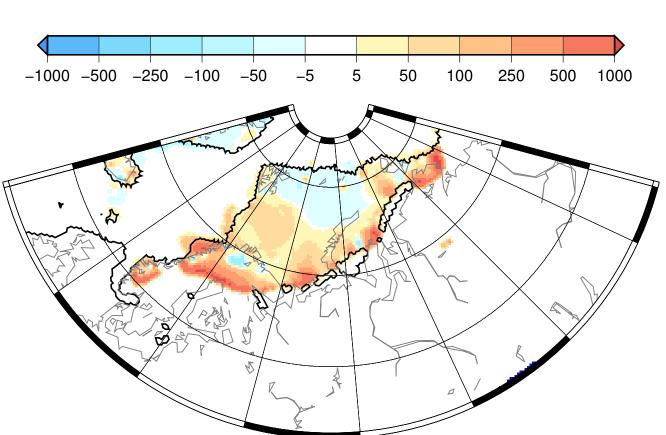


Figure 17

1

Precipitation (mm/day) and 850 hPa winds





00dhLFIS – 100dhLFIS